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AFWL-TR-67-131, VOL II

AFWL-TR-
67-131,
VOL II

NUCLEAR EXPLOSION INTERACTION STUDIES

Volume II

Two-Dimensional Code Development

J. R. Triplett et al.

Gulf General Atomic Incorporated
San Diego, California 92112
Contract No. F29601-67-C-0014



TECHNICAL REPORT NO. AFWL-TR-67-131, VOL II

April 1968



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Air Force Systems Command
Kirtland Air Force Base
New Mexico

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FOREWORD

This report was prepared by Gulf General Atomic Incorporated, San Diego, California, under Contract F29601-67-C-0014. The research was funded by DASA under Project 5710, Subtask 07.017, Program Element 6.16.46.01H, and by ARPA Order 313, Program Element 6.25.03.01R.

Inclusive dates of research were 29 September 1966 to 27 October 1967. The report was submitted 13 March 1968 by the Air Force Weapons Laboratory Project Officer, Major John Bode (WLRT).

This report is published in four volumes: Volume I, Laser Phenomenology (classified CONFIDENTIAL); Volume II, Two-Dimensional Code Development; Volume III, The OUTPUT Code; and Volume IV, Material Property Codes. The first volume contains a classified report on interaction of laser radiation with solid targets and a brief description of calculations done in conjunction with experiments at the Air Force Weapons Laboratory. The remaining three volumes contain reports of code development efforts in the areas of radiative transfer, hydrodynamics, radiative absorption coefficients, and equations of state.

The projects described in this report are for the most part in an incomplete state of development. This is due in part to the nature of the existing computer programs themselves, which continue in a state of development as long as they are in use, and in part to the time scale involved in bringing new programs to a state of capability for solving real problems.


Gulf General Atomic staff personnel responsible for the direction of the research include J. H. Alexander, R. Brightman, R. S. Englemore, B. E. Freeman, W. B. Lindley, L. Norris, J. T. Palmer, L. M. Schalit, J. R. Triplett, and Mrs. Chris Imes. Contractor's report number is GA-7764, Vol II.


The cooperation of Dr. P. V. Avizonis, Major J. Bode, Capt C. C. David, Major G. Spillman, and Lt L. Stoessel of AFWL is gratefully acknowledged.


Other documents produced under this contract are: GAMD-7592, "A Numerical Scheme for First-Order Compton Scattering," J. T. Palmer, December 13, 1966; GAMD-7846, "Difference Equations for Heat Flow in Two Dimensions," J. R. Triplett, March 2, 1967; GAMD-7879, "A Modified Characteristic Method for Radiative Transfer," J. R. Triplett, March 17, 1967; GAMD-7889, "R D C D. A FORTRAN Input Routine," J. H. Alexander, March 24, 1967; GAMD-8333, "Hydrodynamic Equations

for Multidimensional Problems," J. R. Triplett, October 24, 1967; GAMD-8379, "A Brief Study of the Thermodynamic Properties of Several Low Z Elements at Low Temperature," L. M. Schalit, November 22, 1967.

This technical report has been reviewed and is approved.


JOHN BODE
Major, USAF
Project Officer


TRUMAN L. FRANKLIN
Colonel, USAF
Chief, Theoretical Branch


CLAUDE K. STAMBAUGH
Colonel, USAF
Chief, Research Division

ABSTRACT

(Distribution Limitation Statement No. 2)

The HECTIC code is a two-dimensional Eulerian radiative hydrodynamics code designed for nuclear and laser phenomenology applications involving primarily vapor-phase materials, with provisions for heat conduction and vaporization of condensed matter. In this report, a formulation is presented for an improved method of solving the hydrodynamic equations. Solution of the radiative transfer equations by the long-characteristic method is discussed, and computer codes utilizing this approach are presented. The short-characteristic method is also discussed. The nonequilibrium diffusion approximation in two dimensions is considered, and a study of algorithms for solving the implicit difference equations which arise is reported. Experimental codes utilizing the nonequilibrium diffusion and short-characteristic methods are presented.

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SECTION I

HYDRODYNAMIC EQUATIONS FOR MULTIDIMENSIONAL PROBLEMS

The formulation of a computer method for solving the equations governing the flow of a compressible viscous fluid is presented in this section. The aims of the section are essentially to summarize and generalize currently used methods for treating multidimensional fluid flow, and to write down a recipe for computer code development.

The equations are introduced in a more general form than usual, namely for an arbitrary coordinate system. The special case of cylindrical coordinates is then described, and finally the problem is restricted to the axially symmetric case. Specific developments beyond the material in references 1, 2, and 3 include: (1) the representation of the viscous stress components in cylindrical coordinates; (2) the detailed specification of the Lax-Wendroff-Burstein pseudoviscosity terms in cylindrical geometry; (3) a formulation of the two-step Lax-Wendroff difference equations for cylindrical geometry which differs somewhat from that given, for example, by Burstein (Ref. 3); (4) a number of alternative formulations of the differential equations which may prove helpful in the treatment of specific problems; and (5) a suggested special device for adjusting the coefficient of the pseudoviscosity terms.

GLOSSARY

- A Arbitrary second-rank tensor field (Eq. (30)); coefficient of κ^2 in Eq. (100)
- a Arbitrary vector field (Eqs. (28), (29))
- a_0, a_1, a_2 Coefficients of the expansion of Q in powers of J (Eqs. (68), (69), et seq.)
- B Coefficient of κ in Eq. (100)
- C Constant term in Eq. (100)
- c Adiabatic sound speed (Eq. (16) et seq.)
- D/Dt Lagrangian time derivative operator (Eq. (10) et seq.)
- E Specific internal energy, ergs/gram (Eq. (4) et seq.)
- e Total fluid energy density, ergs/cm³ (Eqs. (3), (4), et seq.)
- F Terms whose r-derivatives appear in the flow equations (Eqs. (39), (40), (47), (48), (88), et seq.), or terms whose divergence appears in the flow equations (Eqs. (55 through (81))
- G Terms whose y-derivatives appear in the flow equations (Eqs. (34), (42), (47), (48), et seq.)
- g The metric tensor (or, without indices, its determinant). Values for cylindrical coordinates are given in Eqs. (32), (33), (34).
- H Terms (arising from the use of non-Cartesian coordinates) which are undifferentiated in the flow equations (Eqs. (39), (43), (47), (48), et seq.)
- I Terms whose θ -derivatives appear in the flow equations (Eqs. (39), (41)); the unit matrix (Eqs. (64) through (82))

i	A spatial index; after Eq. (85), the radial mesh index
J	The Jacobian of F with respect to U (Eq. (57), et seq.)
j	A spatial index; after Eq. (86), the axial mesh index
K	Specific kinetic energy, ergs/gram (Eqs. (4), (19), (24))
k	A spatial index
m	A spatial index (Eq. (30)); momentum density (Eqs. (56) through (59), (83), (84)); radial momentum density (Eqs. (50) through (54), (99), (100))
n	Axial momentum density (Eqs. (50) through (54), (99), (100)); time index (Eq. (73), et seq.)
P	Thermodynamic pressure, ergs/cm ³ (Eqs. (3), (6), et seq.)
Q	The Lax-Wendroff pseudoviscosity matrix (Eqs. (64), (68)); the Lapidus pseudoviscosity matrix (Eq. (82))
q	With superscript, total heat flux component (Eq. (3), et seq.)
q ₁ , q ₂ , q ₃	Eigenvalues of Q (Eqs. (64), (65), et seq.)
r	Radial coordinate, cm
S	Specific entropy (Eq. (13), et seq.)
T	Temperature (Eq. (12), et seq.)
t	Time, sec
U	Densities of mass, momentum, and total energy (Eqs. (39), (40), (47), (48), et seq.)
u	With superscript, the fluid velocity (superscript is suppressed in Eqs. (64) through (69)); elsewhere, without superscript, the radial component of fluid velocity (Eq. (37), et seq.)
V	Terms involving viscous stress whose divergence appears in the flow equations (Eq. (83))
v	Axial component of fluid velocity (Eq. (37), et seq.)
w	Aximuthal component of fluid velocity (Eq. (37), et seq.)

x	Arbitrary spatial coordinate, cm
y	Axial coordinate, cm
α, β	Arbitrary unit vector fields (Eqs. (26), (27)); spatial indices (Eqs. (83), (84))
Γ	Derivative of pressure with respect to internal energy density at constant volume, dimensionless (Eq. (17), et seq.); with three indices, a Christoffel symbol (Eqs. (28) through (36))
Δ, ∇	Forward and backward difference operators
δ	Kronecker symbol (Eqs. (31), (57), (62)); central difference operator (Eq. (85), et seq.)
θ	Azimuthal coordinate, dimensionless
κ	Coefficient of pseudoviscosity, dimensionless (Eq. (65), et seq.)
μ	Shear viscosity, gm/(cm sec) (Eq. (4), et seq.)
μ_B	Bulk viscosity, gm/(cm sec) (Eq. (4), et seq.)
ξ	Basic mesh interval (Eq. (65), et seq.)
τ	Viscous stress (Eqs. (2), (5), et seq.)
ρ	Mass density, gm/cm ³
Φ	Terms involving pseudoviscosity whose divergence appears in the flow equations (Eqs. (66), (67))
ϕ	Thermal dissipation rate, ergs/(cm ³ sec) (Eq. (12) et seq.)
ψ	Viscous dissipation rate, ergs/(cm ³ sec) (Eqs. (9), (11), et seq.)

DIFFERENTIAL EQUATIONS

The following equations express conservation of mass, momentum, and total energy for a viscous fluid in a Euclidean space:

$$\frac{\partial \rho}{\partial t} + \partial_j (\rho u^j) = 0 \quad (1)$$

$$\frac{\partial}{\partial t} (\rho u^i) + \partial_j (\rho u^i u^j + g^{ij} P - \tau^{ij}) = 0 \quad (2)$$

$$\frac{\partial e}{\partial t} + \partial_j [u^j (e + P) - g_{ik} \tau^{kj} u^i + q^j] = 0 \quad (3)$$

where ρ is mass density, e is total energy density, P is thermodynamic pressure, u^i is fluid velocity, q^i is total heat flux, g^{ij} is the metric, and τ^{ij} is the residual viscous stress. It is convenient also to introduce the specific internal energy E , the shear viscosity μ , and the bulk viscosity μ_B , in terms of which one writes the energy relation

$$e = \rho(E + 1/2 g_{ij} u^i u^j) = \rho(E + K) \quad (4)$$

the constitutive relation

$$\tau^{ij} = \mu (g^{ik} u^j_{,k} + g^{jk} u^i_{,k}) + (\mu_B - 2/3 \mu) g^{ij} u^k_{,k} \quad (5)$$

and the state equation

$$P = P(\rho, E) \quad (6)$$

The symbols ∂_j and $,j$ are alternative notations for covariant differentiation with respect to the coordinate x^j .

Alternative forms of these equations are as follows:

$$\frac{D\rho}{Dt} + \rho u^j_{,j} = 0 \quad (7)$$

$$\rho \frac{Du^i}{Dt} + \partial_j (g^{ij} P - \tau^{ij}) = 0 \quad (8)$$

$$\rho \frac{DE}{Dt} + P u^j_{,j} - \Psi + q^j_{,j} = 0 \quad (9)$$

where D/Dt is the Lagrangian time derivative operator

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + u^j \partial_j \quad (10)$$

and the quantity

$$\begin{aligned}\Psi = g_{ik} \tau^{kj} u^i_{,j} = & 2/3 \mu [u^1_{,1} - u^2_{,2}]^2 + (u^2_{,2} - u^3_{,3})^2 + (u^3_{,3} - u^1_{,1})^2] \\ & + \mu [(u^1_{,2} + u^2_{,1})^2 + (u^2_{,3} + u^3_{,2})^2 + (u^3_{,1} + u^1_{,3})^2] \\ & + \mu_B (u^j_{,j})^2\end{aligned}\quad (11)$$

is the viscous dissipation rate. With the introduction of specific entropy S and temperature $T = (\partial E / \partial S)_\rho$, the thermal dissipation rate is

$$\phi = - \frac{1}{T} q^i \frac{\partial T}{\partial x^i} \quad (12)$$

The specific entropy S may be divided into reversible and irreversible parts S_{rev} and S_{irr} , such that

$$\rho T \frac{DS_{rev}}{Dt} = -q^j_{,j} - \phi \quad (13)$$

$$\rho T \frac{DS_{irr}}{Dt} = \Psi + \phi \quad (14)$$

The total heating rate is the sum of these:

$$\rho T \frac{DS}{Dt} = \rho \frac{DE}{Dt} + P u^j_{,j} = \Psi - q^j_{,j} \quad (15)$$

Let the adiabatic sound speed be denoted by c , where

$$c^2 = \left. \frac{\partial P}{\partial \rho} \right|_S \quad (16)$$

and a parameter Γ be defined by

$$\Gamma = \left. \frac{1}{\rho} \frac{\partial P}{\partial E} \right|_\rho = \frac{1}{\rho T} \left. \frac{\partial P}{\partial S} \right|_\rho \quad (17)$$

Then the rate equation for P is

$$\frac{DP}{Dt} + \rho c^2 u^j_{,j} = \Gamma (\Psi - q^j_{,j}) \quad (18)$$

Finally, the rate equation for specific kinetic energy K is

$$\rho \frac{DK}{Dt} + u^j P_{,j} - g_{ij} u^j \tau_{,k}^{ik} = 0 \quad (19)$$

Modified forms of Eqs. (7), (9), (15), (18), and (19) may be written for the logarithms of the positive definite scalars ρ , E , S , P , and K , respectively:

$$\frac{D}{Dt} \log \rho + u^j_{,j} = 0 \quad (20)$$

$$\frac{D}{Dt} \log E + \frac{1}{\rho E} (P u^j_{,j} - \Psi + q^j_{,j}) = 0 \quad (21)$$

$$\frac{D}{Dt} \log S = \frac{1}{\rho T S} (\Psi - q^j_{,j}) \quad (22)$$

$$\frac{D}{Dt} \log P + \frac{\rho c^2}{P} u^j_{,j} = \frac{\Gamma}{P} (\Psi - q^j_{,j}) \quad (23)$$

$$\frac{D}{Dt} \log K + \frac{1}{\rho K} (u^j P_{,j} - g_{ij} u^j \tau_{,k}^{ik}) = 0 \quad (24)$$

The characteristic equations are the following combinations of Eqs. (7), (8), and (18):

$$\frac{DP}{Dt} - c^2 \frac{D\rho}{Dt} = \Gamma(\Psi - q^j_{,j}) \quad (25)$$

$$\alpha_i \left[\frac{Du^i}{Dt} + \frac{1}{\rho} \partial_j (g^{ij} P - \tau^{ij}) \right] = 0 \quad (26)$$

$$\rho c \left[\beta_i \frac{D}{Dt} + c \partial_i \right] u^i + \left[\frac{D}{Dt} + c \beta_i g^{ij} \partial_j \right] P = c \beta_i \tau_{,j}^{ij} + \Gamma(\Psi - q^j_{,j}) \quad (27)$$

where α_i and β_i are arbitrary spatial unit vector fields.

In the remainder of this report, the solution of Eqs. (1) through (6) will be discussed.

CYLINDRICAL COORDINATES

The following formulae of tensor analysis are needed:

$$a^i_{,k} = \frac{\partial}{\partial x^k} a^i + \Gamma^i_{jk} a^j \quad (28)$$

$$a^k_{,k} = \frac{\partial}{\partial x^k} a^k + \Gamma^k_{jk} a^j = \frac{1}{\sqrt{g}} \frac{\partial}{\partial x^k} \sqrt{g} a^k \quad (29)$$

$$A^{ij}_{,j} = \frac{\partial}{\partial x^j} A^{ij} + \Gamma^j_{mj} A^{im} + \Gamma^i_{mj} A^{mj} = \frac{1}{\sqrt{g}} \frac{\partial}{\partial x^j} \sqrt{g} A^{ij} + \Gamma^i_{mj} A^{mj} \quad (30)$$

where g is the metric determinant and Γ^i_{jk} is a Euclidean Christoffel symbol. The following identities are also noted:

$$g_{ij} g^{jk} = \delta^k_i$$

$$g^{ij}_{,k} = 0 \quad (31)$$

$$\delta^i_{j,k} = 0$$

For cylindrical coordinates, let $x^1 = r$, $x^2 = \theta$, $x^3 = z$. Then $\sqrt{g} = r$, and the nonvanishing components of g^{ij} and Γ^i_{jk} are

$$g_{11} = g^{11} = g_{33} = g^{33} = 1 \quad (32)$$

$$g_{22} = r^2 \quad (33)$$

$$g^{22} = \frac{1}{r^2} \quad (34)$$

$$\Gamma^1_{22} = -r \quad (35)$$

$$\Gamma^2_{12} = \Gamma^2_{21} = \frac{1}{r} \quad (36)$$

Let the velocity components be denoted by $u^1 = u$, $u^2 = w$, $u^3 = v$. The

velocity gradient components and the velocity divergence are then

$$\begin{aligned}
 u_{,1}^1 &= \frac{\partial u}{\partial r} & u_{,1}^2 &= \frac{\partial w}{\partial r} + \frac{w}{r} & u_{,1}^3 &= \frac{\partial v}{\partial r} \\
 u_{,2}^1 &= \frac{\partial u}{\partial \theta} - rw & u_{,2}^2 &= \frac{\partial w}{\partial \theta} + \frac{u}{r} & u_{,2}^3 &= \frac{\partial v}{\partial \theta} \\
 u_{,3}^1 &= \frac{\partial u}{\partial y} & u_{,3}^2 &= \frac{\partial w}{\partial y} & u_{,3}^3 &= \frac{\partial v}{\partial y} \\
 u_{,k}^k &= \frac{\partial u}{\partial r} + \frac{u}{r} + \frac{\partial w}{\partial \theta} + \frac{\partial v}{\partial y}
 \end{aligned} \tag{37}$$

The components of the residual stress are

$$\begin{aligned}
 \tau^{11} &= 2\mu \frac{\partial u}{\partial r} + \left(\mu_B - \frac{2}{3}\mu \right) u_{,k}^k \\
 \tau^{22} &= r^{-2} \left[2\mu \left(\frac{\partial w}{\partial \theta} + \frac{u}{r} \right) + \left(\mu_B - \frac{2}{3}\mu \right) u_{,k}^k \right] \\
 \tau^{33} &= 2\mu \frac{\partial v}{\partial y} + \left(\mu_B - \frac{2}{3}\mu \right) u_{,k}^k \\
 \tau^{12} &= \tau^{21} = \mu \left(\frac{\partial w}{\partial r} + r^{-2} \frac{\partial u}{\partial \theta} \right) \\
 \tau^{13} &= \tau^{31} = \mu \left(\frac{\partial v}{\partial r} + \frac{\partial u}{\partial y} \right) \\
 \tau^{23} &= \tau^{32} = \mu \left(r^{-2} \frac{\partial v}{\partial \theta} + \frac{\partial w}{\partial y} \right)
 \end{aligned} \tag{38}$$

The conservation equations, (1), (2), and (3), may now be written in matrix form:

$$\frac{\partial U}{\partial t} + \frac{\partial F}{\partial r} + \frac{\partial I}{\partial \theta} + \frac{\partial G}{\partial y} + H = 0 \tag{39}$$

where

$$U = \begin{bmatrix} \rho \\ \rho u \\ \rho w \\ \rho v \\ e \end{bmatrix} \quad F = \begin{bmatrix} \rho u \\ \rho u^2 + P - \tau^{11} \\ \rho u w - \tau^{21} \\ \rho v u - \tau^{31} \\ u(e + P - \tau^{11}) - r^2 \tau^{21} w - \tau^{31} v + q^1 \end{bmatrix} \quad (40)$$

$$I = \begin{bmatrix} \rho w \\ \rho u w - \tau^{12} \\ \rho w^2 + r^{-2} P - \tau^{22} \\ \rho v w - \tau^{32} \\ w(e + P - r^2 \tau^{22}) - \tau^{12} u - \tau^{32} v + q^2 \end{bmatrix} \quad (41)$$

$$G = \begin{bmatrix} \rho v \\ \rho u v - \tau^{13} \\ \rho w v - \tau^{23} \\ \rho v^2 + P - \tau^{33} \\ v(e + P - \tau^{33}) - \tau^{13} u - \tau^{23} w + q^3 \end{bmatrix} \quad (42)$$

$$H = \frac{1}{r} \begin{bmatrix} \rho u \\ \rho u^2 - \tau^{11} - \rho r^2 w^2 + r^2 \tau^{22} \\ 3(\rho u w - \tau^{12}) \\ \rho v u - \tau^{31} \\ u(e + P - \tau^{11}) - r^2 \tau^{21} w - \tau^{31} v + q^1 \end{bmatrix} \quad (43)$$

For axially symmetric problems without rotation,

$$w = \frac{\partial u}{\partial \theta} = \frac{\partial w}{\partial \theta} = \frac{\partial v}{\partial \theta} = \frac{\partial P}{\partial \theta} = 0 \quad (44)$$

The equations then simplify to

$$\begin{aligned}
 u_{,1}^1 &= \frac{\partial u}{\partial r} & u_{,1}^3 &= \frac{\partial v}{\partial r} \\
 u_{,2}^2 &= \frac{u}{r} & u_{,3}^1 &= \frac{\partial u}{\partial y} \\
 u_{,3}^3 &= \frac{\partial v}{\partial y} & u_{,k}^k &= \frac{\partial u}{\partial r} + \frac{u}{r} + \frac{\partial v}{\partial y}
 \end{aligned} \tag{45}$$

and

$$\begin{aligned}
 \tau^{11} &= 2\mu \frac{\partial u}{\partial r} + \left(\mu_B - \frac{2}{3}\mu \right) u_{,k}^k \\
 \tau^{22} &= r^{-2} \left[\frac{2\mu u}{r} + \left(\mu_B - \frac{2}{3}\mu \right) u_{,k}^k \right] \\
 \tau^{33} &= 2\mu \frac{\partial v}{\partial y} + \left(\mu_B - \frac{2}{3}\mu \right) u_{,k}^k \\
 \tau^{13} &= \tau^{31} = \mu \left(\frac{\partial v}{\partial r} + \frac{\partial u}{\partial y} \right)
 \end{aligned} \tag{46}$$

and

$$\frac{\partial U}{\partial t} + \frac{\partial F}{\partial r} + \frac{\partial G}{\partial y} + H = 0 \tag{47}$$

with

$$\begin{aligned}
 U &= \begin{bmatrix} \rho \\ \rho u \\ \rho v \\ e \end{bmatrix} & F &= \begin{bmatrix} \rho u \\ \rho u^2 + P - \tau^{11} \\ \rho vu - \tau^{31} \\ u(e + P - \tau^{11}) - \tau^{31}v + q \end{bmatrix} \\
 G &= \begin{bmatrix} \rho v \\ \rho uv - \tau^{13} \\ \rho v^2 + P - \tau^{33} \\ v(e + P - \tau^{33}) - \tau^{13}u + q \end{bmatrix} & H &= \frac{1}{r} \begin{bmatrix} \rho u \\ \rho u^2 - \tau^{11} + r^2 \tau^{22} \\ \rho vu - \tau^{31} \\ u(e + P - \tau^{11}) - \tau^{31}v + q \end{bmatrix}
 \end{aligned} \tag{48}$$

From Eqs. (16) and (17),

$$\left. \frac{\partial P}{\partial E} \right|_{\rho} = \rho \Gamma \quad \left. \frac{\partial P}{\partial \rho} \right|_E = c^2 - \Gamma \frac{F}{\rho} \quad (49)$$

Let $m = \rho u$ and $n = \rho v$; then, from Eq. (4),

$$E = \frac{e}{\rho} - \frac{m^2 + n^2}{2\rho^2} \quad (50)$$

Therefore,

$$\begin{aligned} \left. \frac{\partial P}{\partial \rho} \right|_{m, n, e} &= \left. \frac{\partial P}{\partial \rho} \right|_E + \left. \frac{\partial P}{\partial E} \right|_{\rho} \left. \frac{\partial E}{\partial \rho} \right|_{m, n, e} \\ &= c^2 - \frac{\Gamma}{\rho} (P + e) + \Gamma(u^2 + v^2) \end{aligned} \quad (51)$$

$$\left. \frac{\partial P}{\partial m} \right|_{\rho, n, e} = -\Gamma u \quad (52)$$

$$\left. \frac{\partial P}{\partial n} \right|_{\rho, m, e} = -\Gamma v \quad (53)$$

$$\left. \frac{\partial P}{\partial e} \right|_{\rho, m, n} = \Gamma \quad (54)$$

JACOBIAN AND PSEUDOVISCOSITY MATRICES

Let Eqs. (1), (2), and (3), with viscous stress $\tau^{ij} = 0$ and source $q^j = 0$, be written as

$$\frac{\partial U}{\partial t} + F^j_{,j} = 0 \quad (55)$$

$$U = \begin{bmatrix} \rho \\ [m^i] \\ e \end{bmatrix} \quad F^j = \begin{bmatrix} m^j \\ \left[\frac{m^i m^j}{\rho} + g^{ij} P \right] \\ \frac{m^j}{\rho} (e + P) \end{bmatrix} \quad (56)$$

The Jacobian matrix, defined as $J^j = \partial F^j / \partial U$, is

$$J^j = \begin{bmatrix} 0 & [\delta_k^j] & 0 \\ \left[-\frac{m^i m^j}{\rho^2} + g^{ij} \frac{\partial P}{\partial \rho} \right] & \left[\frac{m^i}{\rho} \delta_k^j + \frac{m^j}{\rho} \delta_k^i + g^{ij} \frac{\partial P}{\partial m^k} \right] & \left[g^{ij} \frac{\partial P}{\partial e} \right] \\ -\frac{m^j}{\rho^2} (e + P) + \frac{m^j}{\rho} \frac{\partial P}{\partial \rho} & \left[\frac{e + P}{\rho} \delta_k^j + \frac{m^j}{\rho} \frac{\partial P}{\partial m^k} \right] & \frac{m^j}{\rho} \left(1 + \frac{\partial P}{\partial e} \right) \end{bmatrix} \quad (57)$$

where, corresponding to Eqs. (51) through (54),

$$\frac{\partial P}{\partial \rho} = c^2 - \frac{\Gamma}{\rho} (P + e) + \frac{\Gamma}{2} g_{ij} m^i m^j = c^2 - \Gamma \frac{e + P}{\rho} + \Gamma g_{ij} u^i u^j \quad (58)$$

$$\frac{\partial P}{\partial m^k} = -\frac{\Gamma}{\rho} g_{ki} m^i = -\Gamma g_{ki} u^i \quad (59)$$

$$\frac{\partial P}{\partial e} = \Gamma \quad (60)$$

The matrix for J^j given in Eq. (57) is compound, with submatrices (except in the corners) for which i and k label rows and columns, respectively.

For a three-dimensional orthogonal coordinate system, the matrices U ,

F , and J may be expanded as follows:

$$U = \begin{bmatrix} \rho \\ \rho u^1 \\ \rho u^2 \\ \rho u^3 \\ e \end{bmatrix} \quad F^j = \begin{bmatrix} \rho u^j \\ \rho u^1 u^j + g^{1j} P \\ \rho u^2 u^j + g^{2j} P \\ \rho u^3 u^j + g^{3j} P \\ u^j (e + P) \end{bmatrix} \quad (61)$$

$$J^j = \begin{bmatrix} 0 & s_1^j & s_2^j & s_3^j & 0 \\ -u^1 u^j + g^{1j} \frac{\partial P}{\partial \rho} & u^j + (1 - \Gamma) s_1^j u^1 & u^1 s_2^j - \Gamma g^{1j} g_{22} u^2 & u^1 s_3^j - \Gamma g^{1j} g_{33} u^3 & g^{1j} \Gamma \\ -u^2 u^j + g^{2j} \frac{\partial P}{\partial \rho} & u^2 s_1^j - \Gamma g^{2j} g_{11} u^1 & u^j + (1 - \Gamma) s_2^j u^2 & u^2 s_3^j - \Gamma g^{2j} g_{33} u^3 & g^{2j} \Gamma \\ -u^3 u^j + g^{3j} \frac{\partial P}{\partial \rho} & u^3 s_1^j - \Gamma g^{3j} g_{11} u^1 & u^3 s_2^j - \Gamma g^{3j} g_{22} u^2 & u^j + (1 - \Gamma) s_3^j u^3 & g^{3j} \Gamma \\ -\frac{1}{\rho} u^j (e + P) + u^j \frac{\partial P}{\partial \rho} & \frac{e + P}{\rho} s_1^j - u^j \Gamma g_{11} u^1 & \frac{e + P}{\rho} s_2^j - u^j \Gamma g_{22} u^2 & \frac{e + P}{\rho} s_3^j - u^j \Gamma g_{33} u^3 & u^j (1 + \Gamma) \end{bmatrix} \quad (62)$$

The eigenvalues of J^j are u^j , u^j , u^j , $u^j + c^j$, and $u^j - c^j$, where $c^j = c\sqrt{g^{jj}}$. Then

$$\frac{\partial U}{\partial t} + J^j U_{,j} = 0 \quad (63)$$

represents a further alternative formulation for source-free, nonviscous systems. The linearized hydrodynamic equations result when J^j is taken to be constant.

The Lax-Wendroff pseudoviscosity matrix, developed for the one-dimensional case (Ref. 1), is a quadratic function of J constructed so that the eigencolumns of J corresponding to the eigenvalues u , $u + c$, and $u - c$ are also eigencolumns of Q corresponding to the eigenvalues q_1 , q_2 , and q_3 , respectively:

$$\begin{aligned} Q = & q_1 \left[\frac{J - (u + c)I}{-c} \right] \left[\frac{J - (u - c)I}{c} \right] \\ & + q_2 \left[\frac{J - uI}{c} \right] \left[\frac{J - (u - c)I}{2c} \right] \\ & + q_3 \left[\frac{J - uI}{-c} \right] \left[\frac{J - (u + c)I}{-2c} \right] \end{aligned} \quad (64)$$

where I is the unit matrix. Directional indices are understood for Q , q , J , u , and c , as in Eq. (64), but need not be shown in a one-dimensional formula. The eigenvalues q_1 , q_2 , and q_3 are chosen as positive-definite expressions

$$\begin{aligned}
 q_1 &= \kappa \left| \xi \frac{\partial u}{\partial x} \right| \\
 q_2 &= \kappa \left| \xi \frac{\partial}{\partial x} (u + c) \right| \\
 q_3 &= \kappa \left| \xi \frac{\partial}{\partial x} (u - c) \right|
 \end{aligned} \tag{65}$$

where ξ is a coefficient with the dimensions of x (later to be identified with a basic mesh interval) and $\kappa \sim 1$ is an arbitrary numerical coefficient.

Equation (55) is then modified to read

$$\frac{\partial U}{\partial t} + F_{,x} = \Phi_{,x} = \frac{\partial}{\partial x} \Phi \tag{66}$$

where

$$\Phi = 1/2 Q \xi U_{,x} \tag{67}$$

is to be treated like a scalar under covariant differentiation. Equation (64) may also be written

$$Q = a_0 I + a_1 J + a_2 J^2 \tag{68}$$

where

$$\begin{aligned}
 a_0 &= [-q_1(u^2 - c^2) + 1/2 q_2 u(u - c) + 1/2 q_3 u(u + c)]/c^2 \\
 a_1 &= [2q_1 u - 1/2 q_2 (2u - c) - 1/2 q_3 (2u + c)]/c^2 \\
 a_2 &= [-q_1 + 1/2 q_2 + 1/2 q_3]/c^2
 \end{aligned} \tag{69}$$

In order to avoid calculation of the Jacobian in second-order difference approximations to Eqs. (1), (2), and (3), it is customary to employ two-step time differencing (Ref. 2). In this case, a procedure attributed to Lapidus may be used to eliminate the Jacobian from the pseudoviscosity terms as well. The following derivation shows one way in which this result can be reached.

Equations (66), (67), and (68) yield

$$\frac{\partial U}{\partial t} + F_{,x} = 1/2 \partial_x (a_0 I + a_1 J + a_2 J^2) \xi U_{,x} \tag{70}$$

or

$$\frac{\partial U}{\partial t} + \bar{F}_{,x} = 1/2 \partial_x (a_0 \xi U_{,x}) \quad (71)$$

where

$$\bar{F} = F - 1/2 (a_1 J \xi U_{,x} + a_2 J \xi F_{,x}) \quad (72)$$

A time-difference equation approximating Eq. (71) may be written

$$U^{n+1} = U^n - \Delta t \bar{F}_{,x}^{n+1/2} + 1/2 \Delta t \frac{\partial}{\partial x} \xi (a_0 U_{,x})^{n+1/2} \quad (73)$$

The pseudoviscosity terms, being of order $\xi^2 \Delta t$, need not be centered.

Thus, from Eq. (72),

$$\bar{F}^{n+1/2} - F^{n+1/2} = -1/2 J^n \xi (a_1 U_{,x} + a_2 F_{,x})^n \quad (74)$$

Defining a new quantity $U^{n+1/2}$ such that

$$\bar{F}^{n+1/2} - F^n = J^n (U^{n+1/2} - U^n) \quad (75)$$

and then subtracting Eq. (74) from Eq. (75) yields

$$F^{n+1/2} - F^n = J^n [U^{n+1/2} - U^n + 1/2 \xi (a_1^n U_{,x}^n + a_2^n F_{,x}^n)] \quad (76)$$

Here, $F^{n+1/2} = F(U_0^{n+1/2})$, where U_0 is the solution of Eq. (55), i. e., the undamped case. Therefore,

$$F^{n+1/2} - F^n = \frac{\Delta t}{2} J^n \frac{\partial U_0^n}{\partial t} = - \frac{\Delta t}{2} J^n F_{,x}^n \quad (77)$$

From Eqs. (76) and (77), the difference equation for $U^{n+1/2}$ is obtained:

$$U^{n+1/2} = U^n - 1/2 \Delta t F_{,x}^n - 1/2 \xi (a_1 U_{,x}^n + a_2 F_{,x}^n) \quad (78)$$

Equations (78) and (73), with the relation

$$\bar{F}^{n+1/2} = F(U^{n+1/2}) \quad (79)$$

form an explicit time-difference equation system with Lax-Wendroff pseudoviscosity, in the two-step form, for the one-dimensional case.

Unfortunately, since the J and Q matrices corresponding to different directions fail to commute, the generalization of the Lax-Wendroff pseudoviscosity to two or more space dimensions is not straightforward. By superposition of the above one-dimensional forms (Ref. 3), the system can be written (with sum convention suppressed) as

$$U^{n+\frac{1}{2}} = \left[U - \frac{1}{2}\Delta t \sum_j \partial_j F^j - \frac{1}{2} \sum_j \xi^j a_1^j U_{,j} - \frac{1}{2} \sum_j \xi^j a_2^j F^j_{,j} \right]^n \quad (80)$$

$$U^{n+1} = U^n - \Delta t \sum_j \partial_j (F^j - \frac{1}{2} a_0^j \xi^j U_{,j})^{n+\frac{1}{2}} \quad (81)$$

In evaluating the covariant derivatives for the axially symmetric case, it should be noted that $a_0^{(2)} = a_1^{(2)} = a_2^{(2)} = 0$. Therefore, in this case the derivatives within the pseudoviscosity terms, denoted by commas, are simple derivatives containing no curvature terms. The expression $a_0^j \xi^j U$, like Φ^j , is not to be treated as a vector or tensor field when being differentiated. The a_p^j are defined by Eqs. (65) and (69), with a directional superscript j (not subject to the sum convention) applied to u , c , x , q , ξ , and perhaps κ .

A simplified formulation, attributed (Ref. 2) to Lapidus, is obtained by annulling a_1^j and a_2^j and setting a_0^j equal to q_1^j :

$$Q^j = a_0^j I = \kappa \left| \xi^j \frac{\partial}{\partial x^j} u^j \right| I \quad (\text{no sum convention}) \quad (82)$$

The viscous stress τ^{ij} can be accounted for in the definition of F as in Eqs. (40) and (48), or by introducing on the right side of Eqs. (55) and (63) a dissipative term $V^j_{,j}$, where

$$V^j = \begin{bmatrix} 0 \\ [\tau^{ij}] \\ g_{\alpha\beta} \frac{m^\beta}{\rho} \tau^{\alpha j} \end{bmatrix} \quad (83)$$

The Jacobian $J_\tau^j = \partial V^j / \partial U$ is

$$J_\tau^j = \begin{bmatrix} 0 & [0] & 0 \\ [\partial \tau^{ij} / \partial \rho] & [\partial \tau^{ij} / \partial m^k] & [\partial \tau^{ij} / \partial e] \\ g_{\alpha\beta} \frac{m^\beta}{\rho} \left(\frac{\partial \tau^{\alpha j}}{\partial \rho} - \frac{\tau^{\alpha j}}{\rho} \right) & \left[g_{\alpha\beta} \frac{m^\beta}{\rho} \frac{\partial \tau^{\alpha j}}{\partial m^k} + g_{\alpha k} \frac{\tau^{\alpha j}}{\rho} \right] & g_{\alpha\beta} \frac{m^\beta}{\rho} \partial \tau^{ij} / \partial e \end{bmatrix} \quad (84)$$

It is apparent that J_τ^j has at least one zero eigenvalue, due to the absence of τ^{ij} in Eq. (1). A pseudoviscosity of this form, based upon a suitable redefinition of τ^{ij} , is thus relatively ineffective for stabilizing Eulerian difference equations against nonlinear instabilities such as those investigated by Burstein (Ref. 3). However, the Richtmyer-von Neumann pseudoviscosity, which is of this form, has long been used for Lagrangian equations, and variants have been successfully used for Eulerian cases, such as the following:

$$\tau^{11} = 2\nu^{(1)} \frac{\partial u}{\partial r}, \quad r^2 \tau^{22} = 2\nu^{(1)} \frac{u}{r}, \quad \tau^{33} = 2\nu^{(3)} \frac{\partial v}{\partial y}$$

where

$$\nu^{(1)} = \kappa \rho \xi^{(1)} \text{Max} \left(0, -\xi^{(1)} \frac{\partial u}{\partial r} \right), \quad \nu^{(3)} = \kappa \rho \xi^{(3)} \text{Max} \left(0, -\xi^{(3)} \frac{\partial v}{\partial y} \right)$$

TWO-STEP LAX-WENDROFF DIFFERENCE EQUATIONS WITH AXIAL SYMMETRY

In formulating the space-time difference equations, it is convenient to avoid half-integer indices by ascribing even integer indices to the basic mesh quantities and odd indices to the midpoints. The basic time difference

Δt introduced in the preceding section will therefore be replaced by $2\Delta t$ in this and subsequent sections, and the space differences ξ^1 and ξ^3 by

$$2\delta r_i = r_{i+1} - r_{i-1} \quad (85)$$

$$2\delta y_j = y_{j+1} - y_{j-1} \quad (86)$$

respectively. The following equations are written for even values of the space-time parity index $i+j+n$; however, it will be noted that both curvature and pseudoviscosity terms depend upon odd-parity values of U , such as $U_{i+1,j}^n$. The definitions of F , G , and H given in Eq. (48) permit elimination of tensor notation. For a uniform rectangular spatial mesh,

$$U_{i+1,j}^n = \frac{1}{4} (U_{i+2,j}^n + U_{ij}^n + U_{i+1,j+1}^n + U_{i+1,j-1}^n) \quad (87)$$

$$\begin{aligned} U_{i+1,j}^{n+1} = & U_{i+1,j}^n - 1/2 \left[\frac{\Delta t}{\delta r_{i+1}} + (a_2^{(1)})_{i+1,j}^n \right] (F_{i+2,j}^n - F_{ij}^n) - 1/2 \left[\frac{\Delta t}{\delta y_j} + (a_2^{(3)})_{i+1,j}^n \right] \\ & (G_{i+1,j+1}^n - G_{i+1,j-1}^n) - 1/2 \Delta t H_{i+1,j}^n - 1/2 (a_1^{(1)})_{i+1,j}^n (U_{i+2,j}^n - U_{ij}^n) \\ & - 1/2 (a_1^{(3)})_{i+1,j}^n (U_{i+1,j+1}^n - U_{i+1,j-1}^n) \end{aligned} \quad (88)$$

$$U_{ij}^{n+1} = 1/4 (U_{i+1,j}^{n+1} + U_{i-1,j}^{n+1} + U_{i,j+1}^{n+1} + U_{i,j-1}^{n+1}) \quad (89)$$

$$\begin{aligned} U_{ij}^{n+2} = & U_{ij}^n - \frac{\Delta t}{\delta r_i} (F_{i+1,j}^{n+1} - F_{i-1,j}^{n+1}) - \frac{\Delta t}{\delta y_i} (G_{i,j+1}^{n+1} - G_{i,j-1}^{n+1}) - \Delta t H_{ij}^{n+1} \\ & + 1/2 \frac{\Delta t}{\delta r_i} [(a_0^{(1)})_{i+1,j}^{n+1} (U_{i+2,j}^{n+1} - U_{ij}^{n+1}) - (a_0^{(1)})_{i-1,j}^{n+1} (U_{ij}^{n+1} - U_{i-2,j}^{n+1})] \\ & + 1/2 \frac{\Delta t}{\delta y_i} [(a_0^{(3)})_{i,j+1}^{n+1} (U_{i,j+2}^{n+1} - U_{ij}^{n+1}) - (a_0^{(3)})_{i,j-1}^{n+1} (U_{ij}^{n+1} - U_{i,j-2}^{n+1})] \end{aligned} \quad (90)$$

Equations (87) and (89) may be used to eliminate odd-parity quantities from Eqs. (88) and (90), or, alternatively, the four equations may be solved in sequence. In the latter case, both even- and odd-parity values of U must

be simultaneously available, but the amount of computing required is reduced substantially. A third possibility is to modify the equations so that differences of odd-parity quantities, for example, are replaced by differences of even-parity quantities. This last approach is more complex in some ways and will therefore not be followed in this report.

A related problem is generalization to a nonuniform rectangular mesh.

Let

$$\begin{aligned}\Delta r_i &= r_{i+1} - r_i & \nabla r_i &= r_i - r_{i-1} = \Delta r_{i-1} \\ \Delta y_j &= y_{j+1} - y_j & \nabla y_j &= y_j - y_{j-1} = \Delta y_{j-1}\end{aligned}\tag{91}$$

Then Eqs. (87) and (89) become, respectively,

$$\begin{aligned}U_{i+1,j}^n &= \left(\frac{1}{\Delta r_{i+1}} + \frac{1}{\nabla r_{i+1}} + \frac{1}{\Delta y_j} + \frac{1}{\nabla y_j} \right)^{-1} \left(\frac{1}{\Delta r_{i+1}} U_{i+2,j}^n + \frac{1}{\nabla r_{i+1}} U_{ij}^n \right. \\ &\quad \left. + \frac{1}{\Delta y_j} U_{i+1,j+1}^n + \frac{1}{\nabla y_j} U_{i+1,j-1}^n \right)\end{aligned}\tag{92}$$

and

$$\begin{aligned}U_{ij}^{n+1} &= \left(\frac{1}{\Delta r_i} + \frac{1}{\nabla r_i} + \frac{1}{\Delta y_j} + \frac{1}{\nabla y_j} \right)^{-1} \left(\frac{1}{\Delta r_i} U_{i+1,j}^{n+1} + \frac{1}{\nabla r_i} U_{i-1,j}^{n+1} \right. \\ &\quad \left. + \frac{1}{\Delta y_j} U_{i,j+1}^{n+1} + \frac{1}{\nabla y_j} U_{i,j-1}^{n+1} \right)\end{aligned}\tag{93}$$

These equations, like Eqs. (87) and (89), are accurate to first order in the mesh interval.

In Eq. (88), the quantities H , $a_1^{(1)}$, $a_1^{(3)}$, $a_2^{(1)}$, and $a_2^{(3)}$ appear with odd parity. H may be most easily evaluated from the odd-parity U given by Eq. (92). From Eqs. (65) and (69), the a 's consist of differences of even-parity values of u , v , and c , with odd-parity coefficients involving the same quantities. These coefficients may be evaluated as one-directional weighted means of the adjacent even-parity quantities:

$$(a_1^{(1)})_{ij}^{n+1} = \{ [2q_1^{(1)} \bar{u} - 1/2q_2^{(1)} (2\bar{u} - \bar{c}) - 1/2q_3^{(1)} (2\bar{u} + \bar{c})] / \bar{c}^2 \}_{ij}^{n+1} \quad (94)$$

etc., where the bar denotes the appropriate one-directional mean:

$$\bar{u}_{ij}^{n+1} = \left(\frac{1}{\Delta r_i} + \frac{1}{\nabla r_i} \right)^{-1} \left(\frac{1}{\Delta r_i} u_{i+1,j}^{n+1} + \frac{1}{\nabla r_i} u_{i-1,j}^{n+1} \right) \quad (95)$$

etc. In Eq. (90), the odd-parity H_{ij}^{n+1} may be computed from the U_{ij}^{n+1} given by Eq. (93), differences of which also appear as coefficients of even-parity a_0 . These a_0 depend upon even-parity q 's which may be defined by differencing \bar{u} , \bar{v} , and \bar{c} as given by formulae like Eq. (95). For the viscous stress components, contributions to F and G are always of even parity and contributions to H of odd parity. The differencing, in this case, should always be based upon the U values with the opposite parity and the same time index.

BOUNDARY CONDITIONS FOR EXPLICIT DIFFERENCE EQUATIONS

On reflective boundaries (including always the symmetry axis), the normal component of velocity vanishes, while the tangential component of velocity and the "scalars" $a_0^{(j)}$, $a_1^{(i)}$, $a_2^{(j)}$, ρ , e , P , etc., all have vanishing normal derivatives. On the axis, the quantity u/r is equal to $\partial u / \partial r$. Therefore, the first, third, and fourth components of F and the second components of U , G , and H vanish on the axis.

Equations (88), (90), (92), and (93) explicitly refer to points which may lie outside the active mesh. Therefore, a single row of mesh points may be provided outside the active region in order to avoid rewriting these equations for boundary points. Values of quantities assigned to these external points are identical to those one mesh spacing within, except that for reflective boundaries the normal component of velocity is reversed in sign. Values of H are not required at the external points. Other quantities which vanish on the reflective boundary are antisymmetric with respect to the boundary, while the remaining quantities (whose normal derivatives

vanish on the boundary) are symmetric. Differentiation (ordinary or covariant) of a quantity in the normal direction reverses its symmetry character at a reflective boundary.

Boundaries other than the symmetry axis may be defined in various ways; a free boundary is one on which all quantities at external points vanish, whereas a symmetric transmissive boundary is one on which all quantities have vanishing normal derivatives, and an extrapolative boundary is one on which the normal second derivatives of all quantities vanish.

STABILITY OF LAX-WENDROFF DIFFERENCE EQUATIONS

The stability of the linearized Lax-Wendroff equations (Eq. (63) with constant J^j) is analyzed in reference 2, where a sufficient condition for stability is found to be

$$(|u| + c) \frac{\Delta t}{\Delta x} < \frac{1}{\sqrt{2}} \quad (96)$$

for the case in which $\delta y = \delta r = \Delta x$. The analysis is more complex for the case of a nonuniform rectangular mesh, but the pair of conditions

$$(|u| + c) \frac{\Delta t}{\delta r} < \frac{1}{\sqrt{2}} \quad (97)$$

$$(|v| + c) \frac{\Delta t}{\delta y} < \frac{1}{\sqrt{2}} \quad (98)$$

can be taken as essentially equivalent. The effect of the viscosity terms is very complex, but Burstein has shown (Ref. 3) that the a_2 contributions probably increase stability slightly, whereas it is shown in Ref. 2 that the effect of a constant a_0 is slightly destabilizing. The simple viscosity given by Eq. (82) is always stabilizing, however, and the same is probably true of the complete formulation in the nonlinear case.

The coefficient κ which appears in the pseudoviscous terms can be modified if the calculation runs into difficulties such as negative internal

energy. In order for this to be possible, the contribution of the pseudo-viscous terms at each point must be temporarily saved while the internal energy is calculated. Since each component of U is linear in κ , it is possible to write

$$\begin{aligned} \rho &= \rho_0 + \kappa \frac{\partial \rho}{\partial \kappa} & n &= n_0 + \kappa \frac{\partial n}{\partial \kappa} \\ m &= m_0 + \kappa \frac{\partial m}{\partial \kappa} & e &= e_0 + \kappa \frac{\partial e}{\partial \kappa} \end{aligned} \quad (99)$$

where the subscript zero denotes the value without pseudoviscosity. Then

$$\begin{aligned} 2\rho e - m^2 - n^2 &= 2\rho_0 e_0 - m_0^2 - n_0^2 + 2\kappa \left(\rho_0 \frac{\partial e}{\partial \kappa} + e_0 \frac{\partial \rho}{\partial \kappa} - m_0 \frac{\partial m}{\partial \kappa} - n_0 \frac{\partial n}{\partial \kappa} \right) \\ &\quad + \kappa^2 \left[2 \frac{\partial \rho}{\partial \kappa} \frac{\partial e}{\partial \kappa} - \left(\frac{\partial m}{\partial \kappa} \right)^2 - \left(\frac{\partial n}{\partial \kappa} \right)^2 \right] \\ &= C + 2\kappa B + \kappa^2 A \end{aligned} \quad (100)$$

The value of κ required to bring the internal energy to zero is thus

$$\kappa = \frac{-B \pm \sqrt{B^2 - AC}}{A} \quad (101)$$

where, assuming that the pseudoviscosity has a genuine dissipative effect, the sign can be chosen so that κ is positive as well as real. A higher positive value of κ will result in a positive internal energy; a difference approximation to Eq. (24) could be used to estimate an appropriate value.

SUMMARY AND CONCLUSION

The basic equations of the method are Eqs. (88) and (90). The former is solved on odd-numbered time steps, the latter on even-numbered time steps. Values of U , i.e., the four conserved densities ρ , $\rho\mu$, ρv , and e , are computed for just half of the points in the spatial mesh in each time step, namely the points with indices i, j in time step n such that the "parity index" $i + j + n$ is an even integer. Values of U for odd-parity points are

treated differently; where needed, they are computed from the even-parity values by interpolation, through the use of Eqs. (92) and (93).

The quantities F , G , and H appearing in Eqs. (88) and (90) are defined in Eq. (48). They depend upon the components of U ; upon pressure P , which depends upon components of U (cf. Eqs. (6) and (4)); upon viscous stress τ , which depends upon differences of components of U (cf. Eq. (5)) and upon viscosities μ and μ_B , which may be functions of state and kinematic variables; and finally upon heat flux \vec{q} , which must be independently prescribed but may also depend upon state and kinematic variables.

The pseudoviscosity quantities a_0 , a_1 , and a_2 appearing in Eqs. (88) and (90) depend upon velocity components and their differences; upon the sound speed; and upon a parameter κ , which may be constant or under program control. The definitions are given in Eqs. (65) and (69). A simpler formulation is presented in Eq. (82). The purpose of these terms is to improve the stability characteristics of the difference equations; they have no physical significance.

The entire method outlined in this section is the subject of current research at many laboratories and is far from being a panacea for all problems which arise in the calculation of two-dimensional flow. However, development of this approach, as an alternative to the method employed in the OIL family of codes, is expected to improve the treatment of vapor-phase flow which can undergo very large density changes.

SECTION II

RADIATIVE TRANSFER IN HECTICINTRODUCTION

The immediate objective of the current work on radiative transfer in two-dimensional geometries has been to develop numerical methods that can be applied in HECTIC. The first step is to find a way to solve the equation

$$\vec{\Omega} \cdot \nabla I(\vec{r}, \vec{\Omega}) + \sigma(\vec{r}) I(\vec{r}, \vec{\Omega}) = S(\vec{r}, \vec{\Omega}) \quad (102)$$

where I , S , and σ are, respectively, intensity, source, and absorption coefficients as functions of position r and direction Ω . Thus, at this step, polarization, frequency, and retardation are ignored. The source S is assumed known, so such problems as scattering are also deferred.

Descriptions of two codes that were partially developed for solving Eq. (102) as subroutines of HECTIC are given in reference 4. The first, TDRAD, uses the nonequilibrium diffusion method and the second, LONG2, uses the method of long characteristics. The current work, in general terms, has been devoted to:

1. Extending and refining the long-characteristics code by allowing a more flexible scheme for distributing rays and by providing a means of treating regions in which the diffusion approximation is valid.
2. Surveying methods for solving the elliptic difference equations that arise in two-dimensional versions of the nonequilibrium diffusion method.
3. Developing a new method, called the "view-factor" method, which avoids the most serious difficulties encountered in the other two.

The present report contains detailed descriptions of LONG2, the code that uses the extended method of characteristics, and SHORT2, the code that implements the view-factor method. Work performed on TDRAD is also reported. Specifically, descriptions of three algorithms for solving the nonequilibrium diffusion equations and the results of the tests to which they were put are presented.

An attempt has been made to keep the report reasonably self-contained, even at the expense of repeating much information that can be found elsewhere.

DEFINITIONS AND NOTATION

It is appropriate to establish at the outset various notations and conventions used in the discussion of the various routines for calculating radiative transfer in HECTIC. The physical systems these routines are designed to handle must be axially symmetric. Let the axis of symmetry be the z-axis of an x-y-z cartesian coordinate system. Let $r = (x^2 + y^2)^{1/2}$ denote the distance from the axis of the point with coordinates (x, y, z). The z-axis is called "vertical" and grad z is directed upward. Planes $z = \text{constant}$ are horizontal. As might be expected, grad r points outward and its negative, inward. The region of interest lies inside a cylinder $r = R$ and between two planes $z = z_{\min}$ and $z = z_{\max}$. This region is divided into a number of subregions, called "zones", by a set of cylinders $r = r_i$, where $1 \leq i \leq i_{\max}$, and a set of planes $z = z_j$, where $0 \leq j \leq j_{\max}$. As usual, $r_{i-1} < r_i$ and $z_{j-1} < z_j$.

Thus, a zone is a toroid of rectangular cross section. The zones fill the region so $r_{i_{\max}} = R$, $z_0 = z_{\min}$, and $z_{j_{\max}} = z_{\max}$. In the current FORTRAN code for HECTIC, the zones are numbered from 2 to $IMAX * JMAX + 1$, where $IMAX = i_{\max}$ and $JMAX = j_{\max}$. But in this part of the report, a double subscript is used and $Z_{i,j}$ denotes the zone bounded by

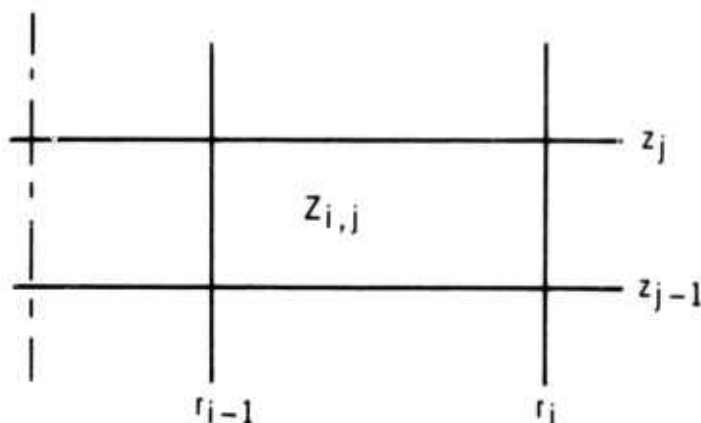


Figure 1. A Typical Zone in HECTIC

the cylinders $r = r_{i-1}$ and $r = r_i$ and the planes $z = z_{j-1}$ and $z = z_j$. (See Fig. 1.) The surface of zone $Z_{i,j}$ has four connected pieces, called "boundaries":

$BL_{i,j}$, lying on the cylinder $r = r_{i-1}$

$BR_{i,j}$, lying on the cylinder $r = r_i$

$BB_{i,j}$, lying on the plane $z = z_{j-1}$

$BA_{i,j}$, lying on the plane $z = z_j$

Let $\hat{z} = \text{grad } z$, a unit vector parallel to the axis in the system, let $\hat{r} = \text{grad } r$, a unit vector orthogonal to \hat{z} and directed outward along a line through the axis, and let $\hat{z} \times \hat{r}$ be the third unit vector of the local basis.

If \vec{Q} is the vector connecting the origin of coordinates to the point in question, then

$$\vec{P} = \vec{Q} + s\vec{\Omega} \quad (103)$$

is the parametric equation of a ray through the point. The vector $\vec{\Omega}$ is the direction of the ray. Directions are required to be unit vectors, and consequently the parameter s in Eq. (103) is the distance from Q to $P(s)$.

With respect to the local basis about the point $P(s)$,

$$\vec{\Omega} = \mu \hat{z} - (1 - \mu^2)^{1/2} (\hat{r} \cos \phi + \hat{r} \times \hat{z} \sin \phi) \quad (104)$$

where

$$\begin{aligned} \mu &= \vec{\Omega} \cdot \hat{z} \\ \phi &= \arctan (\vec{\Omega} \cdot \hat{r} \times \hat{z} / \sqrt{\vec{\Omega} \cdot \vec{\Omega}}) \end{aligned} \quad (105)$$

μ is called the "axial projection" of the ray and ϕ is called its "azimuth" at the point P. Since the direction of the local basis vector \hat{r} depends on the coordinates of P, the magnitude of the azimuth ϕ depends on s , as may be seen from figure 2. When s is negative and large in magnitude, the horizontal component of $\vec{\Omega}$ points almost directly inward and ϕ has a small positive value. As s increases, so does ϕ . When $\phi < \pi/2$, the point is approaching the axis, and when $\phi > \pi/2$, it is receding from the axis. On the other hand, the axial projection μ of the ray depends only on its direction $\vec{\Omega}$ and does not vary with s .

The range $-\pi \leq \phi \leq 0$ is occupied by the azimuths of rays that pass the axis moving clockwise as viewed from above. In view of the assumed axial symmetry of the physical system, the variation of radiant intensity along such a ray is the same as along its reflection in the plane $y = 0$ (or any other plane containing the axis for that matter). Consequently, such rays need not be considered.

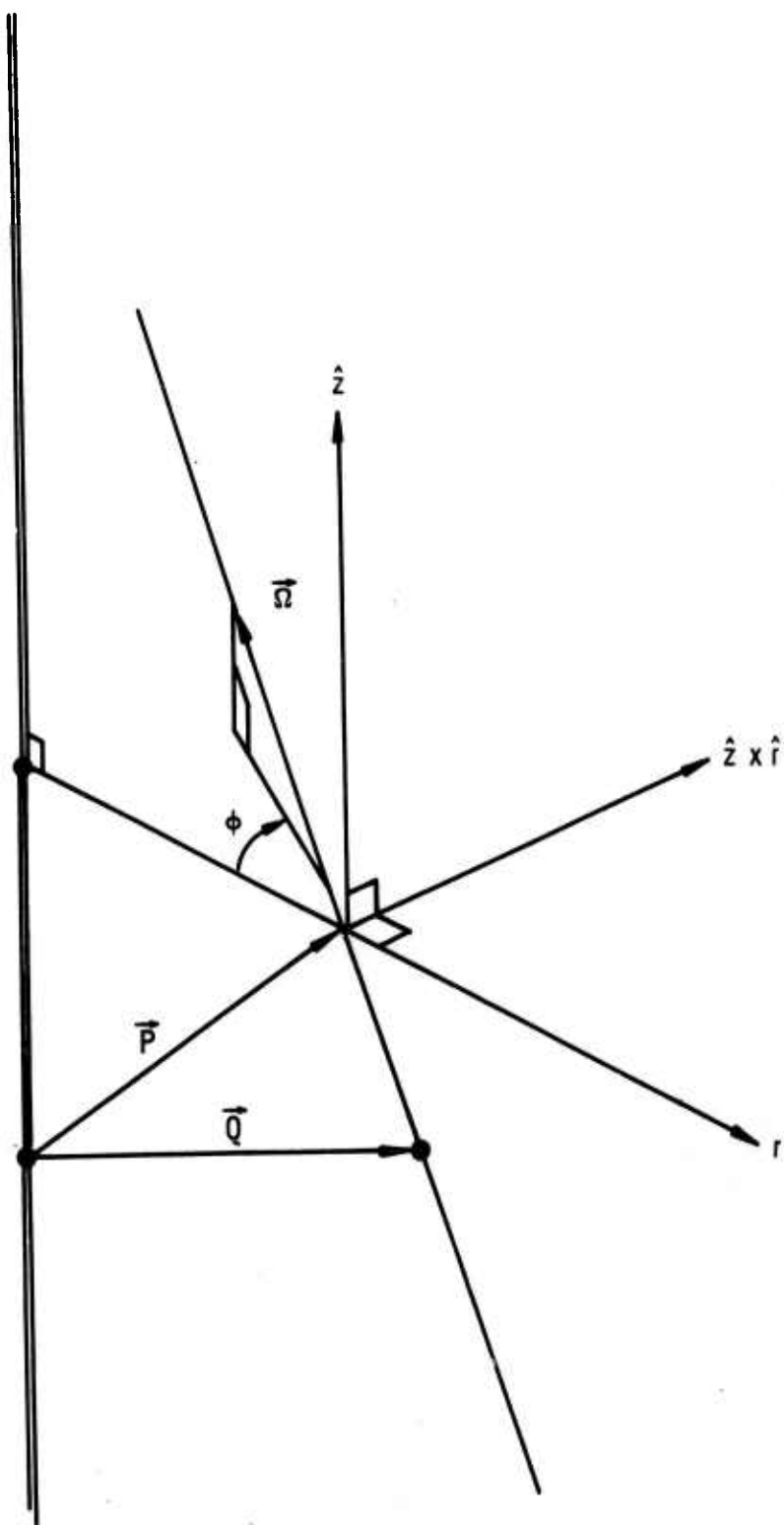


Figure 2. The Local Basis Vectors

SECTION III

THE LONG-CHARACTERISTIC METHOD

The idea behind the long-characteristic method is to calculate radiant intensity by means of the formula

$$I(\vec{b} + s\vec{\Omega}, \vec{\Omega}) = I(\vec{b}, \vec{\Omega}) e^{-\int_0^s \sigma(\vec{b} + u\vec{\Omega}) du} + \int_0^s S(\vec{b} + u\vec{\Omega}, \vec{\Omega}) e^{-\int_u^s \sigma(\vec{b} + v\vec{\Omega}) dv} du \quad (106)$$

where $I(\vec{r}, \vec{\Omega})$ is the intensity of radiation at \vec{r} in direction $\vec{\Omega}$, $\sigma(\vec{r})$ is the absorption coefficient at \vec{r} , and $S(\vec{r}, \vec{\Omega})$ is the strength of the source at \vec{r} in direction $\vec{\Omega}$. The specification of the intensity $I(\vec{b}, \vec{\Omega})$ at a point \vec{b} on the boundary of the system in a direction $\vec{\Omega}$ that points into the system is a boundary condition. If, as is done in LONG2, the absorption coefficients and source strengths are taken to be constant within the hydro zones of HECTIC, the integrals of Eq. (106) are elementary. The problem that then remains is to determine the rate of energy deposition from a given distribution of intensities. The solution adopted in the subroutines of LONG2 is described below.

A SIMPLIFIED VERSION

Before the details of the long-characteristic method as implemented in LONG2 are discussed, a simplified version will be presented. The first step is to select a set of rays that are distributed through the physical

system more or less uniformly both in space and direction. In the simple version, this is done as follows. Let N be a large integer and let d be a length that is small compared with the dimensions of the zones in the HECTIC hydro mesh. Let

$$\mu_n = \frac{N-n+\frac{1}{2}}{N}, \quad \eta_n = (1 - \mu_n^2)^{\frac{1}{2}} \quad (107)$$

for $1 \leq n \leq 2N$. Then to each triplet of integers, ℓ, m, n for which $1 \leq n \leq N$, a ray $R_{\ell, m, n}$ is associated whose parametric equations are

$$\begin{aligned} x &= \eta_n s \\ y &= (\ell - \frac{1}{2}) d \\ z &= \frac{m d}{\eta_n} + \mu_n s \end{aligned} \quad (108)$$

the parameter s being arc length. The restriction $\ell \leq R/d$ on ℓ guarantees that the ray comes close enough to the axis to pierce the cylinder $r = R$. There is some complicated restriction on m that will ensure that the ray actually does penetrate the region of interest, but it serves no useful purpose to write it down. The set of rays corresponding to a fixed value of n is called a "grid." The rays of a grid are all parallel and intersect any plane to which they are normal in a square lattice of points with a separation of d between nearest neighbors. The set of rays corresponding to a fixed value of n and ℓ forms a comb.

To verify that the complete set of rays is uniform, their distribution over some small element of volume will be analyzed. In view of the axial symmetry of the system, the element of volume should be a figure of revolution. Let it be a toroid of rectangular cross section: $r_0 \leq r \leq r_0 + \Delta r$, $z_0 \leq z \leq z_0 + \Delta z$, where $\Delta z \ll \Delta r$, $d \ll \Delta z$. Consider those rays whose directions have axial projections lying in the range $\mu_0 \leq \mu \leq \mu_0 + \Delta \mu$ that

intersect the volume element when their azimuths lie in the range $\phi_0 \leq \phi \leq \phi_0 + \Delta\phi$, and calculate the sum of the lengths of these intersections. The length of one such intersection is

$$t = \frac{\Delta r}{(1 - \mu_0^2)^{\frac{1}{2}} \cos \phi_0} \quad (109)$$

It can be seen from figure 3 that $r \sin \phi = y$ and hence that $r \cos \phi \Delta\phi = \Delta y$. Therefore, $\Delta l = \Delta y/d = r_0 \cos \phi_0 \Delta\phi/d$ is the number of l values of rays that have their azimuths in the desired range at $r = r_0$. The vertical separation between $R_{l,m,n}$ and ray $R_{l,m+1,n}$ is d/η_n . Therefore, the number of m values of rays with fixed l and n striking the volume element is $\Delta m = \Delta z/(d/\eta_n) = \eta_n \Delta l/d$. Finally, the number of n -values with μ in the desired range is simply $\Delta n = N\Delta\mu$. Thus, the sum T of the lengths of the intersections is

$$\begin{aligned} T &= t \Delta l \Delta m \Delta n = \Delta r r_0 \Delta\phi \Delta z N \Delta\mu/d^2 \\ &= \frac{N \Delta V \Delta\Omega}{2\pi d^2} \end{aligned} \quad (110)$$

where $\Delta V = 2\pi r_0 \Delta r \Delta z$ is the magnitude of the volume element and $\Delta\Omega = \Delta\phi \Delta\mu$ is the solid angle subtended by the range of directions being considered. Since the sum of the lengths is proportional to $\Delta V \Delta\Omega$ and independent of r_0 , z_0 , μ_0 , and ϕ_0 , one may conclude that the distribution is uniform in both space and direction.

The rate of deposition of radiant energy in a HECTIC hydro zone is calculated as the difference between the inward-going flux and the outward-going flux through the boundaries of the zone. The net inward flux through a vertical boundary of radius r and altitude ranging from z_1 to z_2 is

$$F = 2 \int_0^\pi \int_{-1}^1 \int_{z_1}^{z_2} \eta \cos \phi I(r, z, \vec{\Omega}) 2\pi r dz d\mu d\phi \quad (111)$$

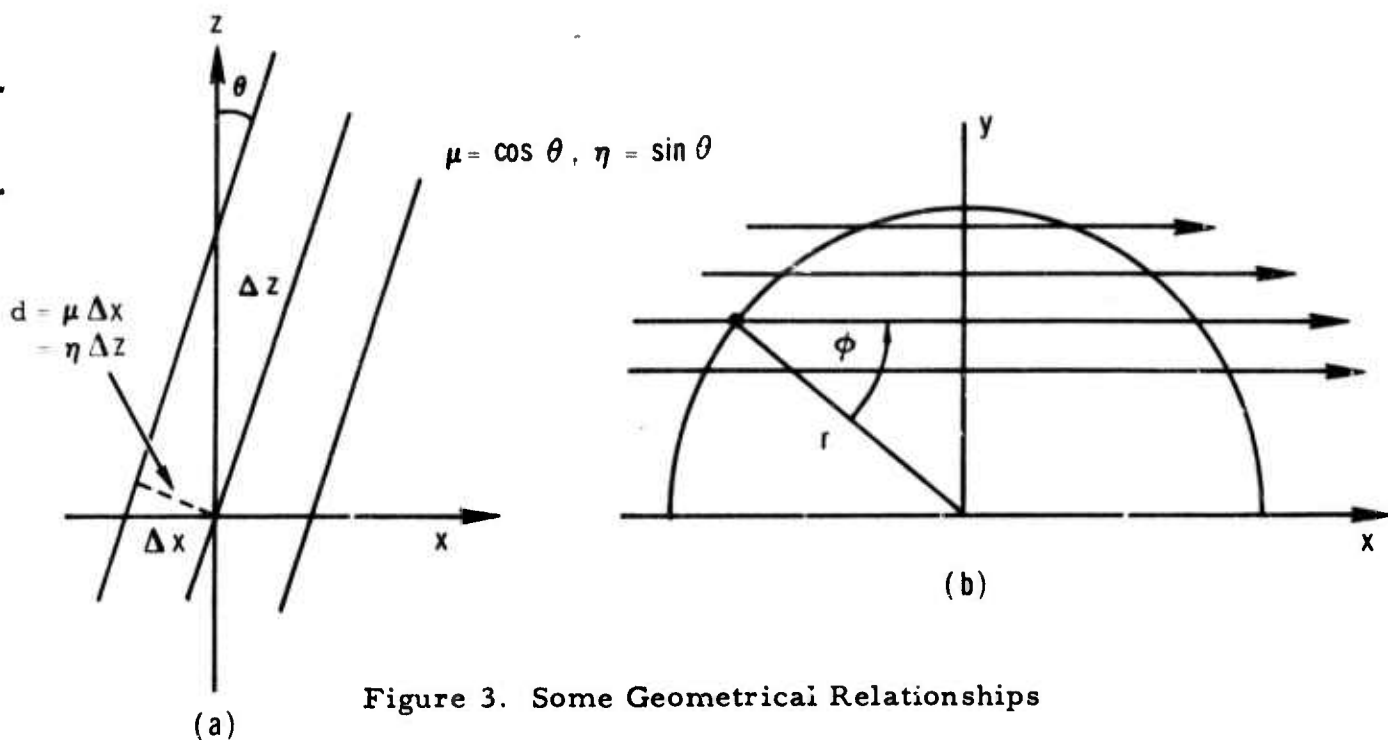


Figure 3. Some Geometrical Relationships

where $\eta = (1 - \mu^2)^{\frac{1}{2}}$. The value of the integral is estimated by the simplest method, i.e., by replacing the integral with finite sum:

$$F = 4\pi r \sum \eta \cos \phi I \Delta z \Delta \mu \Delta \phi \quad (112)$$

where the sum extends over all rays striking the surface. The vertical distance between rays is d/η , so it is appropriate to set $\eta \Delta z = d$ (see Fig. 3a). Next, referring to Eq. (108) and figure 3b, it is seen that $\sin \phi = y/r$, so $|\cos \phi| \Delta \phi = \Delta y/r = d/r$. Finally, $\Delta \mu = 1/N$. Thus, Eq. (112) becomes

$$F = \frac{4 \pi d^2}{N} \sum_{\pm 1} \quad (113)$$

where the sign to use is that of $\cos \phi$: + for inward-going and - for outward-going rays. The net upward flux through a horizontal zone boundary with inner radius r_1 , outer radius r_2 , and altitude z is

$$F = 2 \int_0^\pi \int_{-1}^1 \int_{r_1}^{r_2} \mu I(r, z, \vec{\Omega}) 2 \pi r dr d\mu d\phi \quad (114)$$

Referring again to figure 3b, it can be seen

$$x = -r \cos \phi$$

$$y = r \sin \phi$$

and hence

$$\begin{aligned} \frac{\partial(r, \phi)}{\partial(x, y)} &= \begin{vmatrix} \frac{\partial r}{\partial x} & \frac{\partial r}{\partial y} \\ \frac{\partial \phi}{\partial x} & \frac{\partial \phi}{\partial y} \end{vmatrix} \\ &= \begin{vmatrix} \frac{x}{r} & \frac{y}{r} \\ \frac{y}{r^2} & -\frac{x}{r^2} \end{vmatrix} \\ &= -\frac{x^2}{r^3} - \frac{y^2}{r^3} = -\frac{1}{r} \end{aligned}$$

Therefore, $r dr d\phi = dx dy$. But, as before, $\Delta y = d$ and $|\mu| \Delta x = d$; so making the same sort of estimate of the flux F in Eq. (114) again yields Eq. (113), where this time the sign of I is that of μ : + for upward and - for downward. The rate of deposition of radiant energy in a HECTIC hydro zone is therefore

$$\dot{E} = \frac{4 \pi d^2}{N} \sum (I_{in} - I_{out}) \quad (115)$$

the sum extending over all rays that strike the zone.

THE IMPLEMENTED VERSION

The long-characteristic method, as implemented in LONG2, is a refinement of the simple version described above. Some of the principal changes are (1) more flexibility in the selection of rays, (2) the inclusion of some correction factors, and (3) a special treatment for optically thick zones.

One of the practical difficulties of the long-characteristic method is the large number of arithmetic operations it involves. It is necessary to calculate the change of intensity along every line segment resulting from the intersection of a ray with a hydro zone. It is therefore desirable to allow the widest possible latitude in the distribution of rays so that certain regions may be sampled more thoroughly than others.

In the simple version, a grid consists of rays that are parallel and intersect any orthogonal plane in a square lattice of points. The first generalization incorporated into LONG2 is a facility for having a higher density of rays in some regions than in others. As before, let $2N$ be the number of grids and let n be the index of a grid. For each n , $1 < n \leq 2N$, there are selected two increasing sequences $(y_{l,n})_{l=1}^{L_n}$ and $(z_{m,n}^*)_{m=1}^{M_n}$. With this n and every pair l, m of integers is associated a ray $R_{l,m,n}$ whose parametric equations are

$$\begin{aligned} x &= \eta_n s \\ y &= 1/2(y_{l-1,n} + y_{l,n}) \\ z &= 1/2(z_{m-1,n}^* + z_{m,n}^*) + \mu_n s \end{aligned} \quad (116)$$

which generalize Eq. (108). The convention $y_{0,n} = 0$ is built in. To visualize this arrangement of rays, consider the half-plane $x = 0$, $y \geq 0$ partitioned into rectangles by a series of horizontal lines $z = z_{m,n}^*$ and a series of vertical lines $y = y_{l,n}$. Grid n will then consist of a set of rays parallel to each other and to the plane $y = 0$. They will have axial projection μ_n , and each ray of the grid will pass through the center of one of the rectangles.

If $n > N$, grid n consists of the rays of grid $n - N$ traced in the reverse direction.

The μ_n 's for $1 < n \leq N$ must be distinct and fall within the open interval $0 < \mu < 1$, but otherwise they are arbitrary. No generality is lost by requiring that they be in monotonic sequence, so assume that

$$1 = \mu_1 > \mu_2 > \mu_3 \dots > \mu_n > 0$$

$$\mu_{N+n} = -\mu_n, \quad 1 \leq n \leq N \quad (117)$$

In case $n = 1$, Eq. (116) is replaced by

$$x = 0$$

$$y = 1/2(r_{i-1} + r_i) \quad (118)$$

$$z = s$$

so there are just i_{\max} rays in this grid.

Equation (115) must be modified to handle the generalized distribution of rays properly. The factor d^2 occurring in that equation is the area of the unit square in the lattice of rays having a fixed μ_n . This factor must be replaced by

$$a_{\ell, m, n} = (y_{\ell, n} - y_{\ell-1, n}) (z_{m, n}^* - z_{m-1, n}^*) \eta_n \quad (119)$$

when $1 < n < 2N$, which is the area of rectangle associated with ray $R_{\ell, m, n}$, and

$$a_{i, n} = \pi(r_i^2 - r_{i-1}^2)/2 \quad (120)$$

when $n = 1$ or $2N$, which is the area of that half of radial interval i lying in the half-space $y \geq 0$. The factor $1/N$ in Eq. (115) is just $\Delta\mu_n$. Thus, the equation for the rate of deposition of radiant energy in any given hydro zone is estimated to be

$$\dot{E} = 4\pi \sum \Delta\mu_n a_{\ell, m, n} (I_{\text{in}} - I_{\text{out}})_{\ell, m, n} \quad (121)$$

where the sum extends over all rays that intersect the given zone and

$(I_{in} - I_{out})_{\ell, m, n}$ denotes the difference between the incoming and outgoing intensities along ray $R_{\ell, m, n}$.

In the current version of LONG2, the input routine allows only one value of Δy_n and one value of Δz_n^* to be specified. Then, the parametric equations for ray $R_{\ell, m, n}$ are

$$\begin{aligned} x &= \eta_n s \\ y &= (\ell - 1) \Delta y_n \\ z &= z_{min} + (m + \frac{1}{2}) \Delta z_n^* + \mu_n s \end{aligned} \quad (121a)$$

and

$$a_{\ell, m, n} = \begin{cases} \eta_n \Delta y_n \Delta z_n^* & \ell > 1 \\ \eta_n \Delta y_n \Delta z_n^* / 2 & \ell = 1 \end{cases} \quad (121b)$$

No connection between the μ_n 's and the $\Delta \mu_n$'s has yet been established. What has been tacitly assumed is that the range $1 \geq \mu \geq -1$ is divided into abutting intervals and that the n th interval is of length $\Delta \mu_n$ and contains μ_n in its interior. Thus, if μ_n^* denotes the left endpoint of interval n , then

$$\begin{aligned} \mu_n^* &= \mu_{n-1}^* - \Delta \mu_n \\ \mu_0^* &= 1 \\ \mu_N^* &= 0 \\ \mu_{N+n}^* &= \mu_{N-n}^* \\ \mu_{2N}^* &= -1 \\ \mu_n^* &\leq \mu_n \leq \mu_{n-1}^* , \quad 1 \leq n \leq 2N \end{aligned} \quad (122)$$

and

$$\sum_{n=1}^N \Delta \mu_n = 1$$

An obvious source of error connected with using Eq. (119) for estimating rates of energy deposition is the lack of any close connection between the set of rays passing through the system of interest and the HECTIC hydro mesh. If the set of rays corresponding to a particular n were sufficiently dense, then the volume V of any given zone would be nearly $2 \sum_{l,m,n} t_{l,m,n} a_{l,m,n}$, where $t_{l,m,n}$ is the length of the intersection of ray $R_{l,m,n}$ and the given zone. The proof is simple. The half-space $y \geq 0$ may be thought of as filled with non-overlapping tubes of rectangular cross section, in each of which is a ray and $t_{l,m,n} a_{l,m,n}$ is very nearly the volume of the intersection of the tube containing ray $R_{l,m,n}$ and the given zone, providing that $a_{l,m,n}$ is small enough. Consequently,

$$\begin{aligned} V^* &= \sum_{l,m,n} t_{l,m,n} a_{l,m,n} \Delta \mu_n \\ &\approx \sum_n \frac{1}{2} V \Delta \mu_n \\ &= V \end{aligned} \tag{123}$$

The ratio V/V^* , which is denoted by A , is calculated for each hydro zone. Zones for which $A < 1$ intersect more than their share of rays, whereas other zones intersect less than their share. To compensate to some extent for the fact that $A \neq 1$, the strength of the radiative source in each zone is multiplied by A . As was mentioned earlier, the source strength is assumed to be constant within a hydro zone. Therefore, Eq. (102) becomes

$$I_{out} = I_{in} e^{-\sigma t} + AS \int_0^t e^{-\sigma(t-u)} du$$

which, when integrated, gives

$$I_{out} = I_{in} e^{-\sigma t} + \frac{AS}{\sigma} (1 - e^{-\sigma t}) \quad (124)$$

and

$$I_{in} - I_{out} = (I_{in} - \frac{AS}{\sigma}) (1 - e^{-\sigma t})$$

for the contribution of one summand in Eq. (121). Consider what happens when $\sigma t \ll 1$ and $I_{in} = 0$. Then $1 - e^{-\sigma t} \rightarrow \sigma t$ and

$$\begin{aligned} \dot{E} &= 4\pi \sum \Delta\mu_n a_{\ell, m, n} - \frac{AS}{\sigma} \sigma t_{\ell, m, n} \\ &= -4\pi ASV^* \\ &= -4\pi SV \end{aligned}$$

exactly. Thus, including the factor A in Eq. (124) ensures that the correct amount of radiation is emitted in each zone.

The quantity

$$w_{\ell, m, n} = \Delta\mu_n a_{\ell, m, n} \quad (125)$$

is called the "weight" of the ray and

$$J = 4\pi wI \quad (126)$$

is (loosely) called the "power" of a ray of weight w at a point where its intensity is I. Multiplying Eq. (124) by $4\pi w$ yields

$$J_{out} = J_{in} e^{-\sigma t} + \frac{4\pi w AS}{\sigma} (1 - e^{-\sigma t}) \quad (127)$$

and Eq. (121) becomes simply

$$\dot{E} = \sum (J_{in} - J_{out}) \quad (128)$$

The performance of the simplified version of the method is poor when the thickness of the hydro zones is greater than a few mean free paths, because the assumption of a constant source strength within each zone implies a discontinuity at the zone boundaries which, in turn, causes the estimated rate of transfer of energy to be much too high. In LONG2, there is a special treatment of the transfer of energy between thick zones, a thick zone being one for which $\sigma(r_i - r_{i-1})$ and $\sigma(z_j - z_{j-1})$ are both larger than some input criterion C_λ , usually 5 mfp. In the event that two thick zones share a common boundary, all rays crossing that boundary are ignored and the rate of energy transfer between those two zones is estimated to be

$$\dot{E}_d = -\frac{4\pi}{3} A_b \frac{\Delta S/\sigma}{\Delta\tau} \quad (129)$$

where A_b is the area of the common boundary, $\Delta(S/\sigma)$ is the difference in zone centered values of S/σ , $\Delta\tau$ is the optical distance between the zone centers, and the minus sign indicates that the direction of flow is counter to the source gradient.

Consider a thick-thin boundary, i.e., the interface between two zone, one of which is thick and the other is not. First, it may be stated that the treatment of a thin zone is unaffected by the thickness or thinness of its neighbors, Eq. (121) being used in all such cases. However, where a ray passes from a thin zone into a thick one, its power at that point is counted as a contribution to \dot{E} for the thick zone. Where a ray emerges from a thick zone into a thin one, its power is set to

$$J = 4\pi w AS/\sigma \quad (130)$$

where AS/σ is evaluated at the center of the thick zone and $-J$ is added to \dot{E} for the thick zone.

SOURCES

Two types of sources are treated by LONG2: distributed sources and boundary sources. The distributed sources are derived from temperatures calculated in the hydro subroutines by the formula

$$S = \frac{ac}{4\pi} \theta^4 \sigma \left[P\left(\frac{h\nu_2}{\theta}\right) - P\left(\frac{h\nu_1}{\theta}\right) \right] \quad (131)$$

where $ac/4$ is Stefan's constant, θ is the temperature, σ is the absorption coefficient, $\nu_1 \leq \nu \leq \nu_2$ is the band of frequencies being considered, and

$$P(x) = \frac{15}{\pi^4} \int_0^x \frac{u^3 du}{e^u - 1} \quad (132)$$

is the integral of Planck's radiation function. As has already been mentioned, the absorption coefficient and the source are calculated for each zone using zone-central data and are assumed to be constant over the zone.

In the current version of LONG2, boundary temperatures are input on cards and remain constant from time step to time step, but they could just as well be calculated. In any case, the surface temperature is specified for each axial interval $z_{j-1} \leq z \leq z_j$ on the outer cylindrical surface $r = R$ of the system and for each radial interval $r_{i-1} \leq r \leq r_i$ on the bottom $z = z_{\min}$ and top $z = z_{\max}$. The rate at which energy enters the system through each of these intervals is taken to be

$$\dot{E}_{\text{in}} = \frac{ac}{4} \theta_{\text{out}}^4 \alpha \left[P\left(\frac{h\nu_2}{\theta_{\text{out}}}\right) - P\left(\frac{h\nu_1}{\theta_{\text{out}}}\right) \right] \quad (133)$$

where α is the area of the interval. The power carried into the system by a ray originating in some interval is assumed to be

$$J = \frac{w \dot{E}_{\text{in}}}{\sum w} \quad (134)$$

where $\sum w$ is the sum of the weights of all rays originating in that interval.

GENERAL DESCRIPTION OF LONG2

The program LONG2 consists of two major subroutines of HECTIC whose entry points are DRAW and TRAN2. DRAW is entered only when a new hydro mesh has been generated. The result of calling DRAW is to produce for each ray the sequence $(t_p, k_p)_{p=1}^{p_{\max}}$, where p_{\max} is the number of segments into which the ray is divided by boundaries of the hydro mesh, t_p is the length of segment p , and k_p is the index of the zone in which segment p lies. The order of the sequence is the order in which the segments are traversed as the ray is traced. The reverse of the sequence for one ray is the sequence of another ray running in the opposite direction along the same line. DRAW stores these sequences in an I/O file in blocks of up to 1023 words. There is a three-word prologue to each sequence specifying the weight w of the ray and the boundary intervals in which it begins and ends. When the physical system under consideration has a symmetry plane, it is necessary to consider what happens on only one side of it. In such cases, either $z = z_0$ or $z = z_{j\max}$ may be the symmetry plane -- but assume it is the former. Then DRAW sets $z_0^* = z_0$ so that the system of rays will have the same symmetry plane. Rays hitting the plane $z = z_0$ do not then emerge from the system but are reflected there. In symmetrical problems, rays for which $\mu = \pm 1$ require special treatment because they are their own images in a symmetry plane; TRAN2 is instructed not to retrace such rays.

As the ray sequences are generated, the terms in the summation for V^* are accumulated so that when the DRAW file is complete, an array containing $4\pi A$ for each zone may be established. (See Eq. (123).) In a similar manner, the weights of rays intersecting various outer boundary intervals are accumulated for use in evaluating Eq. (134).

When subroutine TRAN2 is called, it begins by calculating absorption coefficients and sources for each hydro zone and the input power for each interval of the outer boundary of the system. It then runs through the

file of ray sequences, calculating the changes in the power level and contributions to zonal rates of deposition of energy. This is done by first using Eq. (134) to get the initial power in the ray on the basis of information contained in the prologue to the sequence to identify the starting interval. Then for each pair (t,k) in sequence, Eq. (127) is applied to estimate the change in J caused by traveling a distance t in zone k . The order of the sequence is the order in which the segments of path are encountered as the ray passes through the system, so it is unnecessary to store the J 's. When J is calculated at some zone boundary, it suffices to subtract it from the running total for \dot{E} for the zone the ray leaves and to add it to that of the zone it enters. Equation (127) is then used to find the change in J as the latter zone is traversed. When a thick zone is encountered, the current J is added to the \dot{E} of that zone in the normal manner. But then, rather than update J , all segments are skipped until a thin zone is encountered, at which point Eq. (130) is used to estimate J and the normal is resumed. After all the sequences have been processed, the zone boundaries are systematically scanned for thick-thick interfaces which, when found, are assumed to transmit energy at a rate specified by Eq. (129). When all the transfer rates have been estimated, the changes in the internal energy of each zone are calculated along with the energy transfer across the outer boundary of the system. The subroutine concludes with the calculation of various summary data used for checking energy balance.

RELATIONSHIP OF LONG2 TO THE PREVIOUS VERSION

As has been stated above, the current version of LONG2 consists of two subroutines: DRAW and TRAN2. These two subroutines are very similar to those with the same names described in reference 4. There are only two significant differences between the old and new versions of LONG2. The first is the introduction into the current version of a facility for assigning weights to rays which makes it possible to use nonuniform

grids. The second difference is the special diffusion-theoretic treatment of thick-thick interfaces in the current version. Such basic features as the over-all organization, the arrangement of data on I/O files, and the flow of control remain virtually unaltered.

AVENUES FOR FURTHER DEVELOPMENT

In this subsection, the question of what might be done with LONG2 to enable it to perform more efficiently the job for which it is designed is examined. The discussion of possibilities for extending the code to include treatments of scattering, radiation pressure, and retardation is beyond the scope of this report. There are two problems that should receive high priority in any plans to continue the development of LONG2. The first concerns the fact that because the rays travel all the way through the system, it is rarely possible to space them closely enough in regions where close spacing is required without having many more rays than necessary everywhere else, particularly at large radii. As a consequence, an excessive amount of computing effort is spent on relatively uninteresting parts of the geometry. The second major problem arises from the fact that the source of radiation is assumed to be constant over a HECTIC hydro zone, which results in an unrealistically high value for the computed rate of transfer of radiant energy. The so-called "view factor method" discussed in later sections of the present report was designed specifically to overcome these obstacles. The remainder of this subsection is a discussion of ways in which LONG2 could be extended to handle them.

The uniformity with which the rays are distributed through the system is a matter of primary importance in the long-characteristic method. If a HECTIC hydro zone is intersected by fewer than its fair share of rays, it will tend to cool faster than it should because it will fail to receive its full share of the radiant energy emitted by its neighbors. The result of

running a problem with an inadequate supply of rays is thus an unevenness in the temperature distribution that persists through all time steps. On the other hand, if the system is too generously infused with rays, the file of ray data will usually be enormous and the time required for processing it prohibitive. Suppose, for example, that the system is divided into a 40-by-40 HECTIC mesh and that the center of each zone coincides with the center of a ray in each of three grids. In such a mesh, a typical ray would be divided into about 80 segments. At two words per segment, the length of the file ray data would be, roughly, $40 \times 40 \times 80 \times 3 \times 2 = 7.68 \times 10^5$ words, which would take about a minute to read on a fast tape drive. More significant is the fact that $3 \times 80 = 240$ segments/zone is about 20 times more than are needed on the average, even though the inner zones would be penetrated by only five or six rays. The situation is usually even worse for large problems and for problems with thin inner zones, because in such cases the number of segments per ray tends to be higher.

The most direct method for allowing the density of rays to be adjusted more or less arbitrarily and independently in various parts of the system is to provide a means for handling rays that have endpoints in the interior of the system. This could be accomplished by allowing a ray to split into two or more parallel subordinate rays. In TRAN2, when a ray being traced arrived at a point where it was split, the energy that it carried could be divided equally among the new subordinate rays to provide them with a sort of initial condition. In the reverse direction, the energy in the subordinate rays could be accumulated at the point of splitting to provide an initial condition for the reverse trace of the main ray. Such a modification, while simple in principle, would considerably increase the complexity of the subroutine DRAW. However, it seems likely that if such a splitting scheme were instituted, it would be possible to allow subordinate rays to be split again without additional complication. The subroutine TRAN2 would not necessarily be much affected by splitting, since the solutions of all the logical problems solved in DRAW could be transmitted to TRAN2 in the file of ray data.

The second of the above-mentioned problems is associated with the assumption of a flat source over each HECTIC zone. A partial solution has already been implemented in the current version of LONG2, namely, the use of diffusion theory to estimate the transfer rate across thick-thick interfaces. One deficiency in this solution is the fact that the numerical results are too strongly dependent upon the value of the parameter used in the criterion for deciding whether or not a zone is thick. The degree of the deficiency could be reduced by varying the source in zones that are, so to speak, "fairly thick" in such a way that the source gradient at interfaces where the criterion for using the diffusion approximation is nearly met is equal to the diffusion theoretic value. Such measures cannot, however, be expected to provide a totally satisfactory solution to this very difficult problem, which still causes trouble in codes for calculating transport in one-dimensional geometries.

SECTION IV
THE PROGRAM LONG2

The programming details of LONG2 are described in the following paragraphs. The appendix contains a listing of the actual FORTRAN statements.

SUBROUTINE DRAW

Flow of Control

The main function of SUBROUTINE DRAW is to provide SUBROUTINE TRAN2 with representative ray paths, along which radiation transport calculations may be performed for each cycle of a HECTIC problem. Each ray path constructed by DRAW is defined by the indices of the hydro zones pierced by the ray and the segment length of the ray in each of these zones, arranged in an ordered fashion. Since these representative ray paths may be a constant of the HECTIC hydro mesh, SUBROUTINE DRAW has been designed to output this ray-path information onto tape so that only one FORTRAN CALL to SUBROUTINE DRAW need be executed at the beginning of a problem, and none for successive restarts.

The majority of SUBROUTINE DRAW is under control of three nested DO-LOOPS: a θ grid loop, a comb loop, and a ray loop. The θ grid loop is completely defined by four FORTRAN variables (EMU, DMU, DZ1, and DR1) which are input on cards for each grid to be constructed by DRAW. A $\theta = 0$ grid is not allowed as an input grid and is always constructed by DRAW after all the input grids have been constructed.

Construction of a grid is begun by extending the system boundaries to form "ray" boundaries which provide limits beyond which no rays of the grid can intersect the system of hydro zones. Combs of the grid are constructed, and the tracing of each ray of a comb for all combs of the grid is initiated.

The tracing of each ray is always started at the ray's point of closest approach to the vertical axis, which is called the "center" of the

ray, and the tracing is done in two passes called the "upward" pass and the "downward" pass. The upward pass is begun in the section of coding labeled "BEGIN UPWARD TRACE," and the downward pass is begun in the section of coding labeled "BEGIN DOWNWARD TRACE." An attempt is made to first trace the ray in an upward direction, and if possible it is traced in that direction until it leaves the system. (If the center of the ray is at or below a reflecting system bottom boundary, the upward trace cannot be performed, since the ray never enters the system.) After the ray has been traced in an upward fashion through its possible reflection at the top system boundary and successive downward path and exit, control is again returned to the ray's center and the downward pass is performed. If the ray happens to have its center above the top boundary which is transmissive, no upward pass is performed and the first pass is a downward pass. Execution of the downward pass is similar to that of the upward pass.

The upward and downward tracing within each pass is performed by sections of the code labeled "GENERAL CASE OF INCREASING Z" and "GENERAL CASE OF DECREASING Z," respectively. There are several smaller blocks of coding that feed into the above two sections which initialize parameters for the beginning of an upward or downward trace, or after a reflection. The FORTRAN variable IUP indicates whether an upward pass or a downward pass has been performed, and tracing is complete when the ray has been completely followed through the system.

If a ray enters the system, the weight WT associated with that ray is added into the appropriate location of one of the FORTRAN arrays BETAA, BETAR, or BETAB, depending on whether the ray enters the top (above), side (right), or bottom system boundary, respectively. This information is made available to SUBROUTINE TRAN2 for use in calculating the initial radiative energy a ray possesses when it enters the system via one of the system boundaries which have black-body temperature profiles associated with them. Rays constructed in DRAW which intersect the vertical axis have half the weight that would normally be assigned to them since in

TRAN2 transport is performed in two directions along these rays and, in actuality, these rays represent only one possible set of rays in the system.

When the tracing of a ray through a hydro zone is performed, the segment length of the ray in that hydro zone and the index of the hydro zone are stored in temporary storage arrays XB and KXA, respectively.

When the trace of a ray is completed, if two passes are actually performed it is necessary to connect the two separate strings of data in such a way as to have a continuous chain of zone indices or segment lengths portraying the path of the ray through the system. Before the indices and segments are merged into one output storage array buffer ($ER \equiv NER$) by a packing section of code, each segment length of the path is added to the appropriate location in the XALP array.

As the rays of each grid are traced and the output storage buffer is filled with indices, segments, the weight WT of each ray, and flags to show through which system boundaries the ray enters and exits, the buffer is output onto tape as its contents reach the limit of 1023 words.

After the input grids have been constructed, the construction of the rays of the $\theta = 0$ grid is treated in essentially the same manner as that of the rays of the input grids. One situation in which the $\theta = 0$ grid is exceptional is in the presence of a symmetry plane. When there is a symmetry plane at the bottom, an ordinary ray will have its center above the plane, and its image will be a different ray whose center is below the plane. But a ray that is parallel to the axis is its own image, and it would be redundant to trace it both forward and backward. TRAN2 is apprised of this fact by setting NXIT=-1 in the prologue of rays for which $\theta = 0$ for symmetrical problems.

At the end of DRAW, a terminal record is written on the tape of ray data and control is returned to MAIN of HECTIC.

Input to DRAW

Medium	Variables
CARDS	CLAMDA ATHETA(I), I=1,IMAX RTHETA(J), J=1,JMAX BTHETA(I), I=1,IMAX ANGN (EMU, DMU, DZ1, DR1), IANG = 1, NANG

Error Termination in DRAW

Type	Indicator	Description of Error
INPUT	S1 = 5.0025	$X(0) > 10^{-10}$
INPUT	S1 = 5.0021	$\cos \theta > .99995$
INPUT	S1 = 5.0022	NANG $\sum_{i=1} \text{DMU} > 1.$
INPUT	S1 = 5.0023	$\text{FIX}(\text{IMAX}/\text{DR1} + 0.8) > 100$; that is, the numbers of combs per grid is > 100
INPUT	S1 = 5.0024	BCBTAG and BCATAG < 0
PROGRAMMING	S1 = 5.0075	--
PROGRAMMING	S1 = 5.0082	--
PROGRAMMING	S1 = 5.0105	--
PROGRAMMING	S1 = 5.0112	--
PROGRAMMING	S1 = 5.0135	--
PROGRAMMING	S1 = 5.0160	--

Results of DRAW

Medium	Variables
COMMON	CLAMDA ATHETA(I), I=1,IMAX RTHETA(J), J=1,JMAX BTHETA(I), I=1,IMAX BETAA(I), I=1,IMAX BETAR(J), J=1,JMAX BETAB(I), I=1,IMAX XALP(K), K=2,KMAX
DRAW TAPE OR DISK	Contains for each ray traced: 1) Index of each hydro zone pierced 2) Length of ray segment in each pierced zone 3) System boundary through which ray enters 4) System boundary through which ray exits 5) Weight associated with ray

DRAW Glossary

FORTTRAN Label	Report Label	Description
ANGN	N-1	The number of polar angles θ to be used; or, equivalently, the number of grids to be constructed by DRAW. (input on a card).
ATHETA(I)	θ_{out}	The temperature of the black-body system boundary adjacent to cell (I, JMAX) (in eV) (input on cards) (used in TRAN2)
BCATAG	None	Boundary flags. A and B refer to top (above) and bottom system boundaries, respectively. If $\theta < 0$, boundary is a perfect reflector; otherwise, it is transmittive to radiation.
BCBTAG	None	
BETAA(I)	None	\sum (weight WT of each ray) all rays passing through the system boundary adjacent to zone (I, JMAX)
BETAB(I)	None	\sum (weight WT of each ray) all rays passing through the system boundary adjacent to zone (I, 1)
BETAR(J)	None	\sum (weight WT of each ray) all rays passing through the system boundary adjacent to zone (IMAX, J)
BTHETA(I)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (I, 1) (in eV) (input on cards) (used in TRAN2)
CLAMDA	c_{λ}	Thick-thin zone criterion

		(in number of mean free paths) (input on a card) (used in TRAN2)
COTAN	None	COT θ
DBGPRT	None	The debug print control: DBGPRT $\geq 10^{-20}$ means debug print desired; otherwise, debug print not desired.
DELS	t	In tracing a ray, the length (or segment) of a ray in a hydro zone. (in cm)
DELX	None	The projection of t onto a horizontal plane (in cm)
DELZ	None	The projection of t onto the vertical axis (in cm)
DMU	$\Delta\mu$	$\Delta\cos\theta$ for a θ grid (input on cards)
DR1	Δy	Radial distance between rays of a θ grid (constant for a grid) (in cm) (input on cards)
DXF	None	$1/\sin\theta$
DY(J)	$z_j - z_{j-1}$	Vertical length of hydro zones (I,J), I = 1, IMAX (in cm)
DZ1	Δz^*	Vertical distance between rays of a θ grid (constant for a grid) (in cm) (input on cards)
DZF	None	$1/\cos\theta$
EMU	μ	$\cos\theta$, the input parameter that specifies a θ grid (input on cards)

ER	None	\equiv NER. The output storage buffer array for ray data that is to be written on the ray tape
ETA	η	$\sin \theta$
FIOUT	None	Saved temporarily on disk in DRAW
I	None	A running index
IANG	None	Running index of the θ grid DO-LOOP
IB	None	Temporary storage for a subscript
ICOR	None	In tracing a ray, ICOR = 1 means the ray hits the corner of a hydro zone ICOR = 0 means the ray does not hit the corner of a hydro zone
IDNMN	None	In merging starting segments of the upward and downward passes, IDNMN = 201 means starting segments for each pass are not in the same hydro zone IDNMN = 202 means starting segments for each pass are in the same hydro zone
IEND	None	A DO-LOOP limit in the $\theta = 0$ trace
IGWD	None	In tracing rays, the index of the last word loaded into the output storage buffer array ER (which is limited to 1023 words). When the storage array is filled and written onto tape, $ER(1) = NER(1) = IGWD$
II	None	A running index
II2	None	\equiv ITOT. The total number of hydro zones a ray penetrates
IIR	None	A running index
IITOT	None	The accumulative number of ray segments generated by DRAW
IM	None	A running index
IMAX	i_{\max}	The number of radial hydro zones
INW	None	Temporary storage for a subscript

IQ	None	The number of radial hydro zones that a ray penetrates in traversing the system
IQQ	None	A running index
IRD	None	Running index of the comb DO-LOOP
IRSTRT	None	Index I of the innermost radial hydro zones (I,J), J=1, JMAX intersected by a comb
IRX	L_n	The number of combs in a θ grid
ISEND	None	An error flag used by SUBROUTINE EDIT
ITOP	None	In tracing a ray, ITOP = 1 means the center of a ray lies at or above the plane Y(JMAX) ITOP = 0 means otherwise
ITOT	None	$\equiv II2$. (see II2)
ITR	None	In tracing a ray, the radial hydro zone which contains the current ray segment
ITWD	None	A running index
IUP	None	The subscript of the last entry made in the arrays KXA and XB: IUP \leq 200 on the upward pass IUP $>$ 200 on the downward pass
IUPM	None	A DO-LOOP limit
IUPUP	None	The last value of IUP on the first (upward) pass
IXX	None	Given XTR, IXX is such that $RX(IXX) > XTR$ and $RX(IXX-1) \leq XTR$
IZRAY	None	Running index of the ray DO-LOOP
J	None	A running index
JB	None	A running index
JMAX	j_{max}	In tracing a ray, the vertical hydro zone which contains the current ray segment
JZX	None	The numbers of rays in a comb
K	None	The subscript of a hydro zone
KMAX	None	KMAX-1 is the total number of hydro zones in the system

KMAXA	None	KMAX+1
KTR	None	In tracing a ray, the index of the hydro zone which contains the current ray segment
KX1	None	Temporary storage for integers
KXA	None	KXA(IUP) contains the subscript of the hydro zone containing the current segment (XB(IUP)) of the ray being traced
L1	None	A DO-LOOP limit
L2	None	A DO-LOOP limit
L3	None	In summing all ray segment lengths contained in a zone to compute XALP(K), L3 = 1 means summing first pass segments L3 = 2 means summing second pass segments
NANG	N-1	Number of θ grids to be constructed by DRAW
NER	None	\equiv ER
NPOINT	None	A pointer used to return control from the ray-packing section of DRAW
NRGP	None	The number of current ray traces in the output storage buffer, ER \equiv NER
NTER	None	NTER = (1/2/3) indicates that a ray enters the system through the (bottom/right/top) boundary
NWD	None	The number of new words to be added to the current accumulation of ray data in the output storage buffer, ER \equiv NER
NXIT	None	NXIT > 0 NXIT = (1/2/3) indicates that a ray leaves the system through the (bottom/right/top) boundary NXIT < 0 NXIT < 0 flags a $\theta = 0$ trace for TRAN2
P	None	Saved temporarily on disk in DRAW
PARIM	None	Temporary storage
PI	π	π
QW	None	Temporary storage

RPT(IRD)	None	The radial coordinate of the "center" (see ZCNTR) of the rays of comb IRD (in cm)
RTHETA(J)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (IMAX, J) (in eV) (input on cards) (used in TRAN2)
R _X	None	R _X (I), I=1,IQ, is the horizontal projection of the distance along a ray from the "center" (see ZCNTR) of the ray to successive radial zone boundaries intersected by the ray (being the same for all rays in a comb) (in cm)
S1	None	Error flag for EDIT When ZCNTR of a ray is not between Y(0) and Y(JMAX), SDELX is the horizontal
SDELX	None	projection of the distance along the ray from its "center" (see ZCNTR) to the point where it enters the system, and SDELZ is the axial projection (in cm)
SDELZ	None	
SMU	None	$1 - \sum_{i=1, IANG} DMU$
TAU(I)	$\pi(r_i^2 - r_{i-1}^2)$	The area of the annulus defined by radial hydro interval I (in cm ²)
U	None	≡ROSS. Saved temporarily on disk in DRAW
V	None	≡KXA. Saved temporarily on disk in DRAW
WT	$\omega_{l,m,n}$	The weight associated with a given ray: If $\theta \neq 0$, $WT = (\Delta \cos \theta \Delta Z \Delta R \sin \theta)(F), \text{ where}$ $F = 1 \text{ if } IRD \neq 1 \text{ and}$ $F = .5 \text{ if } IRD = 1$ If $\theta = 0$, $WT = 1/2 \sin \theta \text{ TAU}(I)$

WTP	None	Temporary storage, used in the WT calculation for $\theta = 0$ trace
X(I)	r_i	Outer radius of radial hydro zone I (in cm)
XALP(K)	$4\pi A$	$= 2 * \text{TAU}(I) * \Delta y * \pi / \sum_{\text{all rays passing through zone K}} [(\text{length of ray in zone K}) * (\text{weight WT of ray})]$
XB	None	XB(IUP) contains the segment length in the current zone (KXA(IUP)) of the ray being traced, traced
XTR	None	In tracing a ray, XTR is the horizontal component of the distance along a ray from the "center" (see ZCNTR) of the ray to the beginning of the current ray segment (in cm)
XX2	None	Temporary storage
Y(J)	z_j	The vertical coordinate of the upper boundary of hydro zones (I,J), $I=1, \text{IMAX}$ (in cm)
ZCNTR	None	The axial coordinate of the point of closest approach of a ray to the vertical axis, or the "center" of a ray (in cm)
ZEND	None	The axial coordinate of the "center" (see ZCNTR) of ray (IZRAY=JZX) (in cm)
ZSTRT	None	The axial coordinate of the "center" (see ZCNTR) of ray (IZRAY=1) (in cm)
ZTR	None	In tracing a ray, ZTR is the axial coordinate of the beginning of the current ray segment (in cm)
ZXTRA	None	$X(\text{IMAX}) * \cot \theta$, used in the calculation of ZSTRT and ZEND (in cm)

SUBROUTINE TRAN2Flow of Control

The main function of SUBROUTINE TRAN2 is to perform transport calculations along the representative ray paths supplied by SUBROUTINE DRAW, extracting and depositing radiative energy from the HECTIC hydro zones according to the temperature and optical depth of each zone.

Upon entering SUBROUTINE TRAN2 from MAIN, the integrated black-body intensity $B(K)$ for each hydro zone is calculated if the zone is considered "active" by HECTIC standards.

Next, the black-body boundary conditions are calculated. The amount of energy E_{SURF} entering the system boundary through zone K from a black-body boundary at temperature θ (ATHETA, RTHETA, or BTHETA in the code) is given by Eq. (133):

$$E_{SURF} = \frac{ac}{4} \theta^4 \alpha \int_{h\nu_1/\theta}^{h\nu_2/\theta} \frac{15}{\pi^4} \frac{u^3 du}{e^u - 1}$$

└──────────┘
FUNCTION PLNKUT in
the code

where $ac/4$ = Stefan's constant

α = the area which zone K presents to the black-body boundary

ν_1, ν_2 = the frequency limits ($h\nu_1 = 0.001$ eV, $h\nu_2 = 10^6$ eV in the monofrequency problem)

h = Planck's constant

The amount of this energy apportioned to each constructed ray (with weight WT_{RAY}) which enters the system through zone K is given by

$$E_{RAY} = \frac{E_{SURF} \cdot WT_{RAY}}{\sum_{\text{all rays}} WT}$$

passing from the
boundary into zone K

In the section of coding prior to the transport calculations, only the quantities $E_{\text{SURF}}/\Sigma \text{ WT}$ are calculated (in the code, PA, PR1, PB).

Also prior to the transport calculations, the Rosseland mean absorption coefficients are calculated (ROSS(K)) for each hydro zone, and the diffusion flags are set. If a zone is less than CLAMDA (input) mean free paths thick or is inactive, it is considered "optically thin;" otherwise, it is labeled "optically thick" (the flags occupy the ALAMV array).

The two-dimensional transport calculation section is based around three nested loops: the "RAY-TAPE RECORD LOOP," the "RAY LOOP," and the "SEGMENT LOOP." The input storage buffer array ($\text{XB} \equiv \text{KXB}$) is filled with a logical record of ray data from the ray tape which contains data from traces of one or more rays constructed in SUBROUTINE DRAW. The "RAY LOOP" is then entered, and initialization for transport begins for each ray trace in the buffer. If the ray traverses a totally inactive region, the transport calculations are skipped and the next ray trace in the buffer is considered.

In general two "passes" or "transports" are performed for each ray trace supplied by DRAW. The first pass (IPASS = 1) is performed by picking up zone indices from the array KXA in an increasing manner, and the second pass (IPASS = 2) is performed "backwards" in a sense to the first pass along the same ray trace supplied by DRAW. If the ray is vertical ($\theta = 0$, NTER = -1), only one pass is performed. For each pass, the rate at which the hydro zones are gaining energy due to radiation is calculated (ER(K)), the rate at which the system is gaining energy through its boundaries is calculated (EBTM, ETOP, ESIDE), the total rate at which the system is gaining energy is calculated (FOO), and the W2 array is updated.

In the "SEGMENT LOOP," if the ray enters the system through a black-body boundary the rate at which energy is transported into the zone containing its first segment is given by

$$E_{\text{RAY}} = \begin{pmatrix} \text{PB(I)} \\ \text{PRI(J)} \\ \text{or} \\ \text{PA(I)} \end{pmatrix} * \text{WT}_{\text{RAY}}$$

In transporting along the trace of a ray, the diffusion flag of the zone in which its next segment lies is checked. If the zone is considered "thick," then all the energy/(DT) transported by the ray is emptied into that zone. As the trace continues, no deposition calculation is performed and the diffusion flags of the zones are checked until a "thin" zone is encountered. Then the rate at which energy is leaving the last "thick" zone is given by

$$E_{\text{RAY}} = B(K) \cdot \text{XALP}(K) \cdot \text{WT}_{\text{RAY}}$$

The outside of the system is treated as if it were "thin." For a sequence of segments in "thin" zones, the transport is calculated in the section of coding labeled "NORMAL TRANSPORT CALCULATION" under the equation

$$E_{\text{OUT}} = E_{\text{IN}} e^{-\sigma s} + c_1 c_2 (1 - e^{-\sigma s}) \text{WT}_{\text{RAY}}$$

- where E_{OUT} = rate at which energy is leaving zone K containing segment length s
- E_{IN} = rate at which energy is entering zone K containing segment length s
- σ = Rosseland mean absorption coefficient for zone K (ROSS(K))
- c_1 = XALP(K)
- c_2 = B(K)
- WT_{RAY} = weight associated with the ray
- s = segment length in zone K

After transport calculations have been performed along all the rays constructed by DRAW, an explicit diffusion calculation is performed across all "thick-thick" zone interfaces.

At the end of SUBROUTINE TRAN2, the ER array is checked to see if radiation transport in time DT will change the internal energy of a hydro zone by more than the allowed amount (SLUG); if it will, DT is reduced in TRAN2 so as to comply with this restriction. Finally, the internal energy of the system, ETH, is updated to include energy changes by radiation transport, and the variables BACC, SACC, and TACC are updated.

The outside loop in TRAN2 is the frequency loop which is closely patterned after the one in SPUTTER (Ref. 5). The multifrequency option has never been tried, however, in any of the radiation subroutines (TDRAD, TRAN2/LONG2, and TRAN2/SHORT2) that have been incorporated into HECTIC.

Input to TRAN2

Results from SUBROUTINE DRAW are transferred to TRAN2. (See the subsection entitled "Results of DRAW.")

Error Termination in TRAN2

Type	Indicator	Description of Error
INPUT	S1 = 7.0225	MERGE \leq 0 and multifrequency problem
INPUT	S1 = 7.0235	IHNU < 1
INPUT	S1 = 7.0250	Q \geq 0.6 or Q \leq 0.4
INPUT	S1 = 7.0320	XA(K) < 0.
EXECUTION/ INPUT	S1 = 7.0522	An indefinite number has been created before the call to DVCHK
INPUT	S1 = 7.0541	There are more than 210 ray traces for an input buffer-load of ray data from DRAW (it is impossible for DRAW to create correctly more than 170 traces in one buffer-load)

INPUT	S1 = 7.0550	In transporting along a ray, NTER is not 1, 2, or 3
INPUT	S1 = 7.0556	NTER = 1 and BCBTAG < 0.
EXECUTION/ INPUT	S1 = 7.0810	An indefinite number has been created before the call to DVCHK
INPUT	S1 = 7.1010	IHNU > NHNU
EXECUTION	S1 = 7.1068	DT < FFB, the minimum time-step control

Results of TRAN2

Medium	Variables
COMMON	(A possible contribution to)DT (A contribution to)ETH ER(K), K=2, KMAX W2(I), I=1, IMAX

TRAN2 Glossary

FORTTRAN Label	Report Label	Description
AH	None	In the explicit diffusion calculation, a variable used in the calculation of net diffusion in the horizontal direction
AIX(K)	None	The total internal energy of zone K per unit mass (does not include kinetic energy) (in ergs/gm)
AJ	J	In tracing a ray through a hydro zone of the system, the rate at which energy is being transported into or out of the zone (in ergs/sec)
AJN	J	Similar in function to AJ (in ergs/sec)

ALAMV(K)	None	For the current problem cycle, ALAMV(K) = 1 means zone K is optically thick and active ALAMV(K) = 0 means zone K is optically thin or inactive
AMX(K)	None	The mass of zone K (in gm)
ATHETA(I)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (I, JMAX) (in eV) (input in DRAW)
AV	None	In the explicit diffusion calculation, a variable used in the calculation of net diffusion in the vertical direction
B(K)	$\frac{S}{\sigma}$	In the monofrequency calculation, the inte- grated black-body intensity associated with zone K (in $\frac{\text{ergs}}{\text{sec-steradian-cm}^2}$)
	None	In the multifrequency calculation, temporary storage
BACC	None	$\int_{t'=0}^{t'=t}$ (rate at which energy is leaving the system bottom boundary due to radiation) dt' (in ergs)
BCATAG	None	Boundary flags. A, B, and R refer to top (above), bottom, and side (right) system boundaries, respectively. If $\theta < 0$, boundary is a perfect reflector; otherwise, it is trans- mittive to radiation (the side cannot be a reflector)
BCBTAG	None	
BCRTAG	None	
BETA	$\frac{h\nu_1}{\theta}$	In the multifrequency calculation, the lower limit of the current frequency band

BETAA(I)	None	\sum (weight of each ray) all rays passing through the system boundary adjacent to zone (I, JMAX) (calculated in DRAW)
BETAB(I)	None	\sum (weight of each ray) all rays passing through the system boundary adjacent to zone (I, 1) (calculated in DRAW)
BETAP	$\frac{h\nu_2}{\theta}$	In the multifrequency calculation, the upper limit of the current frequency band
BETAR(J)	None	\sum (weight of each ray) all rays passing through the system boundary adjacent to zone (IMAX, J) (calculated in DRAW)
BTHETA(I)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (I, 1) (in eV) (input in DRAW)
CLAMDA	c_λ	Thick-thin zone criterion: If an active zone K has vertical and radial dimensions \geq CLAMDA/ ROSS(K), then it is considered optically thick; otherwise, it is considered optically thin (in number of mean free paths) (input in DRAW)
DBGPRT	None	The debug print control: DBGPRT \geq 10 ⁻²⁰ means debug desired; otherwise, debug print not desired
DE	None	In calculating DT in TRAN2, the amount of energy per unit mass which a zone will lose or gain in time DT due to radiation (in ergs/gm)

DERH	\dot{E}_d	In the explicit diffusion calculation, the net diffusion rate in a horizontal direction between zones (in ergs/sec)
DERV	\dot{E}_d	In the explicit diffusion calculation, the net diffusion rate in a vertical direction between zones (in ergs/sec)
DFB	None	In the multifrequency calculation, $\frac{15}{\pi^4} \int_{\text{BETA}}^{\text{BETAP}} \frac{u^3}{e^u - 1} du$
DHNU	None	In the multifrequency calculation, the width of the current frequency band
DT	None	Time step (which can be controlled by TRAN2) (in sec)
DTEMP	None	Temporary storage
DX(I)	$r_i - r_{i-1}$	Radial zone length of zones (I, J), J=1, JMAX
DY(J)	$z_j - z_{j-1}$	Vertical zone length of zones (I, J), I=1, IMAX
E	None	In the temperature iteration calculation, the total internal energy per unit mass updated by radiation heating only (in ergs/gm)
EBTM	None	The rate at which energy is entering the system through the system bottom boundary in time DT due to radiation (in ergs/sec)
EMB1	None	In the multifrequency calculation, temporary storage
EMB2	None	Same as above
ER(K)	None	The rate at which zone K is gaining energy due to radiation in time DT (in ergs/sec)

	None	In the multifrequency calculation, temporary storage
ESIDE	None	The rate at which energy is entering the system through the system side boundary in time DT due to radiation (in ergs/sec)
ESS	None	Temporary storage used in the "normal transport calculation"
ETH	None	The total energy in the system at time T (in ergs)
ETOP	None	The rate at which energy is entering the system through the system top boundary in time DT due to radiation (in ergs/sec)
FACTOR	None	Temporary storage used in the calculation of ROSS(K) and PLANCK(K)
FFB	None	The minimum time-step allowed by HECTIC (in sec)
FIOUT	None	= OLDTH Saved temporarily on disk in DRAW and restored at the end of TRAN2
FOO	None	The amount of energy gain by the system in time DT due to radiation (in ergs)
GG	None	A dummy argument in the call to ES in TRAN2
HNU	$h\nu_1$	In the multifrequency calculation, the lower frequency of the current frequency band (in eV)
HNU4	None	= HNU ⁴ .
HNUP	$h\nu_2$	In the multifrequency calculation, the upper frequency of the current frequency band. (in eV)
HNUP4	None	= HNUP ⁴ .
I	None	A running index.

IACT	None	<p>An activity flag:</p> <p>IACT = 1 means all the zones traversed by a ray constructed in DRAW are inactive, and no transport is performed</p> <p>IACT = 2 means some or all of the zones traversed by a ray constructed in DRAW are active, and transport is performed</p>
IDY	None	Running index of the segment loop for a ray trace
IGRP	None	Running index of the ray-trace loop
IGWD	None	In tracing a ray within the ray-tape record loop, a pointer to the end of the ray data in the input storage array XB (\equiv KXB) for the current ray
IHNU	None	In the multifrequency calculation, the index of the current frequency band
II	None	A running index
IMAX	i_{\max}	The number of radial hydro zones
IPASS	None	<p>In transporting along a ray,</p> <p>IPASS = 1 means transport is being performed along a ray in a direction specified by the increasing order of zone indices specified in the KXB array</p> <p>IPASS = 2 means transport is in the opposite direction to that of IPASS = 1</p>
ISEND	None	An error flag used by SUBROUTINE EDIT
ITAG	None	<p>Flag for the temperature iteration calculation:</p> <p>= 0 means perform temperature iteration</p> <p>\neq 0 means do not perform temperature iteration</p>
ITOT	None	In transporting along a ray, the number of hydro zones that a ray penetrates in one pass through the system
ITR	None	In transporting along a ray in the ray segment loop, the pointer to the location in the KXA (\equiv XA) array containing the current zone index (ITR+400 points to the location in the XA (\equiv KXA) array containing the corresponding ray segment length)

J	None	A running index
JMAX	j_{\max}	The number of vertical hydro zones
K	None	A running index
KDMY	None	Argument returned by the calls to FUNCTION D'CHK: = 1 means an error = 2 means no error
KFIT	None	An array constructed by HECTIC such that: JMR(KFIT(K),2) = 1 means zone K is active JMR(KFIT(K),2) \neq 1 means zone K is inactive
KMAX	None	KMAX-1 is the total number of hydro zones in the system
KMAXA	None	KMAX+1
KOLD	None	In transporting along a ray in the ray-segment loop, the index of the zone which contains the previous ray segment
KPP	None	In the explicit diffusion calculation, the index of the zone above zone K
KSTRT	None	In transporting along a ray, the index of the first hydro zone which the ray intersects upon entering the system
KT	None	A running index
KXA	None	\equiv XA. In transporting along a ray, the storage array which contains the ray data for the current ray. KYA(K) contains the index of the hydro zone pierced by the ray (the corresponding ray segment length is contained in XA(K+400))
KXB	None	\equiv XB. The input storage buffer array for DRAW ray data (contains traces of one or more rays)
M	None	Not used
MERGE	None	In the multifrequency calculation, a variable controlling the merging of frequency bands

MFTAG	None	Multifrequency flag: = 0 means monofrequency problem ≠ 0 means multifrequency problem
N	None	Not used
NC	None	Integer value of cycle number
NGRP	None	The number of ray traces in the current buffer-load of ray data (XB ≡ KXB)
NGWD	None	The number of words in the current buffer-load of ray data (XB ≡ KXB)
NHNU	None	In the multifrequency calculation, the total number of frequency bands
NTER	None	NTER = (1/2/3) indicates that a ray "enters" (depending on the pass) the system through the (bottom/side/top) boundary
NVEZ	None	In the temperature iteration calculation, the temperature iteration counter
NXIT	None	NXIT > 0 NXIT = (1/2/3) indicates that a ray "leaves" (depending on the pass) the system through the (bottom/side/top) boundary NXIT < 0 means θ = 0 trace
NY	None	Not used
OLDTH	None	≡ FIOU In the temperature iteration calculation, temporary storage
P	None	≡ PLANCK. Saved temporarily on disk in DRAW and restored at the end of TRAN2
PA(I)	None	The rate at which energy is entering zone (I, JMAX) along a ray from the adjacent black-body system boundary at temperature ATHETA(I), divided by the weight WT of the ray (in ergs/sec)

PE(I)	None	The rate at which energy is entering zone (I, 1) along a ray from the adjacent black-body system boundary at temperature BTHETA(I), divided by the weight WT of the ray (in ergs/sec)
PI	π	π
PLANCK(K)	None	In transporting along a ray through hydro zone K, the rate at which energy is entering zone K from the zone previously traversed (in ergs/sec) Also, the PLANCK mean opacity across the merged frequency band (in 1/cm)
PRI(J)	None	The rate at which energy is entering zone (IMAX, J) along a ray from the adjacent black-body system boundary at temperature RTHETA(J), divided by the weight WT of the ray (in ergs/sec)
PUR	None	Not used
PUZ	None	Not used
Q	None	In the calculation of DT in TRAN2, the amount by which the internal energy per unit mass of a zone is allowed to change due to radiation transport (in ergs/gm) Also, in the multifrequency calculation, a parameter used in the merging of frequency bands
ROSS(K)	σ	The Rosseland mean absorption coefficient for zone K (in 1/cm)
RPTAG	None	Absorption coefficient flag: $\neq 0$ means set the Planck mean absorption coefficient equal to the Rosseland mean absorption coefficient $= 0$ means keep both absorption coefficients distinct

RTHETA(J)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (IMAX, J) (in eV) (input in DRAW)
RUR	None	Not used
RUZ	None	Not used
S1	None	Error flag for EDIT
SACC	None	$\int_{t'=0}^{t'=t}$ (rate at which energy is leaving the system bottom boundary due to radiation) dt' (in ergs)
SIGS	None	Temporary storage
SLUG	None	In the calculation fo DT in TRAN2, the fraction by which the internal energy per unit mass of a zone is allowed to change due to radiation transport
SV	None	Temporary storage
T4	None	In merging frequency bands, = THETA(K) ⁴ (in eV ⁴)
TACC	None	$\int_{t'=0}^{t'=t}$ (rate at which energy is leaving the system top boundary due to radiation) dt' (in ergs)
TAU(I)	$\pi(r_i^2 - r_{i-1}^2)$	The area of the annulus defined by radial hydro interval I (in cm ²)
TEMP(1)	None	In changing DT in TRAN2, the computed fraction by which DT is reduced to comply with the restrictive variable SLUG
THETA(K)	θ	Temperature of zone K (in eV)

THTAMX	None	In the multifrequency calculation, the highest temperature in the system used in deciding whether to merge frequency bands (in eV)
U	None	\equiv ROSS. Saved temporarily on disk in DRAW and restored at the end of TRAN2
V	None	\equiv KXA \equiv XA. Saved temporarily on disk in DRAW and restored at the end of TRAN2
VEZ	None	Not used
W2(I)	None	The rate at which energy is leaving the system top boundary adjacent to zone (I, JMAX) due to radiation, divided by TAU(I) (in ergs/cm ² sec)
WT	$\omega_{l,m,n}$	The weight associated with a given ray (see "DRAW Glossary").
X(I)	r_i	Outer radius of radial hydro zone I (in cm)
XA	None	\equiv KXA (see KXA)
XALP(K)	$4\pi A$	$= 2 * \text{TAU}(I) * \Delta y * \pi / \sum$ (length of ray in zone K) all rays *(weight WT of ray) passing through zone K
XB	None	\equiv KXB. (see KXB)
Y(J)	z_j	The vertical coordinate of the upper boundary of hydro zones (I, J), I=1, IMAX (in cm)

SECTION V

THE VIEW-FACTOR METHOD

INTRODUCTION

The view-factor method is proposed as an economical method for calculating radiative transport of energy in situations where diffusion theory is inapplicable and higher-order methods are unfeasible. The problem of distributing rays does not exist because the concept of rays as characteristics is not used. A procedure for interpolating the source strength between zone centers that is consistent with diffusion theory is a feature of the method.

Let Z be a hydro zone in HECTIC and let B_b , $b = 1, 2, 3, 4$, denote its boundaries. The sphere of directions of rays passing through any point of the system is partitioned into bundles, called "cones" and denoted by C_a , $a = 1, 2, \dots, A$. Currently, the assignment of directions to cones at various points is the same for all points of the system relative to the local coordinate system. This means that if at point P a direction with axial projection μ and azimuth ϕ is in cone C_a , then at any other point, say Q , the direction with the same axial projection and azimuth is also in cone C_a . Cones are to be constructed so that the rays of a cone are all rising or all descending and all approaching or all receding from the axis as they pass through the vertex. The function of this restriction is to ensure that if one ray in cone C_a with vertex P on B_b enters Z at P , then all rays of that cone do the same. The specification of cones used in SHORT2 is discussed in the next subsection.

The basic construct of the view-factor method is an object called a "fan." The fan $F_{a,b}$ consists of all rays that intersect boundary B_b of zone Z and belong to the cone C_a with vertex at the point of intersection. If fan $F_{a,b}$ leaves zone Z through boundary B_b and fan $F_{a',b'}$ enters zone Z

through boundary $B_{b'}$, then the view factor $\eta_{a,b,a',b'}$ is, roughly, the ratio of the number of rays in fan $F_{a',b'}$ to the number that are in both fans. (Rays that leave and reenter zone Z while passing from $B_{b'}$ to B_b are not counted.) The meaning of "number of rays in fan" is straightforward: the number of rays passing through an element of surface dS with directions lying in an infinitesimal cone of solid angle $d\Omega$ that makes an angle θ with the normal to the surface is $N_0 \cos \theta d\Omega dS$, where N_0 is the density of rays, which is assumed to be constant over the system.

To illustrate the concepts "fan" and "view factor," let Z be a HECTIC hydro zone consisting of those points of the system with radii lying in the interval $r_1 < r < r_2$ and aptitudes lying in the interval $z_1 \leq z < z_2$. Let B_b and $B_{b'}$ denote, respectively, the inner and outer boundaries of Z . For any point P on the cylinder $r = r_1$, let the cone C_a with vertex at P consist of those rays through P with axial projections in the range $\mu_1 \geq \mu \geq \mu_2$ that intersect the cylinder $r = \rho_2$ but not the cylinder $r = \rho_1$, where $\rho_1 \leq \rho_2 \leq r$, and that are going inward as they pass through P . Then the fan $F_{a,b}$ consists of all rays that intersect B_b going inward, that have axial projection in the range $\mu_1 \geq \mu \geq \mu_2$, and that intersect the cylinder $r = \rho_2$ but not the cylinder $r = \rho_1$. The number of rays in fan $F_{a,b}$ is

$$N = 2\pi r_1 N_0 \int_{\mu_1}^{\mu_2} \int_{\phi_1}^{\phi_2} (1 - \mu^2)^{\frac{1}{2}} \cos \phi \, d\phi d\mu$$

where $r_1 \sin \phi_k = \rho_k$. For any point P' on the cylinder $r = r_2$, let the cone $C_{a'}$ with vertex at P' consist of those rays through P' with axial projections in the range $\mu_1 \geq \mu \geq \mu_2$ that intersect the cylinder $r = \rho_2'$ but not the cylinder $r = \rho_1'$, where $\rho_1 \leq \rho_1' \leq \rho_2 \leq \rho_2' \leq r_2$, and that are going inward as they pass through P' . The fan $F_{a',b'}$ consists of all rays that intersect $B_{b'}$ going inward, that have axial projection in the range $\mu_1 \geq \mu \geq \mu_2$, and that intersect the cylinder $r = \rho_2'$ but not the cylinder $r = \rho_1'$. Let N' be the number of rays in fan $F_{a',b'}$ and let M be the number of rays common to both fans, i. e., the number of rays that intersect both B_b and $B_{b'}$ going

inward, that have axial projections in the range $\mu_1 \geq \mu \geq \mu_2$, and that intersect the cylinder $r = \rho_2$ but not the cylinder $r = \rho_1$. Then $\eta_{a,b,a',b'} = M/N'$ is the view factor connecting the two fans.

Two properties of view factors follow from the definition:

$$\sum_{a,b} \eta_{a,b; a',b'} = 1 \quad (135)$$

and

$$\sum_{a',b'} \eta_{a,b; a',b'} k_{a',b'} = k_{a,b} \quad (136)$$

where $k_{a,b}$ is the number of rays in fan $F_{a,b}$.

Suppose that zone Z is filled with material having the absorption coefficient σ , and let $J_{a,b}$ be the amount of radiant energy crossing boundary B_b along rays in fan $F_{a,b}$. Then the view-factor method estimates $J_{a,b}$ using the formula

$$J_{a,b} = e^{-\sigma t_{a,b}} \sum_{a',b'} \eta_{a,b; a',b'} J_{a',b'} + S_{a,b} \quad (137)$$

where $S_{a,b}$ is the rate at which energy radiated in zone Z leaves through boundary B_b along rays in fan $F_{a,b}$, and $t_{a,b}$ is an average distance traveled by rays in fan $F_{a,b}$ in traversing zone Z. The rate of energy deposition in zone Z is then calculated to be

$$\dot{E} = \sum_{a,b} \pm J_{a,b} \quad (138)$$

where the + sign is used if fan $F_{a,b}$ enters zone Z through boundary B_b and the - sign is used if it leaves.

It will be shown later that when $J_{a,b}$ is given for incoming fans associated with the outer boundary of the system, then all others may be determined by a systematic application of Eq. (137).

Some of the principal advantages of the method are as follows:

1. The equations can be solved by a straightforward, noniterative method.
2. The coefficients may be chosen so that the diffusion limit is approached in a continuous, natural manner.
3. The coefficients are all positive. Rapid variation of intensity with direction causes no problems, and negative fluxes cannot occur.

While the view-factor method allows much latitude in the assignment of directions to cones, a simple, rather uneconomical scheme has been chosen for the first version of SHORT2 described here. For some fixed positive integer M , the sphere of directions at each point P of the system is divided into cones $C_{l,m}$, $1 \leq l, m \leq 2M$, consisting of those directions $\vec{\Omega}$ for which

$$\pi \frac{l-1}{2M} \leq \phi \leq \frac{\pi l}{2M} \quad (139)$$

$$\cos \pi \frac{m-1}{2M} \geq \mu \geq \cos \frac{\pi m}{2M}$$

where ϕ is the azimuth and μ the axial projection of the ray passing through P in direction $\vec{\Omega}$. As in the long-characteristic method, the hemisphere of directions of negative azimuth is ignored. The main disadvantage to this type of subdivision is its lack of uniformity; the cones which are more or less vertical have smaller solid angles than do the horizontal ones.

A scheme due to Carlson (Ref. 6) that yields a set of cones of equal solid angle is as follows. With M as before, let $1 = \mu_0 > \mu_1 \cdots > \mu_M = 0$, and let

$$\mu_{m-1} - \mu_m = m(1 - \mu_1) \quad (140)$$

Thus

$$\begin{aligned} 1 &= \mu_0 - \mu_M \\ &= \sum_{m=1}^M \mu_{m-1} - \mu_m \\ &= (1 - \mu_1) \sum_{m=1}^M m \\ &= \frac{1}{2} M(M+1) (1 - \mu_1) \end{aligned}$$

or

$$1 - \mu_1 = \frac{2}{M(M+1)} \quad (141)$$

For $M < m \leq 2M$, let

$$\mu_{2M-m} = -\mu_m \quad (142)$$

Then, for $l \leq 2m \leq 2M$, let

$$\phi_{l,m} = \frac{\pi l}{2m} \quad (143)$$

and define cone $C_{l,m}$ to be those directions $\vec{\Omega}$ for which

$$\phi_{l-1,m} \leq \phi \leq \phi_{l,m} \quad (144)$$

$$\mu_m \leq \mu \leq \mu_{m-1}$$

If $M < m \leq 2M$, let $m' = 2M - m + 1$; and for $l \leq 2m'$, define cone $C_{l,m}$ to be the reflection of $C_{l,m'}$ in a horizontal plane through P. Thus,

$$\phi_{l-1,m'} \leq \phi \leq \phi_{l,m'} \quad (145)$$

$$\mu_m \leq \mu \leq \mu_{m-1}$$

These cones all subtend the same solid angle because

$$\Delta \vec{\Omega} = \Delta \mu \Delta \phi = m(1 - \mu_1) \frac{\pi}{2m} = \frac{\pi}{M(M+1)}$$

for $m \leq M$ and similarly for $m > M$. Figure 4 shows how the octant $0 \leq \mu \leq 1$, $0 \leq \phi \leq \frac{1}{2}\pi$ is subdivided when $M = 4$ in a projection in which meridians, $\phi = \text{const}$, become concurrent lines and parallels, $\mu = \text{const}$, become parallel lines.

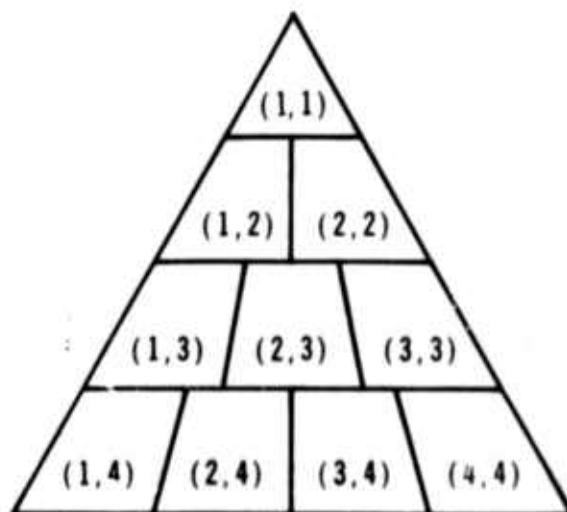


Figure 4. Subdivision of an Octant of the Sphere of Directions

THE MODIFIED DEFINITION OF FAN

The definition of fan given previously is precise but more restrictive than desirable or necessary. Basically, (1) a fan is a set of rays that cross some associated zone boundary in directions that, at the point of crossing, have roughly the same axial projection and azimuth, and (2) fans associated with a given zone boundary are nonoverlapping. The particular method used in SHORT2 for assigning rays to fans is a very slight modification of the one already described. It leads to fewer terms in the summation of Eq. (137) and simplifies the problem of estimating view factors. Only those fans associated with horizontal boundaries are affected. The fan F associated with cone $C_{l,m}$ and boundary $BA_{i,j}$ consists of those rays which intersect $BA_{i,j}$ and which cross the cylinder $r = r_i$ in cone $C_{l,m}$. Since each ray that penetrates that cylinder crosses it twice, there is an ambiguity to be resolved; if $C_{l,m}$ points $\begin{pmatrix} \text{inward} \\ \text{outward} \end{pmatrix}$, then rays of F cross $BA_{i,j}$ while going $\begin{pmatrix} \text{inward} \\ \text{outward} \end{pmatrix}$. The fan associated with vertical zone boundary $BR_{i,j}$ and cone $C_{l,m}$ is still the set containing any ray that intersects $BR_{i,j}$ at a point where its direction is in cone $C_{l,m}$.

It will prove convenient to have a more compact notation available for the ensuing considerations. Let q denote an index combination for fans, and let F_q be the fan associated with cone

$$C_q = C_{l_q, m_q}$$

and boundary

$$B_q = BX_{i_q, j_q}^q,$$

where $X^q = A$ or R . ($BA_{i,0} \equiv BB_{i,1}$.) Let Z_q^i denote the zone that the rays of fan F_q enter as they cross B_q , and let Z_q^o denote the zone they leave. Equation (137) in this notation, becomes

$$J_q = e^{-\sigma t} \sum_{q'} \eta_{q,q'} J_{q'} + S_q \quad (146)$$

where J_q is the rate of transfer of radiant energy across boundary B_q along rays of fan F_q , and S_q is the rate at which energy radiated in zone Z_q^0 crosses boundary B_q along rays of fan F_q .

The View-Factor Integral

Let $k_{q'}$ denote the number of rays in fan $F_{q'}$, assuming that the density of rays is unity. Then

$$k_{q'} = \iint_{F_{q'}} \vec{\Omega} \cdot \vec{n} \, d\Omega \, dS \quad (147)$$

where the surface integral extends over $B_{q'}$ and for each point on $B_{q'}$, the direction integral extends over directions in cone $C_{q'}$ and \vec{n} is a unit vector normal to $B_{q'}$. If F_q is another fan and if $Z_{q'}^1 = Z_q^0$ (i.e., if the zone that $F_{q'}$ enters through $B_{q'}$ is the zone F_q leaves through $B_{q'}$), then

$$\eta_{q,q'} = \frac{1}{k_{q'}} \iint_{F_q \cap F_{q'}} \vec{\Omega} \cdot \vec{n} \, d\Omega \, dS \quad (148)$$

where, for each point P on $B_{q'}$, the direction integral now extends over directions of rays through P that are common to both fans.

Estimation of the View Factors

In principle, the view factors can be calculated using Eq. (148). Such a procedure involves performing a four-fold integral, two of which are not at all trivial. Evaluating such integrals for each pair of fans would consume much more time than can be justified by the application for which the

results are intended. The current version of the DRAW subroutine of SHORT2 estimates view factors on the assumption that all directions $\vec{\Omega}$ in cone $C_{l,m}$ have the same projection $\mu_{m-\frac{1}{2}}$ on the axis of the system, where $\mu_m < \mu_{m-\frac{1}{2}} < \mu_{m-1}$, which eliminates one of the two troublesome integrals. The problem of how to optimize the choice of $\mu_{m-\frac{1}{2}}$ has not yet received any attention. The currently used choice is

$$\mu_{m-\frac{1}{2}} = \cos \frac{2m-1}{4M} \pi$$

which is the cosine of the polar angle half-way between the polar angles whose cosines are μ_m and μ_{m-1} . Equations (147) and (148) are thus replaced by

$$k_q = \delta_q \int_{B_q} \int_{\phi_{l-1}}^{\phi_l} \vec{\Omega} \cdot \vec{n} d\phi dS \quad (149)$$

where $\delta_q = \mu_{m_q-1} - \mu_{m_q}$ and

$$\eta_{q,q'} = \frac{\delta_{q'}}{k_{q'}} \int_{B_{q'}} \int_{\alpha(P)}^{\omega(P)} \vec{\Omega} \cdot \vec{n} d\phi dS \quad (150)$$

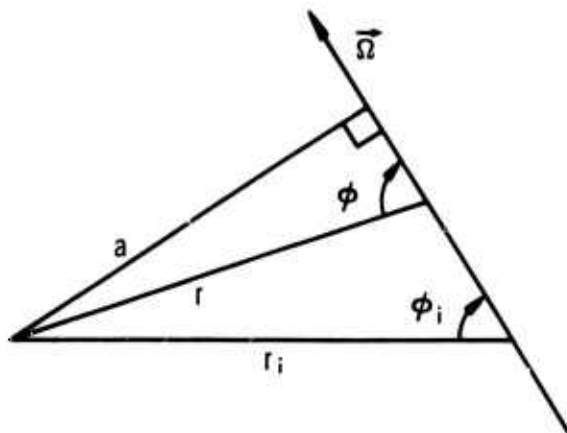
where $\left(\frac{\omega(P)}{\alpha(P)} \right)$ is the $\left(\begin{smallmatrix} \text{upper} \\ \text{lower} \end{smallmatrix} \right)$ limit on the azimuths of rays through P that have axial projection $\mu_{m-\frac{1}{2}}$ and belong to both fans.

If F_q is the fan associated with cone $C_{l,m}$ and a vertical zone boundary of radius r and axial height h , then

$$k_q = 2\pi r h \delta_q (1 - \mu_{m-\frac{1}{2}}^2)^{\frac{1}{2}} (\sin \phi_l - \sin \phi_{l-1}) \quad (151)$$

On the other hand, let the zone boundary associated with F_q be a horizontal one with inner radius r_{i-1} and outer radius r_i . As a point passes through the system along a ray with direction $\vec{\Omega}$, the distance r of the point from

the axis, the azimuth ϕ of $\vec{\Omega}$, and the distance a of the ray from the axis satisfy the relation $r \sin \phi = a$. (See Fig. 5.) It follows that the azimuth ϕ of a ray of F_q will range between $\arcsin (r_i \sin \phi_{l-1})/r$ and $\arcsin (r_i \sin \phi_l)/r \wedge + 1$ at a distance r from the axis. (Here $a \wedge b$ is

Figure 5. The Dependence of ϕ on r

the lesser and $a \vee b$ is the greater of the two quantities a and b .) It is assumed that $r \geq r_i \sin \phi_{l-1}$, for otherwise no ray of F_q will come within a distance r of the axis. The appropriate form for k_q is thus

$$k_q = \delta_q \int_{r_{i-1} \vee r_i \sin \phi_{l-1}}^{r_i} 2 \pi r \mu_{m-\frac{1}{2}} \left[\arcsin \left(\frac{r_i \sin \phi_l}{r} \vee 1 \right) - \arcsin \frac{r_i \sin \phi_{l-1}}{r} \right] dr \quad (152)$$

assuming that $\phi_l \leq (1/2) \pi$. When $\phi_l > (1/2) \pi$, the equation is still correct provided that the appropriate branch of the arcsine is used. In order to write a formula for k_q in terms of elementary transcendental functions, let $a_l = r_i \sin \phi_l$ and

$$f_l(r) = \int_0^r r' \arcsin \frac{a_l}{r'} dr' = \frac{r^2}{2} \arcsin \frac{a_l}{r} + \frac{a_l}{2} \sqrt{r^2 - a_l^2}$$

Then

$$k_q = 2\pi\mu_{m-\frac{1}{2}} \delta_q [f_l(r_i) - f_l(r_{i-1}) - f_{l-1}(r_i) + f_{l-1}(r_{i-1})]$$

when $r_{i-1} > a_l$ and

$$k_q = 2\pi\mu_{m-\frac{1}{2}} \delta_q [f_l(r_i) - \frac{\pi}{4} (r_{i-1} \vee a_{l-1})^2 - f_{l-1}(r_i) + f_{l-1}(r_{i-1} \vee a_{l-1})]$$

when $r_{i-1} \leq a_l$.

As is true of many geometrical problems, the calculation of the coefficients η requires the consideration of many special cases. Let $F_{q'}$ and F_q be a pair of fans for which $Z_{q'}^i = Z_q^0 = Z_{i,j}$. Note that $\eta_{q,q'} = 0$ unless $m_q = m_{q'}$, because μ does not change as a point travels along a ray. Consequently, let m be given and assume that $m_q = m = m_{q'}$. Also, denote $\mu_{m-\frac{1}{2}}$ simply by μ and $\phi_{l,m}$ by ϕ_l . Further, let $p = \phi_l - \phi_{l-1}$ so that $\phi_l = lp$ and $l_{\max} p = \pi$, where l_{\max} denotes the maximum value of l . Denote l_q simply by l and $l_{q'}$ by l' . Let $h = z_j - z_{j-1}$.

Consider first those cases in which the rays of fan $F_{q'}$ enter zone $Z_{i,j}$ through its outer curved boundary $BR_{i,j}$. Then the integral for the view factor $\eta_{q,q'}$ has the form

$$\eta_{q,q'} = \frac{2\pi r_i (1 - \mu^2)^{\frac{1}{2}} \delta_q}{k_{q'}} \int_0^h \int_{\alpha(\zeta)}^{\omega(\zeta)} \cos \phi d\phi d\zeta \quad (153)$$

where ζ refers to the distance between the point P on $BR_{i,j}$ and the top edge of $BR_{i,j}$ and $\begin{pmatrix} \omega(\zeta) \\ \alpha(\zeta) \end{pmatrix}$ is the $\begin{pmatrix} \text{upper} \\ \text{lower} \end{pmatrix}$ bound on the azimuths of rays through P belonging to both fans. Thus, one obtains

$$\eta_{q,q'} = \frac{2 \pi r_i (1 - \mu^2)^{\frac{1}{2}} \delta_{q'}}{k_{q'}} \int_0^h \left[\sin \omega(\zeta) - \sin \alpha(\zeta) \right] d\zeta \quad (154)$$

Outside-to-Inside. Let the rays common to fans F_q and $F_{q'}$ enter $Z_{i,j}$ through the outside boundary and leave through the inside boundary. Rays belonging to fan F_q leave $Z_{i,j}$ at $r = r_{i-1}$ with azimuths in the range $(l-1)p \leq \phi \leq lp$, which means that at $r = r_i$, where they enter, their azimuths fall in the range

$$\arcsin \frac{r_{i-1} \sin(l-1)p}{r_i} \leq \phi \leq \arcsin \frac{r_{i-1} \sin lp}{r_i}$$

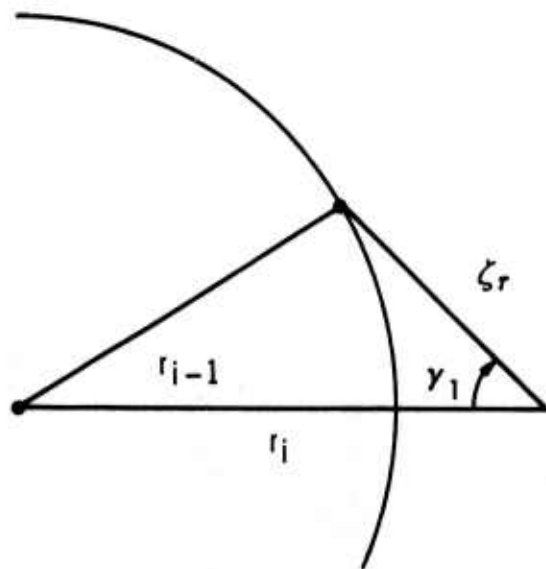
(See Fig. 6.) Consequently, the azimuths of rays common to F_q and $F_{q'}$ fall in the range

$$(l' - 1)p \vee \arcsin \frac{r_{i-1} \sin(l-1)p}{r_i} \leq \phi \leq l'p \wedge \arcsin \frac{r_{i-1} \sin lp}{r_i}$$

It remains to determine which rays in $F_{q'}$ leave zone $Z_{i,j}$ through the inside boundary. Assume that $\mu > 0$ so that the rays have an upward component, and let

$$\tau = \tan \arccos \mu = (1 - \mu^2)^{\frac{1}{2}} / \mu$$

which is the tangent of the angle between the ray and the axis. Then $\tau \zeta$ will be the horizontal distance that a point goes when it rises by an amount ζ along a ray having an axial projection of μ .

Figure 6. Determination of γ_1

Therefore, a ray of fan $F_{q'}$ that enters zone $Z_{i,j}$ through $BR_{i,j}$ at a point ζ below the top will leave through the inside boundary $r = r_{i-1}$ provided that

$$\cos \phi \geq \frac{r_i^2 + \zeta^2 \tau^2 - r_{i-1}^2}{2 r_i \zeta \tau}$$

where ϕ is its azimuth at the outside boundary $r = r_i$, assuming that

$$(r_i - r_{i-1})^2 < \zeta^2 \tau^2 < r_i^2 - r_{i-1}^2$$

If $\zeta \tau < r_i - r_{i-1}$, then the ray will leave through the top or outside regardless of its incoming azimuth; and if $\zeta \tau > (r_i^2 - r_{i-1}^2)^{\frac{1}{2}}$, then it will leave through the inside whenever $\sin \phi < r_{i-1}/r_i$. All this can be summarized by saying that the ray through $BR_{i,j}$ at a point ζ below the top will leave through the inside only if $\phi < \gamma_1(\zeta)$, where

$$\gamma_1(\zeta) = \begin{cases} 0, & \text{if } \zeta \tau < r_i - r_{i-1} \\ \arcsin r_{i-1}/r_i, & \text{if } \zeta^2 \tau^2 > r_i^2 - r_{i-1}^2 \\ \arccos \frac{r_i^2 + \zeta^2 \tau^2 - r_{i-1}^2}{2 r_i \zeta \tau}, & \text{otherwise} \end{cases} \quad (155)$$

It follows that

$$\eta_{q,q'} = \frac{2 \pi r_i (1 - \mu^2)^{\frac{1}{2}}}{k_{q'}} \int_0^h [\sin \omega(\zeta) - \sin \alpha(\zeta)] \vee 0 \, d\zeta \quad (156)$$

where

$$\begin{aligned} \omega(\zeta) &= \ell' p \wedge \gamma_1(\zeta) \wedge \arcsin \frac{r_{i-1} \sin \ell p}{r_i} \\ \alpha(\zeta) &= (\ell' - 1) p \vee \arcsin \frac{r_{i-1} \sin (\ell - 1)p}{r_i} \end{aligned} \quad (157)$$

The part, $\vee 0$, of Eq. (156) is required when $\omega(\zeta) < \alpha(\zeta)$, which will occur, for example, when $\gamma_1(\zeta) < (\ell' - 1)p$, which is the case for small ζ when $\ell' \geq 2$.

Outside-to-Top. Next, let the rays common to fans F_q and $F_{q'}$ enter $Z_{i,j}$ through the outside boundary and leave through the top. This case may be divided into two subcases, depending on whether the rays are traveling inward ($\ell p \leq \pi/2$) or outward ($\ell p > \pi/2$) when they emerge. In either subcase, the azimuths of rays in fan F_q will lie in the range $(\ell - 1)p \leq \phi \leq \ell p$ at $r = r_i$; but in the latter subcase, this will be the situation as they pass out of the cylinder $r = r_i$. It is seen from figure 7 that the incoming and outgoing azimuths are supplementary and hence that

the range of azimuths of rays in F_q where they enter the cylinder $r = r_i$ is, for $\ell p > \frac{1}{2}\pi$,

$$(\ell_{\max} - \ell)p \leq \phi \leq [\ell_{\max} - (\ell - 1)]p$$

Therefore, $\eta_{q,q'} = 0$ unless $\ell' = \ell$ or $\ell' = \ell_{\max} - (\ell - 1)$. Equation (156) is still valid for $\eta_{q,q'}$, but the formulas for $\omega(\zeta)$ and $\alpha(\zeta)$ must be changed. In case $\ell' = \ell$, $\alpha(\zeta) \geq \gamma_1(\zeta)$ (Eq. (155)) to assure that the rays will not emerge through the inside surface and a new constraint on $\omega(\zeta)$ is required to assure that the rays are still going inward when they emerge. In figure 8, it is seen that if a ray has azimuth ϕ at $r = r_i$, then it must go a horizontal distance $d = r_i \cos \phi$ before beginning to go outward. But the horizontal distance that a ray in fan $F_{q'}$ travels in changing its altitude by ζ is $\tau\zeta$, where $\tau = (1 - \mu^2)^{1/2}/\mu$. Therefore,

$$\beta(\zeta) = \arccos \left(\frac{\tau\zeta}{r_i} \wedge 1 \right) \quad (158)$$

is the upper limit of the azimuth of rays entering zone $Z_{i,j}$ in fan $F_{q'}$ at a point lying ζ below the top boundary that are still going inward when they reach the top. The $\wedge 1$ part of the formula forces $\beta(\zeta) = 0$ when $\tau\zeta \geq r_i$, which is to be desired since all rays entering farther from the top than r_i/ζ are outward-going when they reach the top. Thus, when $\ell' = \ell$,

$$\begin{aligned} \alpha(\zeta) &= (\ell' - 1)p \vee \gamma_1(\zeta) \\ \omega(\zeta) &= \ell'p \wedge \beta(\zeta) \end{aligned} \quad (159)$$

When $\ell' = \ell_{\max} - \ell + 1$, it is required that the rays emerge going outward through the top. $\alpha(\zeta) \geq \beta(\zeta)$ will assure that they are going outward. A new restriction is required to eliminate rays that leave through the outside boundary. Recalling that $\tau\zeta$ is the horizontal distance traveled in rising to the top, it is seen in figure 9 that

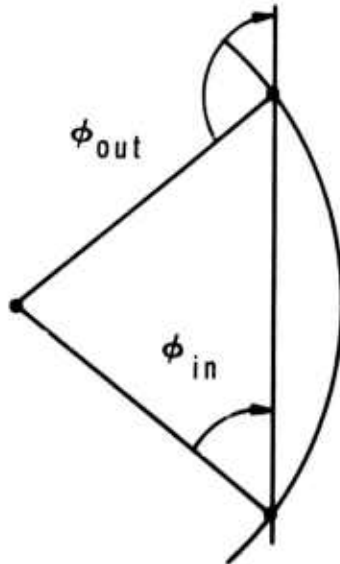


Figure 7. Relation of ϕ_{in} to ϕ_{out}

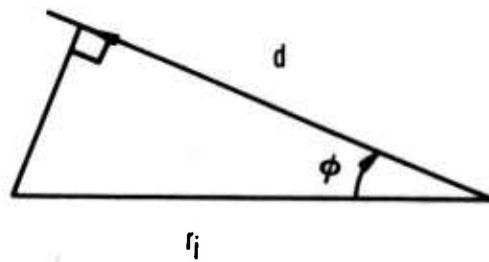


Figure 8. Distance to the Point Closest to the Axis

$$\gamma_2(\zeta) = \arccos \left(\frac{\tau \zeta}{2r_i} \wedge 1 \right) \quad (160)$$

is the limiting azimuth. Again, clipping the cosine at 1 is appropriate because any ray that enters the cylinder $r = r_i$ with axial projection μ will leave before it goes an axial distance of $2r_i/\tau$. Thus, when $\ell' = \ell_{\max} - \ell + 1$,

$$\begin{aligned} \alpha(\zeta) &= (\ell' - 1) \vee \gamma_1(\zeta) \vee \beta(\zeta) \\ \omega(\zeta) &= \ell' \vee \gamma_2(\zeta) \end{aligned} \quad (161)$$

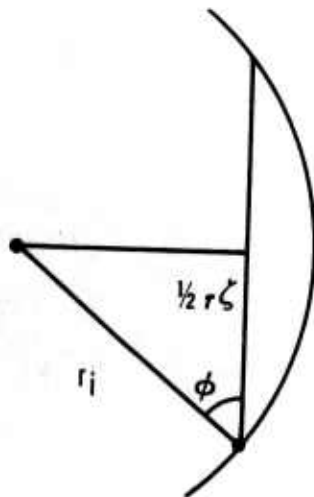


Figure 9. Azimuth of a Ray that Strikes the Upper Outside Corner

Outside-to-Outside. As in the second subcase of outside-to-top, it can be inferred that if $F_{q'}$ consists of rays entering zone $Z_{i,j}$ through the boundary $r = r_i$ and F_q consists of rays leaving through the same boundary, then $\eta_{q,q'} = 0$ unless $\ell' = \ell_{\max} - \ell + 1$. The limit on $\alpha(\zeta)$ is simply to assure that the rays emerge through the outside boundary before reaching the top or inside. Thus,

$$\alpha(\zeta) = (\ell' - 1) p \vee \gamma_1(\zeta) \vee \gamma_2(\zeta)$$

$$\omega(\zeta) = \ell' p.$$

(162)

Inside-to-Outside. As is to be expected, this case bears a strong resemblance to the outside-to-inside case. One difference is that the integrals in the numerator and denominator of Eq. (154) are now negative, although $\eta_{q,q'}$ remains positive. Other differences will appear as the discussion progresses.

Let the rays common to fans F_q and $F_{q'}$ enter zone $Z_{i,j}$ through the inside boundary and leave through the outside boundary. As before, the boundary associated with $F_{q'}$ is the one through which the rays are entering; i.e., the inside boundary associated with F_q is the outside boundary. The rays in fan F_q leave the cylinder $r = r_i$ with azimuths in the range $(\ell - 1) p \leq \phi \leq \ell p$. It is seen from figure 10 that a ray having azimuth ϕ at $r = r_i$ has azimuth $\pi - \arcsin(r_i \sin \phi / r_{i-1})$ at $r = r_{i-1}$.

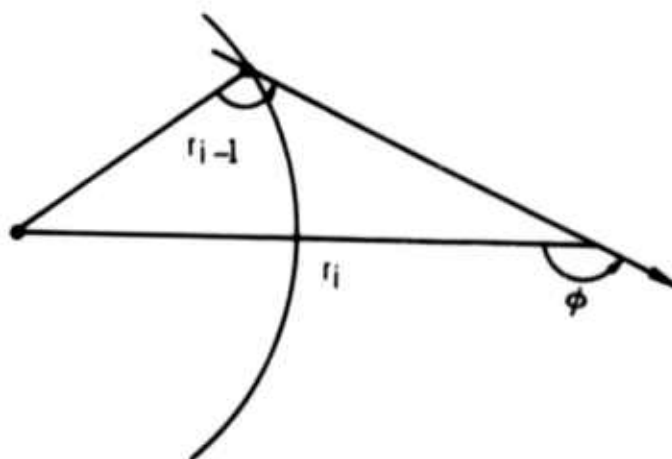


Figure 10. Azimuth of Outgoing Ray

provided that it gets that close to the axis (which is the case if, and only if, $r_i \sin \phi \leq r_{i-1}$). Thus, the azimuths of rays of F_q that come within r_{i-1} of the axis lie at that point in the range

$$\pi - \arcsin \left(\frac{r_i \sin (\ell - 1) p}{r_{i-1}} \wedge 1 \right) \leq \phi \leq \pi - \arcsin \left(\frac{r_i \sin \ell p}{r_i} \wedge 1 \right)$$

As for the condition that the rays hit the outside boundary before they reach the top, assume that $(r_i - r_{i-1})^2 \leq \tau^2 \zeta^2 \leq r_i^2 - r_{i-1}^2$. Then a ray of azimuth

$$\phi = \arccos \frac{r_{i-1}^2 + \tau^2 \zeta^2 - r_i^2}{2 r_{i-1} \tau \zeta}$$

will rise by an amount ζ in going from $r = r_{i-1}$ to $r = r_i$. Consequently, if $\gamma_2(\zeta)$ is redefined as

$$\gamma_3(\zeta) = \begin{cases} \pi & \text{if } \tau \zeta \leq r_i - r_{i-1} \\ \pi/2 & \text{if } \tau^2 \zeta^2 > r_i^2 - r_{i-1}^2 \\ \arccos \frac{r_{i-1}^2 + \tau^2 \zeta^2 - r_i^2}{2 r_{i-1} \tau \zeta} & \text{otherwise,} \end{cases} \quad (163)$$

then it can be stated that a ray with azimuth ϕ at $r = r_{i-1}$ will reach $r = r_i$ before it rises by an amount ζ if, and only if, $\phi > \gamma_3(\zeta)$. Thus, Eq. (156) is still a valid equation for $\eta_{q,q'}$ provided that

$$\begin{aligned} \alpha(\zeta) &= (\ell' - 1) p \vee \left[\pi - \arcsin \left(\frac{r_i \sin(\ell - 1)p}{r_{i-1}} \vee 1 \right) \right] \vee \gamma_3(\zeta) \\ \omega(\zeta) &= \ell' p \wedge \left[\pi - \arcsin \left(\frac{r_i \sin \ell p}{r_{i-1}} \wedge 1 \right) \right] \end{aligned} \quad (164)$$

Inside-to-Top. This case is exactly like the inside-to-outside case except that $\gamma_3(\zeta)$ is now an upper limit:

$$\begin{aligned}\alpha(\zeta) &= (\ell' - 1)p \vee \left[\pi - \arcsin \left(\frac{r_i \sin(\ell - 1)p}{r_{i-1}} \wedge 1 \right) \right] \\ \omega(\zeta) &= \ell' p \wedge \left[\pi - \arcsin \left(\frac{r_i \sin \ell p}{r_{i-1}} \wedge 1 \right) \right] \wedge \gamma_3(\zeta)\end{aligned}\quad (165)$$

Next, consider those cases in which $F_{q'}$ is a fan associated with the bottom of zone $Z_{i,j}$ (assuming, as usual, that $\mu > 0$). In such cases the formula (150) for $\eta_{q,q'}$ becomes

$$\eta_{q,q'} = \frac{2\pi\mu}{k_{q'}} \delta_{q'} \int_{r_{i-1}}^{r_i} ([\omega(r) - \alpha(r)] \vee 0) r \, dr \quad (166)$$

where $\omega(r)$ and $\alpha(r)$ are the upper and lower limits on the azimuths of rays common to fans F_q and $F_{q'}$ that enter the bottom of $Z_{i,j}$ at radius r .

Bottom-to-Inside. Let the rays common to fans F_q and $F_{q'}$ enter zone $Z_{i,j}$ through the bottom and leave through the inside boundary. The azimuth ϕ of a ray in fan $F_{q'}$ lies in the range $(\ell' - 1)p \leq \phi \leq \ell'p$ at $r = r_i$ and hence in the range

$$\arcsin \left(\frac{r_i \sin(\ell' - 1)p}{r} \wedge 1 \right) \leq \phi \leq \arcsin \left(\frac{r_i \sin \ell' p}{r} \wedge 1 \right)$$

at r . The azimuth ϕ of a ray in fan F_q lies in the range $(\ell - 1)p \leq \phi \leq \ell p$ at $r = r_{i-1}$ and hence in the range

$$\arcsin \frac{r_{i-1} \sin(\ell - 1)p}{r} \leq \phi \leq \arcsin \frac{r_{i-1} \sin \ell p}{r}$$

at r . It remains to determine what limit must be placed on the azimuth of rays entering through the bottom to assure their leaving through the inside boundary. In rising by an amount h , a ray goes a horizontal distance of $h\tau$.

Thus, if $\gamma_4(r)$ is defined as

$$\gamma_4(r) = \begin{cases} 0 & \text{if } r > r_{i-1} + h\tau \\ \arcsin r_{i-1}/r & \text{if } r^2 < r_{i-1}^2 + h^2\tau^2 \\ \arccos \frac{r^2 + h^2\tau^2 - r_{i-1}^2}{2rh\tau} & \text{otherwise} \end{cases} \quad (167)$$

then rays entering the bottom at radius r will reach the inside before the top if, and only if, they enter with azimuth less than $\gamma_4(r)$. Thus, Eq. (166) may be used for $\eta_{q,q'}$, provided that

$$\alpha(r) = \arcsin \left(\frac{(r_i \sin(\ell' - 1)p) \vee (r_{i-1} \sin(\ell - 1)p)}{r} \wedge 1 \right) \quad (168)$$

$$\omega(r) = \arcsin \left(\frac{(r_i \sin \ell' p) \vee (r_{i-1} \sin \ell p)}{r} \wedge 1 \right) \wedge \gamma_4(r)$$

Bottom-to-Top. Let $F_{q'}$ still be associated with the bottom of zone $Z_{i,j'}$, but let F_q be associated with the top. As in the outside-to-top case, $\eta_{q,q'} = 0$ unless $\ell' = \ell$ or $\ell + \ell' = \ell_{\max} + 1$, but now $\ell' = \ell$ splits into two subcases: $\ell' = \ell \leq (1/2)\ell_{\max}$ and $\ell' = \ell > (1/2)\ell_{\max}$. Also, it will be necessary to redefine β and γ . Let

$$\beta_1(r) = \arccos \left(\frac{\tau h}{r} \wedge 1 \right) \quad (169)$$

so that rays having azimuth $\phi < \beta_1(r)$ at r will still be going inward after rising by an amount h . And let

$$\gamma_5(r) = \begin{cases} 0 & \text{if } r < h\tau - r_i \\ \pi & \text{if } r < r_i - h\tau \\ \arccos \frac{r^2 + h^2\tau^2 - r_i^2}{2rh\tau} & \text{otherwise} \end{cases} \quad (170)$$

so that rays having azimuth $\phi < \gamma_5(r)$ at r will rise by an amount h before emerging from the cylinder $r = r_i$.

Consider first the subcase $\ell' = \ell \leq (1/2)\ell_{\max}$, which is the one where the rays are still going inward as they leave through the top of the zone. Then

$$\begin{aligned} \alpha(r) &= \arcsin\left(\frac{r_i \sin(\ell' - 1)p}{r} \wedge 1\right) \vee \gamma_4(r) \\ \omega(r) &= \arcsin\left(\frac{r_i \sin \ell' p}{r} \wedge 1\right) \wedge \beta_1(r) \end{aligned} \quad (171)$$

Next, if $\ell' \leq (1/2)\ell_{\max}$ and $\ell = \ell_{\max} - \ell' + 1 (> (1/2)\ell_{\max})$, then

$$\begin{aligned} \alpha(r) &= \arcsin\left(\frac{r_i \sin(\ell' - 1)p}{r} \wedge 1\right) \vee \beta_1(r) \vee \gamma_4(r) \\ \omega(r) &= \arcsin\left(\frac{r_i \sin \ell' p}{r} \wedge 1\right) \wedge \gamma_5(r) \end{aligned} \quad (172)$$

and, finally, if $\ell' = \ell > (1/2)\ell_{\max}$, then

$$\begin{aligned} \alpha(r) &= \pi - \arcsin\left(\frac{r_i \sin(\ell - 1)p}{r} \wedge 1\right) \\ \omega(r) &= \left[\pi - \arcsin\left(\frac{r_i \sin \ell' p}{r} \wedge 1\right) \right] \wedge \gamma_5(r) \end{aligned} \quad (173)$$

Bottom-to-Outside. Let $F_{q'}$ still be a fan associated with the bottom of zone $Z_{i,j}$, but let F_q be associated with the outside boundary. Then $\eta_{q,q'} = 0$ unless $\ell' = \ell$ or $\ell' = \ell_{\max} - \ell + 1$. Since the rays of F_q are assumed to be leaving $Z_{i,j}$ at the outside boundary, $\ell > (1/2)\ell_{\max}$. Thus, for $\ell' = \ell_{\max} - \ell + 1$,

$$\alpha(r) = \arcsin\left(\frac{r_i \sin(\ell' - 1)p}{r} \wedge 1\right) \vee \gamma_5(r)$$

$$\omega(r) = \arcsin\left(\frac{r_i \sin \ell' p}{r} \wedge 1\right)$$
(174)

and for $\ell' = \ell$,

$$\alpha(r) = \left[\pi - \arcsin\left(\frac{r_i \sin(\ell' - 1)p}{r} \wedge 1\right) \right] \vee \gamma_5(r)$$

$$\omega(r) = \pi - \arcsin\left(\frac{r_i \sin \ell' p}{r} \wedge 1\right)$$
(175)

ESTIMATION OF THE LENGTH

The factor t_q occurring in Eq. (146) is a quantity associated with fan F_q having the dimensions of length. It represents something like the mean length of the intersections of rays of fan F_q with the zone it leaves. The procedure used in SUBROUTINE DRAW of SHORT2 is to choose a representative ray of fan F_q and set t_q equal to the length of its intersection with the zone. In order to specify precisely what is done, define the mean direction of cone $C_{\ell,m}$ to be one with axial projection $\mu = \mu_{m-\frac{1}{2}}$ and azimuth $\phi = (\phi_{\ell-1,m} + \phi_{\ell,m})/2$ and the mid-circle of boundary $BA_{i,j}$ to be the circle $z = z_j$, $r = (1/2)(r_{i-1} + r_i)$ and the mid-circle of boundary $BR_{i,j}$ to be the circle $z = (z_{j-1} + z_j)/2$, $r = r_i$. In these terms, if F_q is associated with boundary B_q and cone C_q , then a representative ray R_q of F_q is a ray passing through the mid-circle of B_q in the mean direction

of cone C_q , and t_q is the length of the intersection of R_q with Z_q^0 . The value of t_q is well-defined, because all representative rays of F_q are equivalent with respect to the symmetry of the system, and consequently they intersect Z_q^0 in segments of equal length.

THE SOURCE TERM

The strength of the source at the center of a HECTIC zone is calculated from zone-centered temperatures using Eq. (131), and at other points it is obtained by a linear (in r and z) interpolation of the zone-centered data. The source term S_q in Eq. (146) is then obtained in the following way. Let R_q be the representative ray of F_q as defined above, and let P^i and P^o , respectively, be the points where R_q enters and leaves zone Z_q^0 . The distance between these points has already been given the designation t_q . Let S^α be the strength of the source at P^α , $\alpha = i$ or o . The contribution to the intensity of radiation at points along the segment of R_q between P^i and P^o due to the source in zone Z_q^0 is assumed to satisfy the transport equation,

$$\frac{dI}{ds} = -\sigma I + S^i + \frac{s(S^o - S^i)}{t_q}$$

where s is the distance along R_q from P^i toward P^o and $I(0) = 0$. It follows that

$$I(t_q) = \left[S^i - \frac{S^o - S^i}{\sigma t_q} \right] \frac{1 - e^{-\sigma t_q}}{\sigma} + \frac{S^o - S^i}{\sigma} \quad (176)$$

Finally, the source term is defined to be

$$S_q = I(t_q) k_q \quad (177)$$

where k_q is the number of rays in fan F_q as given in Eq. (147).

Another type of source to be considered is the radiant energy that flows into the system through its outer surface. As in LONG2, a boundary temperature θ_{out} is assigned to each interval of the outer boundary by input cards. On the assumption that the incoming energy is isotropic black-body radiation, the energy flow assigned to an incoming fan F_q is

$$J_q = \frac{ac}{4\pi} \theta_{out}^4 \left[P\left(\frac{h\nu_2}{\theta_{out}}\right) - P\left(\frac{h\nu_1}{\theta_{out}}\right) \right] k_q \quad (178)$$

which is the SHORT2 version of Eqs. (133) and (134). To see the analogy fully, observe that for any outer boundary interval B,

$$\sum_{q \in Q} k_q = \pi \alpha \quad (179)$$

where α is the area of B and $\{F_q\}_{q \in Q}$ is the set of fans F_q for which $B_q = B$ and which enter the system through B.

Serious difficulties have been encountered when the treatment of sources just described has been applied to overly heterogeneous systems. If, for example, a hot opaque zone lies next to a cold transparent one, then representative rays running through the cold zone after emerging from the hot one will, according to Eq. (176) pick up far more intensity than they should because of the large value of S^i . In LONG2, the problem has been relieved to some extent by setting

$$I(t^q) = S^c \frac{1 - e^{-\sigma t^q}}{\sigma} \quad (180)$$

where S^c is the central source strength whenever $S^i \wedge S^o \vee S^c < f(S^i \vee S^o \vee S^c)$, where f is input (usually, $f = 2$).

It must be emphasized that the problem of interpolating a source distribution on the basis of data supplied on a coarse mesh will never be completely solved. But in view of the expense of performing two-dimensional transport calculations, there is a strong economic incentive to find partial solutions of various types.

SOME SPECIAL SOLUTIONS

Two cases of importance are easy to solve with the view-factor method. The first is the case of an evacuated system with uniform isotropic incident radiation. In this case, $S = \sigma = 0$ in all zones, and Eq. (176) becomes indeterminate; but the correct assignment is, of course, $S_q = 0$. Let the energy flow in the incoming fans F_q be

$$J_q = I^0 k_q$$

By virtue of the relation expressed in Eq. (136), $J_q = I^0 k_q$ for all q is the solution to Eq. (146).

Next, suppose that the system is filled with material of constant absorption coefficient and constant source strength $S = \sigma I^0$. If, in addition, radiant energy is flowing into the system with isotropic intensity I^0 , then $J_q = I^0 k_q$ again satisfies Eqs. (146), (176), and (177) with $I(t_q) = I^0 (1 - e^{-\sigma t_q})$.

Another important special case, but one for which there exist no trivial solutions, is the case of the thick zone. To say a zone is thick implies that $\sigma t_q \gg 1$ and

$$I(t_q) \approx \frac{S^0}{\sigma} - \frac{S^0 - S^i}{\sigma^2 t_q}$$

Suppose, for example, that $S = S_0 + z S_1$ so that $S^0 - S^i = \Delta z S_1$ and $t_q = \Delta z / \mu_q$, where Δz is the difference in altitude between P^0 and P^i . Thus,

$$I(t_q) \approx \frac{S^0}{\sigma} - \frac{\mu_q S_1}{\sigma}$$

and

$$J_q \approx \left(\frac{S^0}{\sigma} - \frac{\mu S_1}{\sigma^2} \right) k_q$$

Next, let B be some horizontal zone boundary, and let Q be the set of fan indices q for which $\mu_q > 0$ and $B_q = B$. Then the rate of transfer of radiant energy in the upward direction through B will be, according to the view-factor method,

$$\begin{aligned}\dot{E}_B^{\text{up}} &= \sum_{q \in Q} J_q \\ &= \frac{S^0}{\sigma} \sum_Q k_q - \frac{S_1}{\sigma} \sum_Q \mu_q k_q\end{aligned}$$

To perform the sums, first sum over azimuth. Simply setting $\phi_\ell = \pi/2$ and $\phi_{\ell-1} = 0$ in Eq. (152) gives 1/4 of the total, i.e.,

$$k_q = \frac{\pi \delta_q}{2} (\pi r_i^2 - \pi r_{i-1}^2) \mu_q$$

so the sum over azimuth is

$$\sum_\ell k_q = 2 \pi \delta_q \alpha \mu_q$$

where $\alpha = \pi(r_i^2 - r_{i-1}^2)$, the area of B. The sums over axial projection are then

$$\sum_m \mu_{m-\frac{1}{2}} (\mu_{m-1} - \mu_m) \approx \int_0^1 \mu d\mu = \frac{1}{2}$$

$$\sum_m \mu_{m-\frac{1}{2}}^2 (\mu_{m-1} - \mu_m) \approx \int_0^1 \mu^2 d\mu = \frac{1}{3}$$

which finally leads to

$$\dot{E}_B^{\text{up}} \approx \frac{S^0}{\sigma} \pi \alpha - \frac{S_1}{2} \frac{2\pi\alpha}{3}$$

Performing the same calculation for $\mu < 0$ only changes the sign of the first term. Thus,

$$\dot{E}_B = \dot{E}_B^{\text{up}} + \dot{E}_B^{\text{down}} = -\frac{4\pi\alpha}{3} \frac{S_1}{\sigma^2}$$

or, since

$$\frac{S_1}{\sigma} = \frac{1}{\sigma} \frac{dS}{dz} = \frac{dS}{d\tau}$$

where τ is optical depth, one finally arrives at the diffusion approximation

$$\dot{E}_B = -\frac{4\pi\alpha}{3\sigma} \frac{dS}{d\tau}$$

A similar result is obtained when the gradient of the source is radial, but there are additional complications in that case caused by the curvature of vertical boundaries.

METHOD OF SOLUTION

Solving the transport equation in the view-factor approximation is, on the face of it, a matter of solving a large linear system of simultaneous equations (Eq. (146)). It turns out, however, that it is possible to arrange the equations in a sequence that permits each to be solved in terms of the solution of the preceding ones.

Recall first that $\eta_{q,q'} = 0$ unless $m_q = m_{q'}$. Hence, it is possible to solve subsystems of equations corresponding to the various values of m_q independently of each other. Consider, then, the problem of solving the system represented by Eq. (146) for all q such that $m_q = m$. Suppose, for example, that the sign of $\mu_{m-\frac{1}{2}}$ is positive. The technique is to treat a layer of zones at a time, a layer of zones being $Z_{i,j}$ for $1 \leq i \leq i_{\text{max}}$ and j fixed, starting with the bottom layer ($j = 1$) and proceeding upward. Within a layer, one first calculates the J 's of the incoming fans ($\phi_l \leq \pi/2$), starting at the outside and proceeding inward, and then the outgoing fan ($\phi_l > \pi/2$), starting at the center ($i = 1$) and proceeding outward. To see why this

sequence has all the desired properties, observe first that when $j = 1$, all the J 's for fans entering through the bottom are specified as boundary sources*, which is also true of incoming fans at $i = i_{\max}$. On the inward sweep of the bottom layer, the J 's of incoming fans entering through the bottom and outside of any particular zone $Z_{i,1}$ are all that are needed to calculate the incoming fans leaving through the inside and top. Since the incoming fans leaving $Z_{i,1}$ through $BL_{i,1}$ are the incoming fans entering $Z_{i-1,1}$ through $BR_{i-1,1}$ ($= BL_{i,1}$), the next step is ready to be taken, and the validity of the inward sweep is established. To calculate J 's of the outgoing fans leaving $Z_{i,j}$ through the outside and top, it suffices to have the J 's of the outgoing fans entering through the inside and bottom together with the incoming fans that were obtained on the inward sweep. Thus, an outward sweep does permit the calculation of the outgoing fans. The process produces the J 's of both incoming and outgoing fans through the top boundaries of the layer which provides the bottom boundary condition for the layer above. Thus, the entire upward sweep is carried out. When the upward sweep is done, a downward sweep is performed with the same set of η 's and t 's, but with the layers reversed for the μ which is just the negative of the one involved in the preceeding upward sweep.

The contributions to $\dot{E}_{i,j}$ (Eq. (138)) are accumulated as the J 's are computed. Some storage of J 's is required, but not much; it suffices to have room for all the J 's that enter or leave one layer for one value of μ . When the top boundary is a symmetry boundary, the appropriate boundary condition is $J_{l,m'} = J_{l,m}$, where $\mu_{m'} = -\mu_m$. When the bottom boundary is a symmetry boundary, a similar relationship subsists on the bottom, so the downward sweep must precede the upward one. A pair of symmetry planes, which would imply an axial periodicity, is not currently treated by the code.

* As discussed in the section entitled "The Source Term."

SECTION VI

FLOW OF CONTROL IN SUBROUTINE DRAW AND THE "DRAW SUBROUTINE SEQUENCE"

SUBROUTINE DRAW specifies the overall flow of control through the "DRAW subroutine sequence" (shown in Figs. 11 and 12) that initializes geometric quantities and creates the "DRAW data" for the "transport subroutine sequence" (shown in Fig. 18).

Upon entering SUBROUTINE DRAW the input to the "DRAW subroutine sequence" is read in on cards. This input consists of black-body system boundary temperatures, two tolerances that place restrictions on accuracy of the calculations in the "DRAW subroutine sequence" (TOLER and CHKSUM), and variables LMAX and MUMAX that specify the number of fans of direction to be constructed per hydro zone of a HECTIC problem.

The "DRAW subroutine sequence" creates, for each zone surface of the hydro mesh, a set of fans of direction, each specified by a polar angle $\theta_m(\arccos MU(M))$, an azimuthal angle $\phi_l(PHI(L))$, and a range $\Delta\phi(PP)$ of azimuthal angles that must be the same for all fans associated with all zone surfaces. Each zone surface has associated with it a set of fans, each set consisting of $LMAX*MUMAX*2$ fans. (For each polar angle θ_m of MUMAX polar angles, there are $LMAX*2$ azimuthal fans.) Because of cylindrical symmetry, however, only LMAX azimuthal fans are considered in the "DRAW subroutine sequence."

Given a set of fans for each zone face, the "DRAW subroutine sequence" performs a sweeping operation through the hydro mesh in a specified manner calculating "view factors," or " η 's," which couple

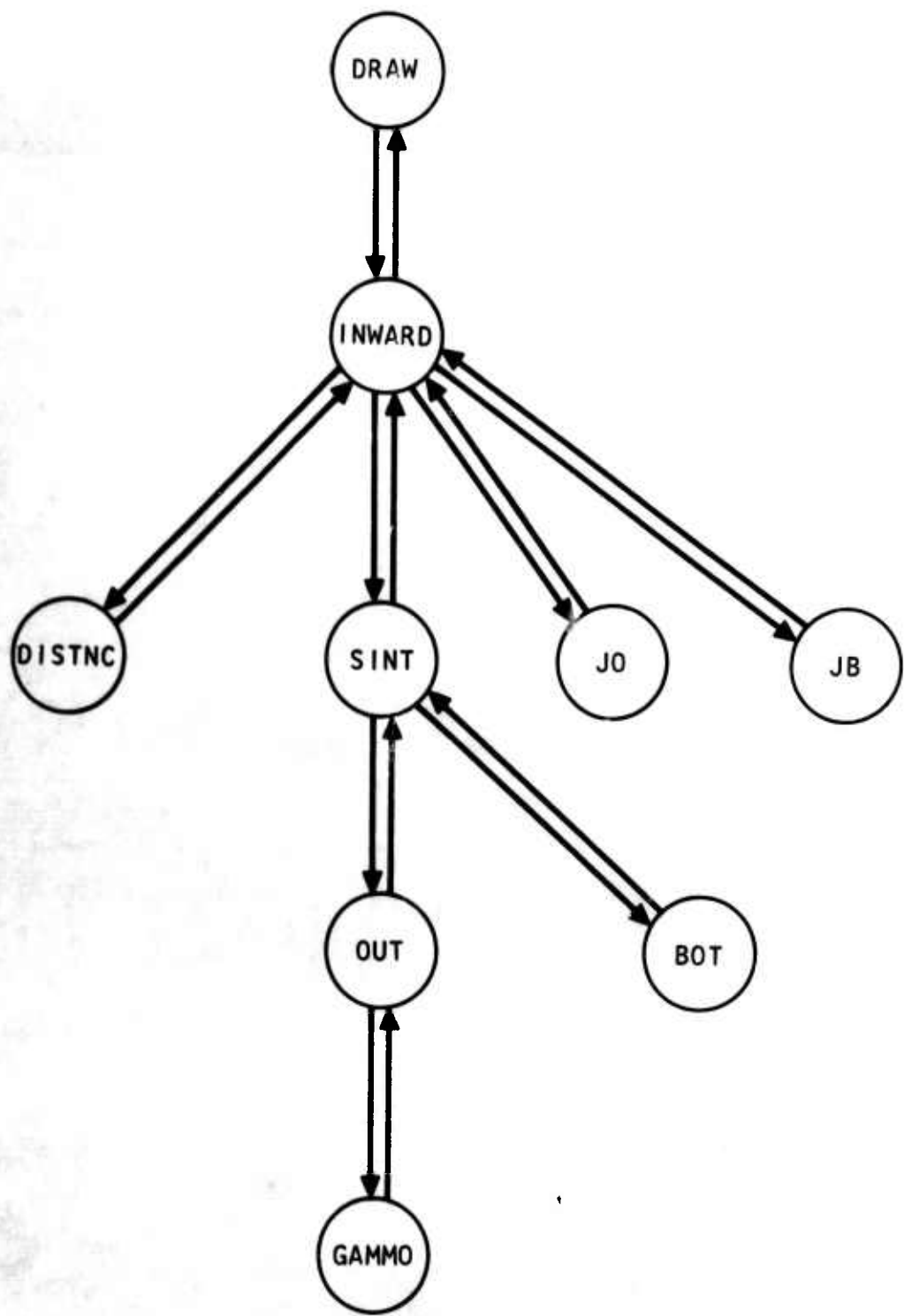


Figure 11. SUBROUTINES of the "DRAW Subroutine Sequence"
Employed in an Inward DRAW Sweep

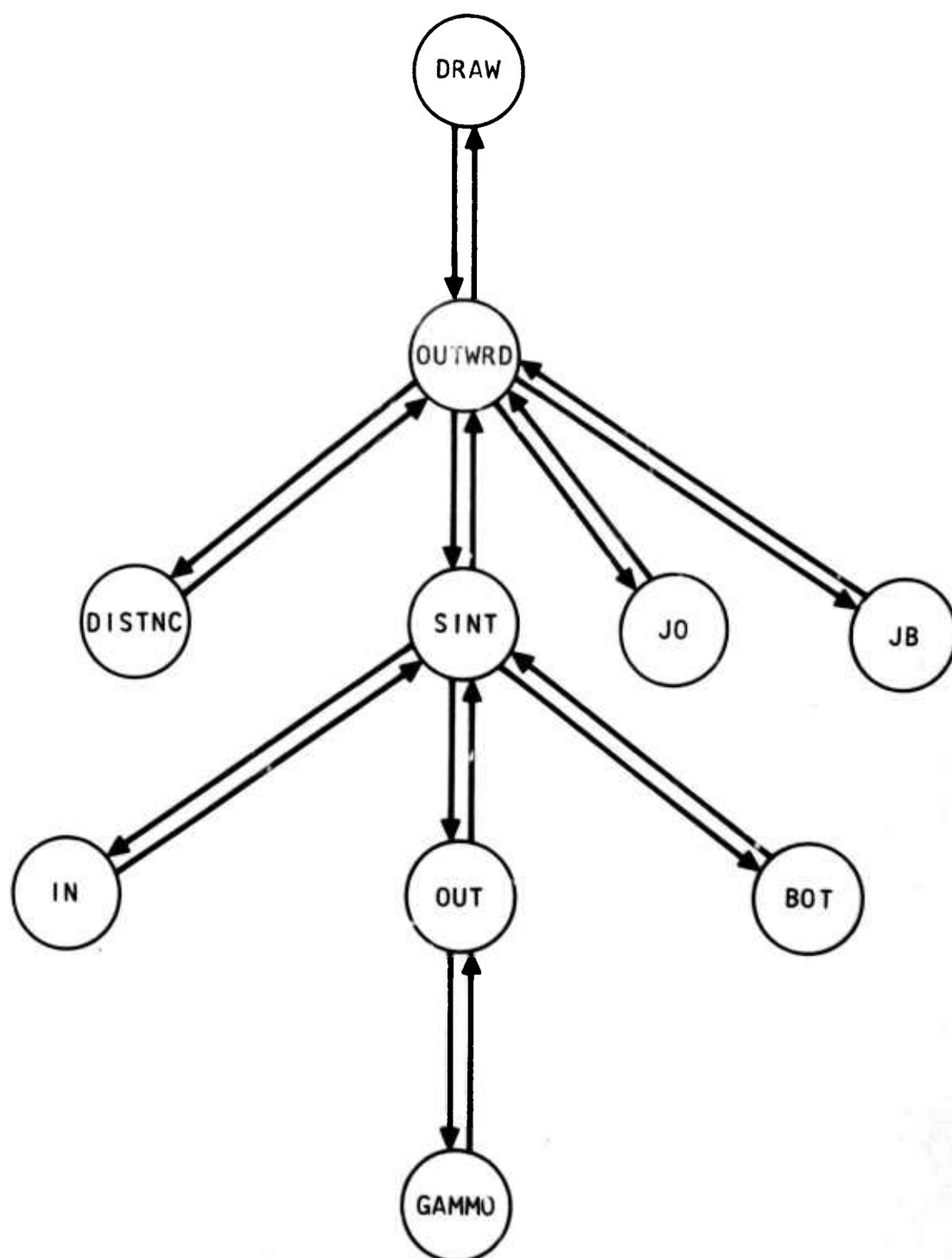


Figure 12. SUBROUTINES of the "DRAW Subroutine Sequence"
Employed in an Outward DRAW Sweep

each fan on each zone face with all the other fans on all the neighboring zone faces. Formally:

Given a particular hydro zone,

$\eta_{S_1 \rightarrow S_2}(M, L, LP)$ = the fraction of (energy entering the zone in time t through face S_1 in fan LP) which leaves the zone through face S_2 in fan L in the case of isotropic, homogeneous radiation $\sigma = S = 0$ with the intensity of the radiation field, $I_0 = 1$.

Notice the neighboring zone faces S_2 , are the faces of the zones which share face S_1 .

The order in which sweeping occurs in DRAW is as follows:

For each polar angle specified in M ($M=1, MUMAX/2$), sweeping is performed upward through the hydro mesh. For each zone horizontal layer, an inward and or outward sweep is performed. On the inward sweep, sweeping begins at the system outside boundary. (For an account of which η 's are calculated on the inward and outward sweeps, see SUBROUTINES INWARD and OUTWARD respectively). Only half the polar angles are considered, because

$$\eta_{S_1 \rightarrow S_2}(M, L, LP) = \eta_{S_1 \rightarrow S_2}(MUMAX+1-M, L, LP)$$

due to the symmetry restrictions placed on the polar angles θ_M (see DRAW Glossary).

Notice that a fan of a given zone face may be associated with any point on that face. For this reason each η calculation consists of an integration over the given zone face, and this integration is performed numerically by FUNCTION SINT under control of the relative error criterion TOLER, which is input in DRAW. It is desirable and reasonable to require that the η 's be calculated such that $\sum \eta(L, LP) = 1.0$ exactly, for each fan LP of a zone face summed over all fans L of neighboring zone faces. However, each η is calculated by an integration

with a certain error; and with the present version of the code the η 's do not sum exactly to 1.0. The sum is checked in DRAW (and partially in SUBROUTINE OUTWRD) for each fan of each zone face and a message is printed if the sums are not within CHKSUM (input) of 1.0. If the printing of this message occurs during execution of SUBROUTINE DRAW, it is an indication (assuming $\text{CHKSUM} \geq 10 * \text{TOLER}$) that there is a programming error in the calculation of the η 's. The η 's along with information specifying the characteristic ray path (see SUBROUTINE DISTNC) associated with each fan on each zone face and other information comprise the "DRAW data" which is output by SUBROUTINES DRAW, INWARD, and OUTWRD to a "DRAW data" file to be used by SUBROUTINE TRAN2 in each cycle of a HECTIC problem. Output of this data is accomplished by the filling and subsequent emptying of an output storage buffer (BUFF2) as the sweeping through the hydro mesh occurs, creating records of data on the file as the buffer is filled and emptied. The smallest unit of data allowed for one buffer-load is an inward or an outward sweep through one zone (one CALL executed to SUBROUTINES INWARD or OUTWRD). The end of an inward or outward sweep of one zone layer is always terminated by the creation of a record on the "DRAW data" file; thus, "DRAW data" of one complete sweep of a zone layer is never mixed with that of the following or preceding zone layer sweep (this ensures ease of handling this "DRAW data" in SUBROUTINE TRAN2).

It is intended that the "DRAW data" file be a magnetic tape. Since the data comprising this file is of a geometric nature and only depends on the HECTIC problem mesh and the variables LMAX, MUMAX, and TOLER, a complete HECTIC problem (comprising many cycles) may be solved with only one FORTRAN CALL to SUBROUTINE DRAW and subsequent CALLS to SUBROUTINE TRAN2 for each cycle of a HECTIC problem. If the "DRAW data" file is a magnetic tape, then SUBROUTINE DRAW need not be called again to regenerate data. (However, parameters transferred to TRAN2 by DRAW through COMMON must be initialized.)

Input to "DRAW SUBROUTINE SEQUENCE"

Medium	Variables
CARDS	LMAX
	MUMAX
	TOLER
	CHKSUM
	ATHETA(I), I=1, IMAX
	BTHETA(I), I=1, IMAX
	RTHETA(J), J=1, JMAX

Conditions That Cause Termination of Execution in "DRAW SUBROUTINE SEQUENCE"

The current version of the coding has only one error condition test. BUFF1 in SUBROUTINES INWARD and OUTWRD is not allowed to overflow--the problem will stop if this occurs (see SUBROUTINES INWARD or OUTWRD). There are, however, restrictions on the variables that are considered input to the "DRAW subroutine sequence" and restrictions on variables that are to be calculated by the sequence. (See SUBROUTINE DRAW Glossary.)

Results from DRAW SUBROUTINE SEQUENCE Transferred to the "TRANSPORT SUBROUTINE SEQUENCE"

<u>Medium</u>	<u>Variables</u>
COMMON	LMAX
	MUMAX
	ATHETA(I), I=1, IMAX
	BTHETA(I), I=1, IMAX
	RTHETA(J), J=1, JMAX

PP
 L2
 L2P1
 PHI(L), L=1, LMAX
 COSPHI(L), L=1, LMAX
 SINLP(L), L=1, LMAX
 SINP2
 M2
 MU(M), M=1, MUMAX
 SMU(M), M=1, MUMAX
 W(M), M=1, MUMAX
 NOBUFF
 NOO
 CH, CVIN, CVOUT
 DH, XH, YH
 DV, XV, YV
 All η 's

"DRAW data" FILE

Glossary of Variables used in SUBROUTINE DRAW

FORTTRAN Label	Report Label	Description
ATHETA(I)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (I, JMAX) (in eV) (input on cards) (used in TRAN2)
BUFF2	None	\equiv P. BUFF2(1) = NOO. Output storage buffer array for the "DRAW data" that is to be written on the "DRAW data" file. (The maximum number of words allowed is 1023, and each buffer-load must contain at least one complete (inward or outward) zone sweep.)
CAP	None	Saved temporarily on disk in DRAW

CHKSUM	None	<p>Error criterion for the sum of the η's:</p> <p>For each polar direction specified by M, each hydro zone (I, J), each face S_1 of these hydro zones, and each fan LP on faces S_1 it is desirable,</p> <p>letting $S = \sum_{\substack{\eta_{S_1 \rightarrow S_2} \\ \text{neighboring} \\ \text{zone faces } S_2 \\ \text{and all fans} \\ L \text{ on each of} \\ \text{these faces.}}} (M, L, LP)$</p> <p>(See SUMB, SUMO)</p> <p>that $S-1$ be $< \text{CHKSUM}$. All failures to meet this criterion will result in a printed message</p> <p>(input on card) (should be $> \text{TOLER}$)</p>
COSPHI(L)	$\text{Cos } \Phi_\ell$ $= [\text{cos } (\ell - \frac{1}{2}) \Delta \Phi]$	$\text{Cos}(\text{PHI}(L))$
FIOUT	None	Saved temporarily on disk in DRAW
I	i	Index of the current radial hydro zone with outer radius X(I)
IBB	None	<p>An indicator:</p> <p>IBB = 0 means for the current M and J, the "bottom η's" sum to 1</p> <p>IBB = 1 means for the current M and J, the "bottom η's" do not sum to 1</p> <p>(See SUMB)</p>
ICHECK	None	<p>An indicator:</p> <p>ICHECK = 0 means for the current M and J, the "inside η's" sum to 1.</p> <p>ICHECK = 1 means for the current M and J, the "inside η's" do not sum to 1.</p> <p>(See SUM in OUTWRD Glossary.)</p>
IMAX	i_{max}	The number of radial hydro zones

IOO	None	<p>An indicator:</p> <p>IOO = 0 means for the current M and J, the "outside η's" sum to 1.</p> <p>IOO = 1 means for the current M and J, the "outside η's" do not sum to 1.</p> <p>(See SUMO)</p>
J	j	Index of the current horizontal layer of hydro zones with upper boundary Y(J)
JMAX	j _{max}	The number of vertical hydro zones
K	None	A running index
KK2	None	The index of the location in the BUFF2 array which contains the last entry of "DRAW data" for the current filling of the array
KMAX	None	KMAX-1 is the total number of hydro zones in the system
KMAXA	None	KMAX+1
L	ℓ	A running index
L2	None	LMAX/2
L2P1	None	L2+1
LMAX	ℓ_{\max}	<p>The total number of fans of directions to be considered in the "DRAW subroutine sequence."</p> <p>(input on card)</p> <p>(must be even and ≤ 6)</p>
M	m	Index of the current polar angle θ_m (arccos MU(M)) being considered in constructing the η 's
M2	None	MUMAX/2

MU(M)	μ_m	<p>$\cos \theta_m$, where θ_m is the mth discrete polar direction (along which the η's are calculated) measured in a clockwise sense from the vertical.</p> <p>$(\theta_m = \frac{\pi(m-\frac{1}{2})}{MUMAX}); 0 < \theta_m < \pi;$</p> <p>$\Delta\theta = \theta_{m+\frac{1}{2}} - \theta_{m-\frac{1}{2}}$</p> <p>is constant for all m; $\theta_m, m=1, MUMAX$ symmetric about $\pi/2$.)</p>
MUMAX	2M	<p>The maximum number of discrete polar directions, θ_m, considered in the construction of the η's.</p> <p>(input on card)</p> <p>(must be even and ≤ 6)</p>
N	None	A running index
NOBUFF	None	The number of output buffer (BUFF2) loads filled per inward and outward sweep of DRAW along one horizontal layer of zones (J) in direction θ_m . (Not a function of M or J)
NOO	None	<p>BUFF2(1)</p> <p>The number of inward and/or outward zone sweeps contained in the current output buffer load BUFF2. (Each buffer load must contain at a minimum one complete (inward or outward) zone sweep)</p>
P	None	Saved temporarily on disk in DRAW
PHI(L)	Φ_ℓ $= (\ell - \frac{1}{2}) \Delta\Phi$	<p>The average azimuthal direction of fan L measured in the horizontal plane containing the vector whose origin is the point in question and terminal point is the axis, in a positive sense away from the vertical axis,</p> <p>$(0 < \Phi_\ell < \pi; \Phi_\ell, \ell=1, LMAX$ symmetric about $\pi/2$; $\Delta\Phi = \Phi_\ell - \Phi_{\ell-1}$ is constant for all ℓ).</p>

PI	π	π
PP	$P = \Delta\Phi = \Phi_{\ell+\frac{1}{2}} - \Phi_{\ell-\frac{1}{2}}$	$\pi/LMAX$, a constant
RTHETA(J)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (IMAX, J) (in eV) (input on cards) (used in TRAN2)
SINLP	$\sin(\ell \cdot \Delta\Phi)$	$\sin(PP \cdot L)$
SINP2	None	$\sin(\frac{PP}{2})$
SMU(M)	$\sqrt{1 - \mu_m^2}$	$\sin(\theta_m)$
STEF	$\frac{ac}{4}$	Stefan's constant = 1.0283×10^{12} (in $\frac{\text{ergs}}{\text{cm}^2 \text{ sec eV}^4}$)
SUMB(I, LP)	None	Given directions specified by M and LP and hydro zone (I, J): $SUMB(I, LP) = \sum_{L=1, L2} ETABI(L, LP) + \sum_{L=1, LMAX} ETABT(L, LP) + \sum_{L=L2P1, LMAX} ETABO(L, LP)$ (See CHKSUM) (should = 1.)
SUMO(I, LP)	None	Given directions specified by M and LP and hydro zone (I, J): $SUMO(I, LP) = \sum_{L=1, L2} ETAOI(L, LP) + \sum_{L=1, LMAX} ETAOT(L, LP) + \sum_{L=L2P1, LMAX} ETAOO(L, LP)$ (See CHKSUM) (should = 1.)

TOLER	None	Relative error criterion imposed upon each integration performed by SUBROUTINE SINT (in decimal fraction) (input on card)
U	None	Saved temporarily on disk in DRAW
V	None	Saved temporarily on disk in DRAW
W(M)	$\frac{ \Delta\mu_m }{2}$	$ \cos[M*\Delta\theta] - \cos[(M-1)*\Delta\theta] * \frac{1}{2}$
	$= \frac{ \mu_{m+\frac{1}{2}} - \mu_{m-\frac{1}{2}} }{2}$	

FLOW OF CONTROL IN SUBROUTINE INWARD

SUBROUTINE INWARD is called by SUBROUTINE DRAW to compute, given polar angle θ_M , the η 's that couple for face S_1 of zone (I,J) all the fans LP on S_1 containing rays entering the zone traveling in an inward direction with fans L (on all the other zone faces) which contain rays leaving zone (I,J) traveling in an inward direction. A ray traveling through the system is considered traveling in an inward direction at point P if point P comes before the point D of closest approach between the ray and the system axis; the ray is considered traveling in an outward direction at point P if P comes after D.

Given below are the η 's that meet the above restriction for the θ_m 's (all $< \pi/2$) considered in the "DRAW subroutine sequence," and which are calculated in this routine:

$\eta_{B \rightarrow I}(L, LP)$

$L=1, L2; LP=1, L2$

Considering rays entering the zone bottom face traveling inward and exiting the zone inside face (rays leaving the zone through the inside face must be traveling inward).

$\eta_{O \rightarrow I}(L, LP)$ $L=1, L2; LP=1, L2$

Considering rays entering the zone outside face (must be traveling inward) and exiting the zone inside face.

 $\eta_{B \rightarrow T}(L, LP)$ $L=1, L2; LP=1, L2$

Considering rays entering the zone bottom face traveling inward and exiting the zone top face traveling inward.

 $\eta_{O \rightarrow T}(L, LP)$ $L=1, L2; LP=1, L2$

Considering rays entering the zone outside face and exiting the zone top face traveling inward

The computed η 's along with other "DRAW data" (see glossary) are stored in storage buffer array BUFF1 as they are calculated. At the end of the inward sweep of zone (I,J), BUFF1 is emptied into the output storage buffer array BUFF2 to be saved for output to the "DRAW data" file. If BUFF2 cannot accommodate the BUFF1 load of new data, it is output to the "DRAW data" file. BUFF1 must be able to store all the "DRAW data" for the inward sweep through zone (I,J); if it cannot, the problem is stopped.

Before control is returned to SUBROUTINE DRAW, preliminary calculations are performed for the sum-check on the η 's (see SUBROUTINE DRAW).

Glossary of Variables Used in SUBROUTINE INWARD

FORTTRAN Label	Report Label	Description
AA(L)	a_l	At a given radius X(I): AA(L) = SINLP(L)*X(I), the distance of closest approach (to the vertical axis) of a horizontal ray making azimuthal angle ($\pi - L*PP$) with the radius vector at X(I) (in cm)

AID	None	Temporary storage
ALPHAL	None	In computing ETABI(L, LP), for a given L and radius X(I): π - (the angle the radius vector makes with ray \vec{R} at X(I)), where ray \vec{R} = a ray traveling inward towards the vertical axis in a horizontal plane making an angle $(\pi - L*PP)$ with a radius vector at X(I-1) (in radians)
ANGL1	None	Temporary storage for information transferred to SUBROUTINES BOT and OUT
ANGL2	None	Same as above
ASN	None	Given I, $ASN = \arcsin \frac{X(I-1)}{X(I)}$ (in radians)
BUFF1	None	$\equiv CAP$ Storage buffer array for each (inward or outward) zone sweep. When each sweep is completed, the contents of BUFF1 are added to the output storage buffer array BUFF2
BUFF2	None	$\equiv P$. BUFF2(1) $\equiv NOO$ Output storage buffer array (see DRAW Glossary)
CAP	None	Saved temporarily on disk in DRAW
CH(LP)	k_q	$= JB(J-1, I, LP, M)$ (in ergs/sec) (part of "DRAW data" transferred to TRAN2)
CV(LP)	k_q	$= JO(J, I, LP, M)$ (in ergs/sec) (part of "DRAW data" transferred to TRAN2)

CVIN(L)	k_q	<p>$= JO(J, I-1, L, M)$</p> <p>(in ergs/sec)</p> <p>(part of "DRAW data" transferred to TRAN2)</p>
DBGPRT	None	<p>The debug print control:</p> <p>DBGPRT $\geq 10^{-20}$ means debug print desired; otherwise, debug print not desired</p>
DELTY	h	<p>Given J, DY(J)</p> <p>(in cm)</p>
DH(L)	t	<p>Given M, J, I, L:</p> <p>The length of the characteristic ray of fan L measured on the top face Y(J) of zone (I, J) (See SUBROUTINE DISTNC)</p> <p>(in cm)</p> <p>(part of "DRAW data" transferred to TRAN2)</p>
DV(L)	t	<p>Given M, J, I, L:</p> <p>The length of the characteristic ray of fan L measured on the vertical face X(I-1) of zone (I, J) (See SUBROUTINE DISTNC)</p> <p>(in cm)</p> <p>(part of "DRAW data" transferred to TRAN2)</p>
DY(J)	$h = z_j - z_{j-1}$	<p>Vertical length of hydro zones (I, J), $I=1, IMAX$</p> <p>(in cm)</p>
EPS	None	<p>Absolute error criterion imposed upon each integration performed by SUBROUTINE SINT</p>

ETABI(L, LP) η

Given M, J, I, L, LP:

The fraction of (energy in fan LP passing into zone (I, J) through its bottom face in 1 sec) that passes out zone (I, J) through its inside face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma = S = 0$

(part of "DRAW data" transferred to TRAN2)

ETABT(L, LP) η

Given M, J, I, L, LP:

The fraction of (energy in fan LP passing into zone (I, J) through its bottom face in 1 sec.) that passes out zone (I, J) through its top face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma = S = 0$

(part of "DRAW data" transferred to TRAN2)

ETAOI(L, LP) η

Given M, J, I, L, LP:

The fraction of (energy in fan LP passing into zone (I, J) through its outside face in 1 sec) that passes out zone (I, J) through its inside face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma = S = 0$

(part of "DRAW data" transferred to TRAN2)

ETAOT(L, LP) η

Given M, J, I, L, LP:

The fraction of (energy in fan LP passing into zone (I, J) through its outside face in 1 sec) that passes out zone (I, J) through its top face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma = S = 0$

(part of "DRAW data" transferred to TRAN2)

FORPI 4π

4π

HELP1 None

Temporary storage for information transferred to SUBROUTINE BOT.

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HELP2	None	Same as above
I	i	Index of the current radial hydro zone with outer radius X(I)
IAID	None	Flag generated for SUBROUTINES BOT and OUT
ISEND	None	Flag generated for SUBROUTINES BOT and OUT
ISWEEP	None	Current direction of DRAW sweep: = 1 means sweeping inward = 2 means sweeping outward
J	j	Axial index of the current hydro zone with upper boundary Y(J)
KK1	None	Index of the last current entry in the storage buffer array BUFF1 (must not exceed 1023)
KK2	None	Index of the last current entry in the output storage buffer array BUFF2
KKK	None	A running index
L	l	Fan index
L2	None	LMAX/2
L2P1	None	L2+1
LHI	None	Complement of the lowest index of the fans for which characteristic ray lengths are to be calculated in SUBROUTINE DISTNC
L1M1	None	The lower limit in some integrations by SUBROUTINE SINT (in cm)
L1M2	None	The upper limit in some integrations by SUBROUTINE SINT (in cm)
LLOW	None	Complement of the highest index of the fans for which characteristic ray lengths are to be calculated in SUBROUTINE DISTNC
LMAX	l_{\max}	The total number of fans of directions to be considered in the "DRAW subroutine sequence" (input on card in DRAW)

LMDA	$h\tau$	Given M and J: The horizontal distance a ray will travel in horizontal zone layer DY(J) between intersections with surface Y(J-1) and Y(J) in direction θ_m ; or $DY(J)*TAN(\theta_m)$ (in cm)
LP	l'	Fan index
M	m	Index of the current polar angle θ_m (arccos MU(M)) being considered in constructing the η 's
MU(M)	μ_m	$\cos \theta_m$. (See DRAW glossary)
NOBUFF	None	The number of output buffer (BUFF2) loads filled per inward and outward sweep of DRAW along one horizontal layer of zones(J) in direction μ_m . (Not a function of M or J)
NOO	None	BUFF2(1) The number of inward and/or outward zone sweeps contained in the current output buffer load BUFF2. (Each buffer load must contain at a minimum one complete (inward or outward) zone sweep) (part of "DRAW data" transferred to TRAN2)
PP	$\Delta\phi = p$ $= \phi_{l+\frac{1}{2}} - \phi_{l-\frac{1}{2}}$	$\pi/LMAX$, a constant
Q	None	Temporary storage
QP	None	Temporary storage
SINLP(L)	None	Sin (PP*L)
SMU(M)	$\sqrt{1-\mu_m^2}$	$SIN(\theta_m)$

SUMB(I, LP)	None	Given M, J, I, LP, in SUBROUTINE INWARD: $\text{SUMB}(I, LP) = \text{SUMB}(I, LP) + \sum_{L=1, L2} \text{ETABI}(L, LP) + \sum_{L=1, L2} \text{ETABT}(L, LP)$
SUMO(I, LP)	None	Given M, J, I, LP, in SUBROUTINE INWARD: $\text{SUMO}(I, LP) = \text{SUMO}(I, LP) + \sum_{L=1, L2} \text{ETAOI}(L, LP) + \sum_{L=1, L2} \text{ETAOT}(L, LP)$
TANN	τ	$\text{TAN } \theta_m$
TOLER	None	Relative error criterion imposed upon each integration performed by SUBROUTINE SINT (in decimal fraction) (input on card in DRAW)
W(M)	$\frac{ \Delta \mu_m }{2}$ $= \frac{ \mu_{m+\frac{1}{2}} - \mu_{m-\frac{1}{2}} }{2}$	$ \cos(M \cdot \Delta \theta) - \cos(M-1) \cdot \Delta \theta \cdot \frac{1}{2}$
X(I)	r_i	Outer radius of radial hydro zones (I, J), J=1, JMAX (in cm)
XH(L)	None	Given M, J, I, L: The radial coordinate of the intersection of the characteristic ray of fan L and the "nearest zone face" (see DH(L)) (in cm) (part of "DRAW data" transferred to TRAN2)
XMAX	r_i	=X(I), the outer radius of the radial hydro zone currently being considered (transferred to SUBROUTINES BOT and OUT) (in cm)

XMIN	r_{i-1}	=X(I-1), the inner radius of the radial hydro zone currently being considered (transferred to SUBROUTINES BOT and OUT) (in cm)
XV(L)	None	Given M, J, I, L: The radial coordinate of the intersection of the characteristic ray of fan L and the "nearest zone surface" (see DV(L)) (in cm) (part of "DRAW data" transferred to TRAN2)
YH(L)	None	Given M, J, I, L: $Y(J - \frac{1}{2})$ - (the axial coordinate corresponding to XH(L)) (in cm) (part of "DRAW data" transferred to TRAN2)
YV(L)	None	Given M, J, I, L: $Y(J - \frac{1}{2})$ - (the axial coordinate corresponding to XV(L)) (in cm) (part of "DRAW data" transferred to TRAN2)

FLOW OF CONTROL IN SUBROUTINE OUTWRD

SUBROUTINE OUTWRD is called by SUBROUTINE DRAW to compute, given polar angle θ_M , the η 's that couple for face S_1 of zone (I, J) all the fans LP on S_1 containing rays entering the zone traveling in an inward or outward direction with the fans L (on all the other zone faces) which contain rays leaving zone (I, J) traveling in an outward direction. (See SUBROUTINE INWARD for the definition of a ray traveling inward or outward at a point P.)

Given below are the η 's that meet the above restriction for the θ_M 's (all $< \pi/2$) considered in the "DRAW subroutine sequence," and which are calculated in this routine:

$\eta_{I \rightarrow O}(L, LP)$

$L=L2P1, LMAX; LP=L2P1, LMAX$

Considering rays entering the zone inside face (must be traveling outward) and exiting the zone outside face (must be traveling outward).

$\eta_{I \rightarrow T}(L, LP)$

$L=L2P1, LMAX; LP=L2P1, LMAX$

Considering rays entering the zone inside face and exiting the zone top face traveling outward.

$\eta_{B \rightarrow O}(L, LP)$

$L=L2P1, LMAX; LP=1, LMAX$

Considering rays entering the zone bottom face traveling inward and outward, and exiting the zone outside face.

$\eta_{O \rightarrow O}(L, LP)$

$L=L2P1, LMAX; LP=1, L2$

Considering rays entering the zone outside face and exiting the zone outside face.

$\eta_{B \rightarrow T}(L, LP)$

$L=L2P1, LMAX; LP=1, LMAX$

Considering rays entering the zone bottom face traveling inward and outward, and exiting the zone top face traveling outward.

$\eta_{O \rightarrow T}(L, LP)$

$L=L2P1, LMAX; LP=1, L2$

Considering rays entering the zone outside face and exiting the zone top face traveling outward.

The computed η 's along with the other "DRAW data" (see glossary) are stored in storage buffer array BUFF1 as they are calculated. At the end of the outward sweep of zone (I,J), BUFF1 is emptied into the output storage buffer array BUFF2 to be saved for output to the "DRAW data" file. If BUFF2 cannot accommodate the BUFF1 load of new data, it is output to the "DRAW data" file. BUFF1 must be able to store all the "DRAW data" for the outward sweep through zone (I,J); if it cannot, the problem is stopped.

The η 's for the rays entering the inside zone face can be checked in SUBROUTINE OUTWRD to see if they sum to 1.0 within CHKSUM (see SUBROUTINE DRAW), since rays entering a zone through the inside face can never exit the zone traveling in an inward direction. Preliminary calculations are performed for the sum-check on the remaining η 's (see SUBROUTINE DRAW).

Glossary of Variables Used in SUBROUTINE OUTWRD

FORTTRAN Label	Report Label	Description
AA(L)	A_l	See INWARD glossary
AID	None	Temporary storage
ANGL1	None	Temporary storage for information transferred to SUBROUTINES IN, OUT, and BOT
ANGL2	None	Same as above
ASN	None	Given I, ASN - $\arcsin \frac{X(I-1)}{X(I)}$
BUFF1	None	\equiv CAP Storage buffer array (see INWARD Glossary)
BUFF2	None	\equiv P. BUFF2(1) \equiv NOO Output storage array (see DRAW Glossary)
CH(LP)	k_q	$= JB(J-1, I, LP, M)$ (in ergs/sec) (part of "DRAW data" transferred to TRAN2)
CHKSUM	None	Error criterion for the sum of the ETA's (see DRAW Glossary) (input on card in DRAW)

CVIN(LP)	k_q	= JO(J, I-1, LP, M) (in ergs/sec) (part of "DRAW data" transferred to TRAN2)
CVOUT(L)	k_q	= JO(J, I, L, M) (in ergs/sec) (part of "DRAW data" transferred in TRAN2)
DBGPRT	None	The debug print control: DBGPRT $\geq 10^{-20}$ means debug print desired; otherwise, debug print not desired
DELT Y	$h = z_j - z_{j-1}$	DY(J) (in cm)
DH(L)	t	Given M, J, I, L: The length of the characteristic ray of fan L measured the top face Y(J) of zone (I, J) (see SUBROUTINE LISTNC) (in cm) (part of "DRAW data" transferred in TRAN2)
DQ	None	Temporary storage
DV(L)	t	Given M, J, I, L: The length of the characteristic ray of fan L measured on vertical face X(I) of zone (I, J) (see SUBROUTINE DISTNC) (in cm) (part of "DRAW data" transferred to TRAN2)
DX(I)	$r_i - r_{i-1}$	Radial zone length of zones (I, J), J=1, JMAX (in cm)
DY(J)	$h = z_j - z_{j-1}$	Vertical length of hydro zones (I, J). I=1, IMAX (in cm)

EPS	None	Absolute error criterion imposed upon each integration performed by SUB-ROUTINE SINT
ETABO(L, LP)	η	<p>Given M, J, I, L, LP:</p> <p>The fraction of (energy in fan LP passing into zone (I, J) through its bottom face in 1 sec) that passes out zone (I, J) through its outside face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma=S=0$</p> <p>(part of "DRAW data" transferred to TRAN2)</p>
ETABT(L, LP)	η	<p>See INWARD glossary</p> <p>(part of "DRAW data" transferred to TRAN2)</p>
ETAIO(L, LP)	η	<p>Given M, J, I, L, LP:</p> <p>The fraction of (energy in fan LP passing into zone (I, J) through its inside face in 1 sec) that passes out zone (I, J) through its outside face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma=S=0$</p> <p>(part of "DRAW data" transferred to TRAN2)</p>
ETAIT(L, LP)	η	<p>Given M, J, I, L, LP:</p> <p>The fraction of (energy in fan LP passing into zone (I, J) through its inside face in 1 sec) that passes out zone (I, J) through its top face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma=S=0$</p> <p>(part of "DRAW data" transferred to TRAN2)</p>

ETAOO(L, LP) η

Given M, J, I, L, LP:

The fraction of (energy in fan LP passing into zone (I, J) through its outside face in 1 sec) that passes back out zone (I, J) through its outside face in fan L, in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma=S=0$.

(part of "DRAW data" transferred to TRAN2)

ETAOT(L, LP) η

See INWARD glossary

(part of "DRAW data" transferred to TRAN2)

FORPI 4π

4π

HELP1 None

Temporary storage for information transferred to SUBROUTINES IN and BOT

HELP2 None

Temporary storage for information transferred to SUBROUTINE BOT

I i

Index of the current radial hydro zone with outer radius X(I)

LAID None

Flag generated for SUBROUTINES BOT and OUT

ICHECK None

An indicator:

ICHECK = 0 means for the current M and J, the "inside" η 's sum to 1.

ICHECK = 1 means for the current M and J, the "inside" η 's do not sum to 1.

(see SUM)

ISEND None

Flag generated for SUBROUTINES BOT and OUT

ISENDD None

Flag for SUBROUTINE DISTNC

ISWEEP None

Current direction of DRAW sweep:

= 1 means sweeping inward

= 2 means sweeping outward

J j

Axial index of the current hydro zone with upper boundary Y(J)

KK1	None	Index of the last current entry in the storage buffer array BUFF1 (must not exceed 1023)
KK2	None	Index of the last current entry in the output storage array BUFF2
KKK	None	A running index
L	l	Fan index
L2	None	LMAX/2
L2P1	None	L2+1
LD	None	A running index.
LHI	None	See INWARD glossary.
LIM1	None	The lower limit in some integrations by SUBROUTINE SINT (in cm)
LIM2	None	The upper limit in some integrations by SUBROUTINE SINT (in cm)
LLOW	None	See INWARD glossary.
LMAX	l_{\max}	The total number of fans of directions to be considered in the "DRAW subroutine sequence". (input on card in DRAW)
LMDA	$h\tau$	Given M and J: The horizontal distance a ray will travel in horizontal zone layer DY(J) between intersections with surface Y(J-1) and Y(J) in direction θ_m ; or, $DY(J) * \tan(\theta_m)$ (in cm)
LP	l'	Fan index.
M	m	Index of the current polar angle θ_m ($\arccos MU(M)$) being considered in constructing the η 's.
MU(M)	μ_m	$\cos \theta_m$. (See DRAW glossary).

NOBUFF	None	The number of output buffer (BUFF2) loads filled per inward and outward sweep of DRAW along one horizontal layer of zones (J) in direction μ_m (Not a function of M or J)
NOO	None	BUFF2(1) The number of inward and/or outward zone sweeps contained in the current output buffer load BUFF2. (Each buffer load must contain at a minimum one complete (inward or outward) zone sweep). (part of "DRAW data" transferred to TRAN2)
PI	π	π
PP	$\Delta\Phi = P$ $= \phi_{l+\frac{1}{2}} - \phi_{l-\frac{1}{2}}$	$\pi / LMAX$, a constant.
Q	None	Temporary storage
QP	None	Temporary storage
R	None	Temporary storage
RP	None	Temporary storage
SINLP(L)		Sin (PP*L)
SMU(M)	$\sqrt{1-\mu_m^2}$	Sin (θ_m)
SUM	None	In checking the sums of the "inside η 's," given M, J, I, LP:

$$SUMOO = \sum_{L=L2P1, LMAX} ETAIO(L, LP)$$

$$SUMTT = \sum_{L=L2P1, LMAX} ETAIT(L, LP)$$

$$SUM = SUMOO + SUMTT$$

(should = 1.)

SUMB(I, LP) None

Given M, J, I, LP, in SUBROUTINE
OUTWRD:

$$\text{SUMB}(I, LP) = \text{SUMB}(I, LP)$$

$$+ \sum_{L=L2P1, LMAX} \text{ETABO}(L, LP)$$

$$+ \sum_{L=L2P1, LMAX} \text{ETABT}(L, LP)$$

SUMO(I, LP) None

Given M, J, I, LP, in SUBROUTINE
OUTW

$$\text{SUMO}(I, LP) = \text{SUMO}(I, LP)$$

$$+ \sum_{L=L2P1, LMAX} \text{ETAOO}(L, LP)$$

$$+ \sum_{L=L2P1, LMAX} \text{ETAOT}(L, LP)$$

SUMOO None

See SUM

SUMTT None

See SUM

TANN τ Tan (θ_m)

TOLER None

Relative error criterion imposed upon
each integration performed by SUB-
ROUTINE SINT

(in decimal fraction)

(input on card in DRAW)

$$W(M) = \frac{|\Delta\mu_m|}{2}$$

$$= \frac{|\mu_{m+\frac{1}{2}} - \mu_{m-\frac{1}{2}}|}{2}$$

$$|\cos(M*\Delta\theta) - \cos((M-1)*\Delta\theta)| * \frac{1}{2}$$

X(I)

 r_i Outer radius of radial hydro cells (I, J),
J=1, JMAX

(in cm)

XH(L)	None	Given M, J, I, L: The radial coordinate of the intersection of the characteristic ray of fan L and the "nearest zone face" (see DH(L)) (in cm) (part of "DRAW data" transferred to TRAN2)
XI	None	Temporary storage
XMAX	r_i	= X(I), the outer radius of the radial hydro zone currently being considered (transferred to SUBROUTINES IN, OUT, and BOT) (in cm)
XMIN	r_{i-1}	= X(I-1), the inner radius of the radial hydro zone currently being considered (transferred to SUBROUTINES IN, OUT and BOT) (in cm)
XV(L)	None	Given M, J, I, L: The radial coordinate of the intersection of the characteristic ray of fan L and the "nearest zone surface" (see DV(L)) (in cm) (part of "DRAW data" transferred to TRAN2)
YH(L)	None	Given M, J, I, L: $Y(J - \frac{1}{2})$ - the axial coordinate corresponding to XH(L) (in cm) (part of "DRAW data" transferred to TRAN2)
YV(L)	None	Given M, J, I, L: $Y(J - \frac{1}{2})$ - the axial coordinate corresponding to XV(L) (in cm) (part of "DRAW data" transferred to TRAN2)

FLOW OF CONTROL IN SUBROUTINE DISTNC

Given a range of fans(LIM1, LIM2), the current polar angle θ_m , the direction of the current sweep in SUBROUTINE DRAW, and the current zone (I,J) under consideration in SUBROUTINE DRAW, SUBROUTINE DISTNC calculates the "length, D, of the characteristic ray" of the specified fans in zone (I,J) and the coordinates of each ray's intersection with "the nearest zone face,"

Given zone (I,J), the following definitions are in order:

A. Fans measured on horizontal face Y(J)

The "characteristic ray" of fan L is a ray measured from the midpoint of the top face of zone (I,J), $[X(I-1/2), Y(J)]$, in a direction specified by azimuthal angle $\Phi = \Phi_{LMAX+1-L}$ and polar angle $\theta = \pi - \arccos [MU(M)]$ (see Fig. 13.)

The "length, D, of the characteristic ray" is the distance from $[X(I-1/2), Y(J)]$ in the above specified direction to the "nearest zone face", which could be the outside face of zone (I,J), the bottom face of zone (I,J) or the inside face of zone (I,J).

B. Fans measured on vertical face X(I-1) or X(I)

1) Sweeping outward, $L > L2$ [Fans measured on face X(I)] .

The "characteristic ray" of fan L is a ray measured from the midpoint of the outside face of zone (I,J), $[X(I), Y(J-1/2)]$, in a direction specified by azimuthal angle $\Phi = \Phi_{LMAX+1-L}$ and polar angle $\theta = \pi - \arccos [MU(M)]$ (See Fig. 14.)

The "length, D, of the characteristic ray" is the distance from $[X(I), Y(J-1/2)]$ in the above specified direction to the "nearest zone face," which could be the bottom face of zone (I,J), the inside face of zone (I,J), or the outside face of zone (I,J).

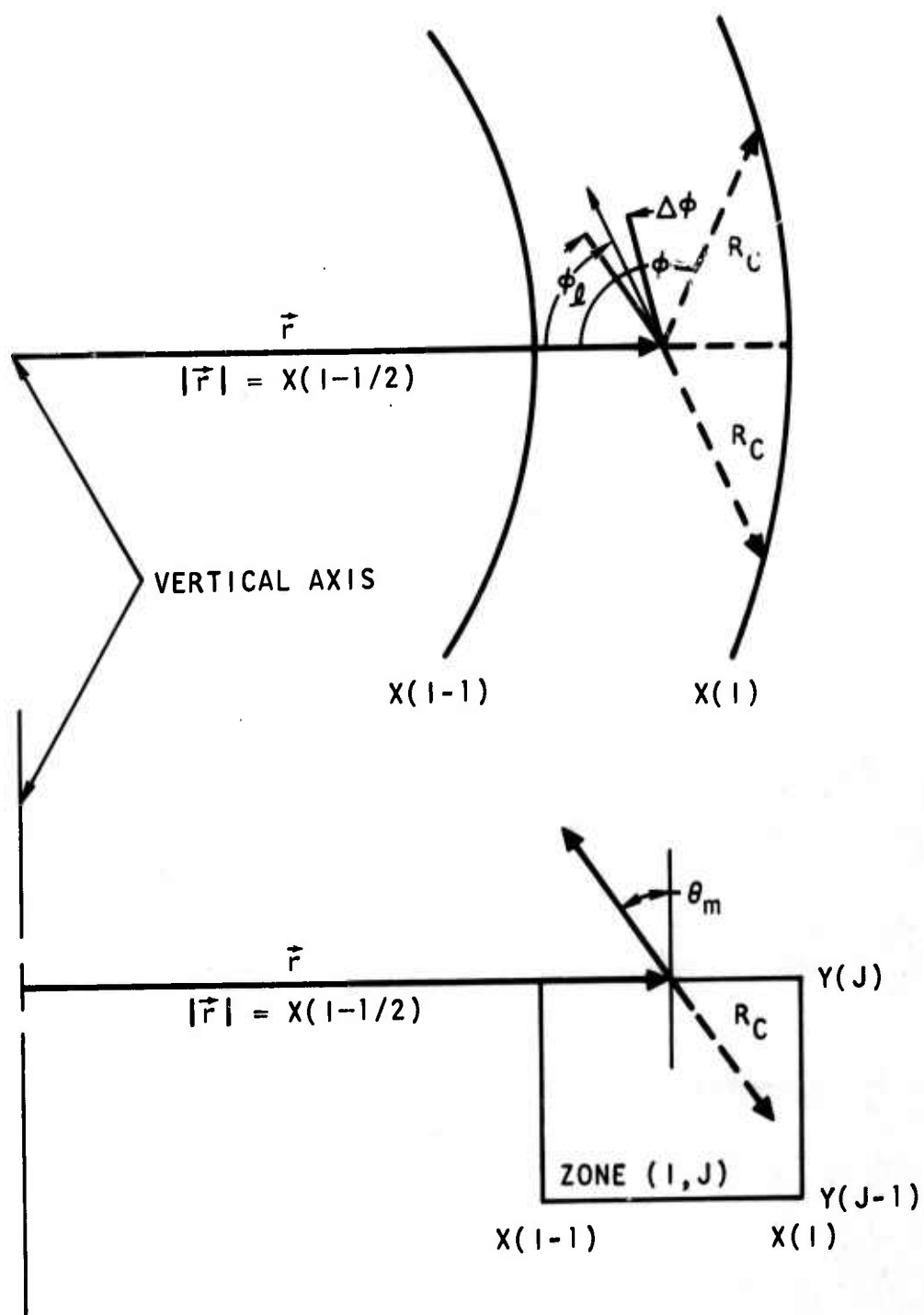


Figure 13. Characteristic Ray, R_C , for Fans Measured on Horizontal Face $Y(J)$ of Zone (I, J)

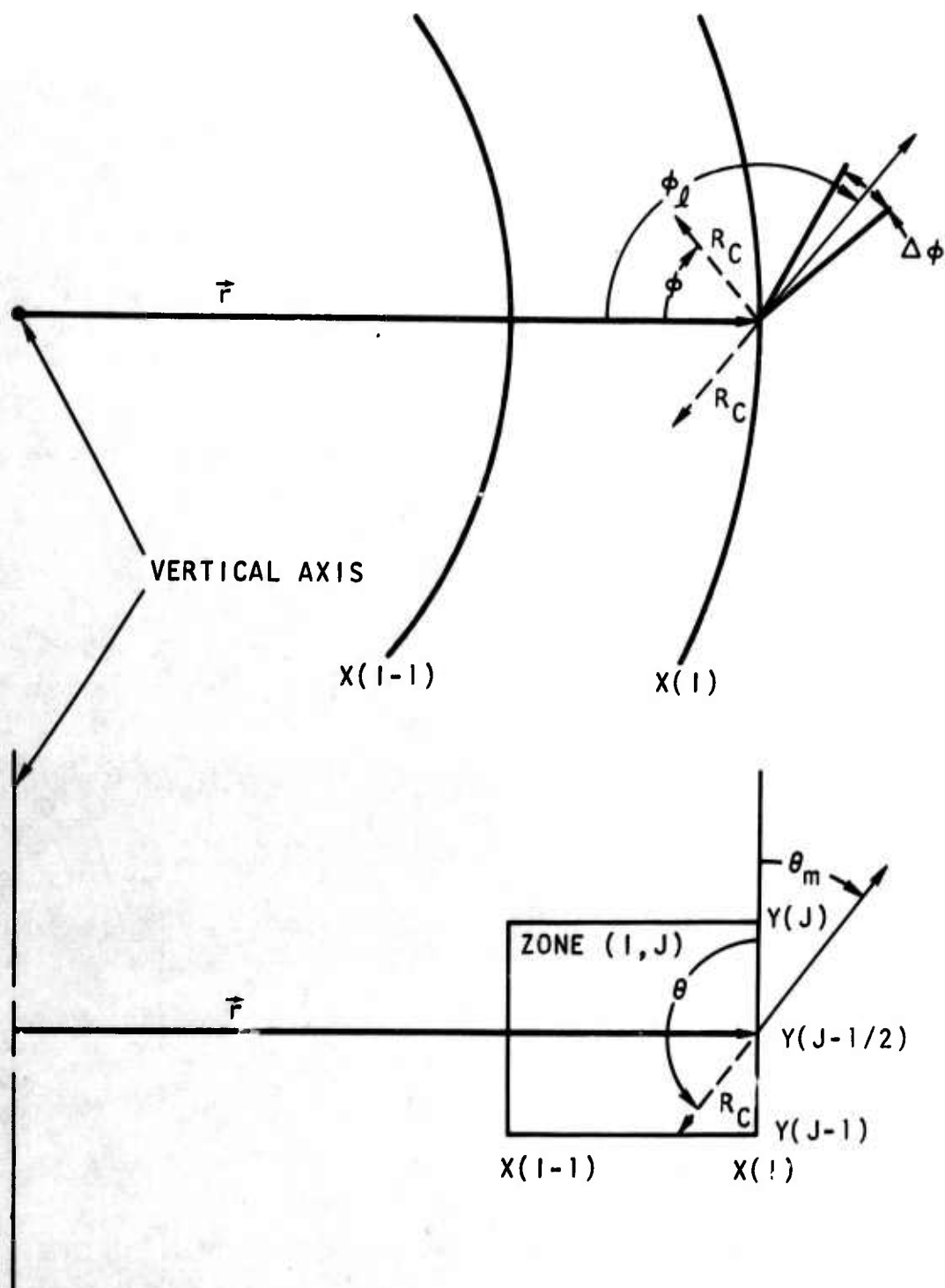


Figure 14. Characteristic Ray, R_C , for Fans Measured on Vertical Face X(I) of Zone (I, J)

- 2) Sweeping inward, $L \leq L2$ [fans measured on face $X(I-1)$].

The "characteristic ray" of fan L is a ray measured from the midpoint of the inside face of zone (I,J) , $[X(I-1), Y(J-1/2)]$, in a direction specified by azimuthal angle $\Phi = \Phi_{LMAX+1} - L$ and polar angle $\theta = \pi - \arccos [MU(M)]$ (see Fig. 15).

The "length, D , of the characteristic ray" is the distance from $[X(I-1), Y(J-1/2)]$ in the above specified direction to the "nearest zone face," which could be the bottom face of zone (I,J) or the outside face of zone (I,J) .

In the coding, the projection of D in the horizontal plane containing the characteristic ray's origin is first calculated, and then the three-dimensional length is computed.

NOTE: The coordinates (x,y) of the intersection of the "characteristic ray" and "the nearest zone face" are calculated; however, $y' = [Y(J-1/2) - y]$ is returned to SUBROUTINES INWARD and OUTWRD instead of the axial coordinate, y . This feature is to enable SUBROUTINE TRAN2 to reconstruct the axial coordinate y_{TRAN2} for polar directions $\theta \leq \pi/2$ (in which case $y_{TRAN2} = Y(J-1/2) - y'$) and $\theta > \pi/2$ (in which case $y_{TRAN2} = Y(J-1/2) + y'$) using the restrictions that the θ_M , $M=1, MUMAX$, must be symmetric about $\theta = \pi/2$.

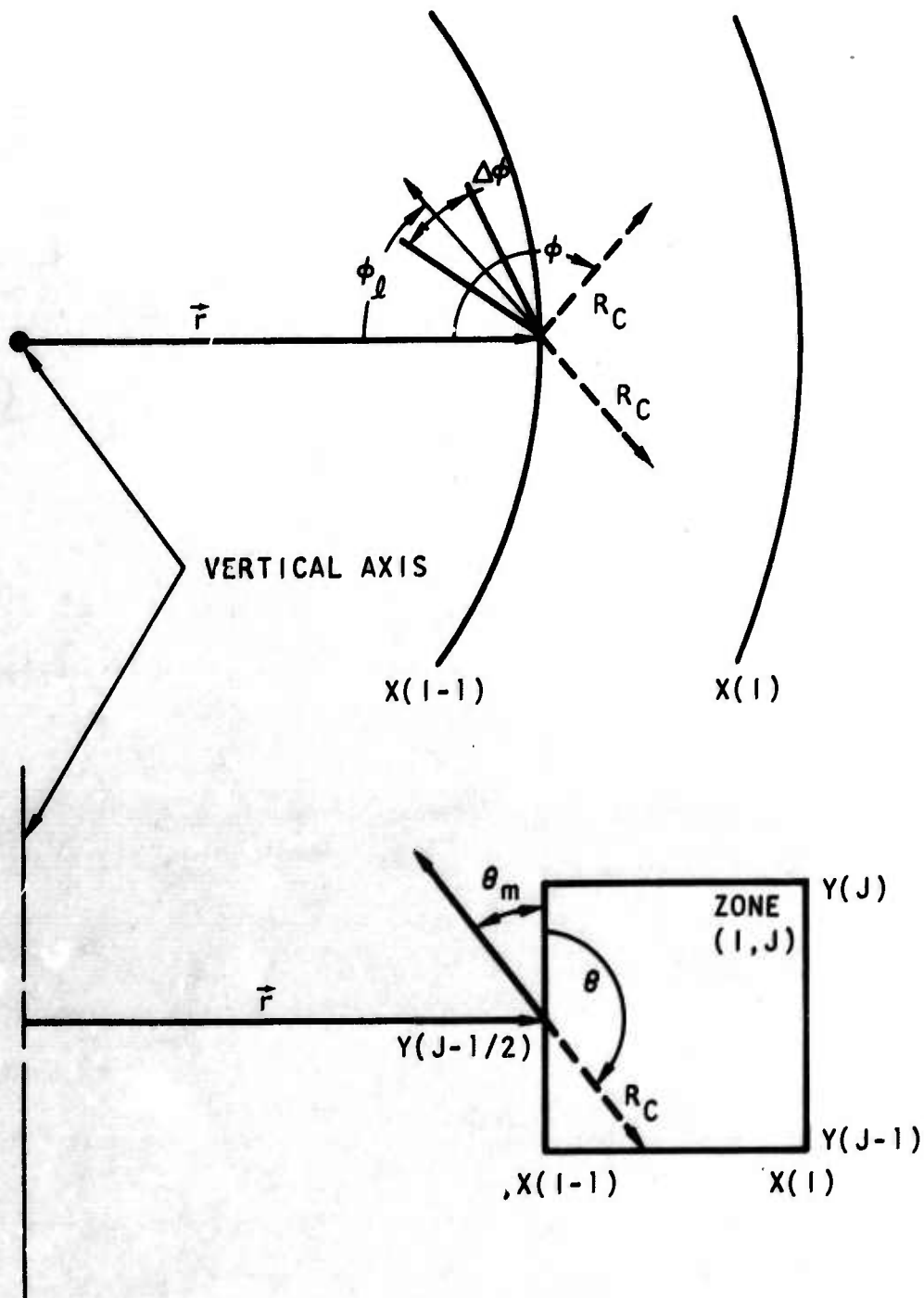


Figure 15. Characteristic Ray, R_C , for Fans Measured on Vertical Face $X(I-1)$ of Zone (I,J)

Glossary of Variables Used in SUBROUTINE DISTNC

<u>FORTTRAN</u> <u>Label</u>	<u>Report</u> <u>Label</u>	<u>Description</u>
A	None	Temporary storage.
B	None	Temporary storage.
BBB	None	Temporary storage.
COSPHI(L)	$\text{Cos}\phi_l$ $= \cos \left[\left(l - \frac{1}{2} \right) \Delta\Phi \right]$	$\text{Cos}(\text{PHI}(L))$.
DBGPRT	None	The debug print control: $\text{DBGPRT} \geq 10^{-20}$ means debug print desired; otherwise, debug print not desired.
DH(L)	t	Given M, J, I, L: The length of the characteristic ray of fan L measured on the horizontal face Y(J) of zone (I, J). (in cm)
DHH	None	Given M, J, I, L: The projection of DH(L) onto plane Y(J). (in cm)
DV(L)	t	Given M, J, I, L: The length of the characteristic ray of fan L measured on vertical face X(I) or X(I-1) of zone (I, J). (in cm)
DVV	None	Given M, J, I, L: The projection of DV(L) onto plane $Y(J - \frac{1}{2})$. (in cm)
DX(I)	$r_i - r_{i-1}$	Radial zone length of hydro zones (I, J), J=1, JMAX (in cm)

DY(J)	$h = z_j - z_{j-1}$	Vertical length of hydro zones (I, J), J=1, JMAX (in cm)
HELP	None	Temporary storage
HIT	None	In computing characteristic ray lengths, the first vertical zone boundary pierced by the current characteristic ray (in cm)
I	None	In calculating the lengths of characteristic rays originating from vertical zone faces: I=II-1 if current sweep is inward I=II if current sweep is outward such that X(I) is the face from which the current characteristic ray is measured. If calculating the characteristic ray lengths originating from horizontal zone faces: I=II
IBACK	None	Flag used in computing XV(L)
II	i	Index of the current radial hydro zone with outer radius X(II) (I of SUBROUTINE INWARD or OUTWRD)
IMAX	i_{max}	The number of radial hydro zones
ISENDD	None	A flag: = 1 means compute lengths of char- acteristic rays originating from a vertical zone face = 2 means compute lengths of char- acteristic rays originating from a horizontal zone face
ISWEEP	None	Current direction of DRAW sweep: = 1 means sweeping inward = 2 means sweeping outward
J	j	Index of the current axial (vertical) hydro zone with upper boundary Y(J)
K	None	A running index

L	ℓ	Fan
L2	None	LMAX/2
LIM1	None	= LLOW in SUBROUTINES INWARD and OUTWRD = complement of highest index of the fans for which characteristic ray lengths are to be calculated
LIM1P	None	An index limit used in a debug print (complement of LIM2)
LIM2	None	= LHI in SUBROUTINES INWARD and OUTWRD = complement of the lowest index of the fans for which characteristic ray lengths are to be calculated
LIM2P	None	An index limit used in a debug print (complement of LIM1)
LMAX	ℓ_{\max}	The total number of fans of directions to be considered in the "DRAW subroutine sequence."
M	m	Index of the current polar angle θ_m (arccos MU(M)) being considered in constructing the η 's.
MU(M)	μ_m	Cos θ_m . (See DRAW glossary.)
PHI(L)	Φ_ℓ	Φ_ℓ . (See DRAW glossary.)
SMU(M)	$\sqrt{1-\mu_m^2}$	Sin (θ_m)
X(I)	r_i	Outer radius of radial hydro zones (I, J), J=1, JMAX (in cm)
XH(L)	None	Given M, J, I, L: The radial coordinate of the intersection of the characteristic ray of fan L and the "nearest zone face" (see DH(L)) (in cm)

XV(L)	None	Given M, J, I, L: The radial coordinate of the intersection of the characteristic ray of fan L and the "nearest zone face" (see DV(L)) (in cm)
YH(L)	None	Given M, J, I, L: $Y(J - \frac{1}{2})$ - [the axial coordinate corresponding to XH(L)] (in cm)
YV(L)	None	Given M, J, I, L: $Y(J - \frac{1}{2})$ - [the axial coordinate corresponding to XV(L)] (in cm)

FLOW OF CONTROL IN FUNCTION JO(J, I, L, M)

Given the vertical cylindrical surface X(I) of height DY(J), the polar angle θ_m ($\arccos MU(M)$) with range $\Delta\mu_m$ ($2*W(M)$), and fan L with range $\Delta\phi_l$, FUNCTION JO calculates the rate at which energy is passing through surface X(I) of area $DY(J)*2*\pi*X(I)$ in fan L in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma=S=0$ divided by I_o , the intensity of the radiation field (or, equivalently, with $I_o = 1$) (See Fig. 16):

$$JO = \frac{1}{I_o} \int_{Y(J-1)}^{Y(J)} \int_{(L-1)*PP}^{(L)*PP} I_o * 2 * \pi * X(I) * \sin \theta_m * |\cos \phi_l| \Delta\mu_m d\phi dz$$

$$= X(I) * \sin \theta_m * 4 * \pi * W(M) * DY(J) * 2 * \sin\left(\frac{PP}{2}\right) \left| \cos \left[L - \frac{1}{2} \right] \right|$$

(See Eq. 151.)

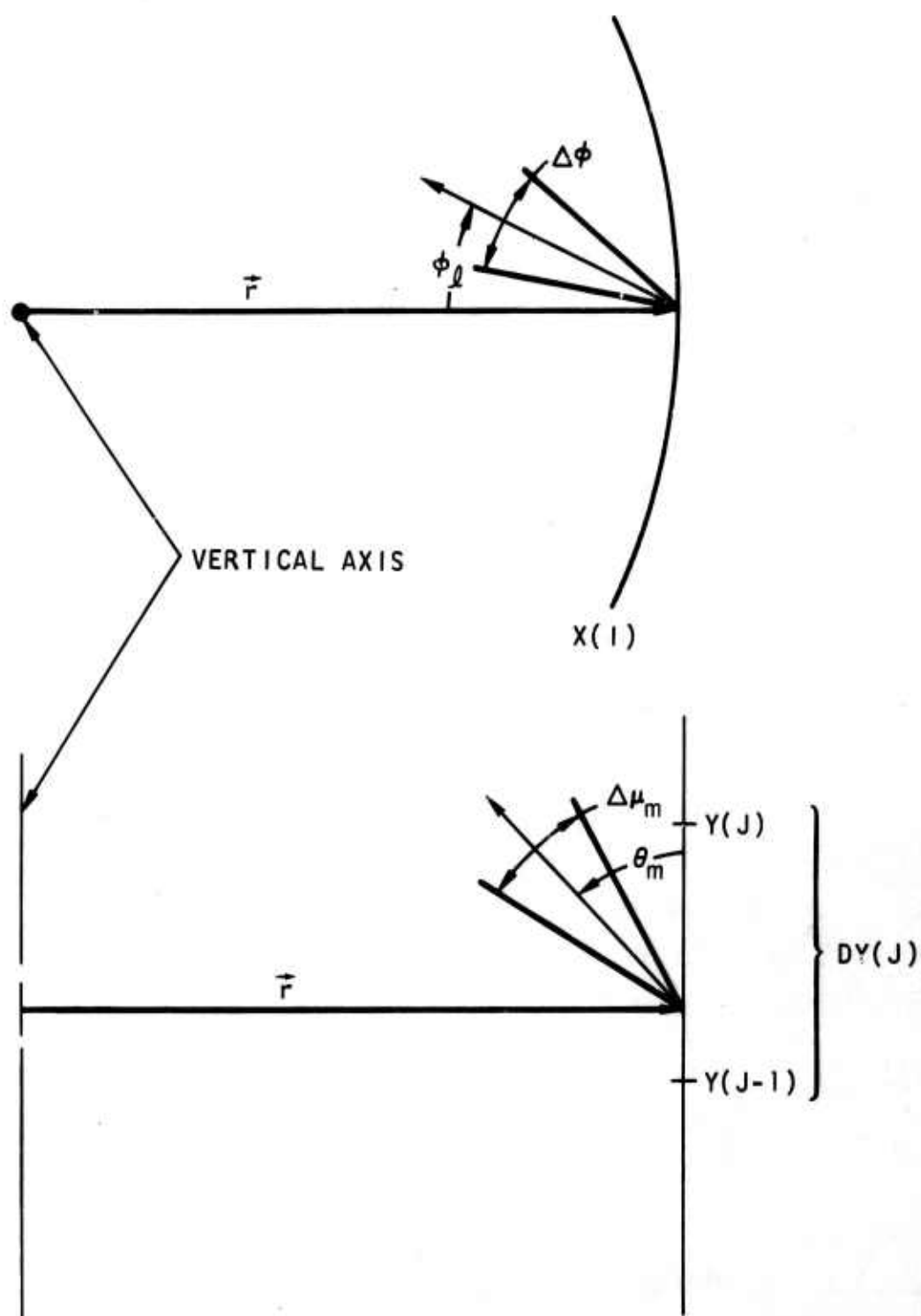


Figure 16. Fan of Directions for FUNCTION JO

Glossary of Variables Used in FUNCTION JO(J,I,L,M)

<u>FORTTRAN Label</u>	<u>Report Label</u>	<u>Description</u>
COSPHI(L)	$\text{Cos} \Phi_l$ $= \cos \left[\left(l - \frac{1}{2} \right) \Delta \Phi_l \right]$	$\cos (\text{PHI}(L))$
DY(J)	h	$Y(J) - Y(J-1)$ (in cm)
FORPI	4π	4π
I	i	Index of the current radial hydro zone with outer radius X(I)
J	j	Index of the current axial (vertical) hydro zone with upper boundary Y(J)
M	m	Index of the current polar angle θ_m (arccos MU(M) being considered)
SINP2	$\sin \frac{\Delta \Phi_l}{2}$	$\sin \left(\frac{PP}{2} \right)$
SMU(M)	$\sqrt{1 - \mu_m^2}$	$\sin (\theta_m)$
W(M)	$\frac{ \Delta \mu_m }{2}$ $= \frac{ \mu_{m+\frac{1}{2}} - \mu_{m-\frac{1}{2}} }{2}$	$ \cos [M * \Delta \theta] - \cos [(M-1) * \Delta \theta] * \frac{1}{2}$
X(I)	r_i	Outer radius of radial hydro zones (I, J), J=1, JMAX (in cm)

FLOW OF CONTROL IN FUNCTION JB(J,I,L,M)

Given horizontal annulus DX(I) of area $\pi[X^2(I) - X^2(I-1)]$, the polar angle θ_m (arccos MU(M)) with range $\Delta \mu_m$ ($2 * W(M)$), and fan L with range $\Delta \Phi_l$, FUNCTION JB calculates the rate at which energy is passing through annulus DX(I) in fan L in direction MU(M) in the case of isotropic,

homogeneous radiation $\sigma=S=0$ divided by I_0 , the intensity of the radiation field (or, equivalently, with $I_0 = 1$) (see Fig. 17):

$$\begin{aligned}
 JB &= \frac{1}{I_0} \int_{X(I-1)}^{X(I)} \int_{(L-1)*PP}^{(L)*PP} I_0 * 2 * \pi * r * |\mu_m| \Delta\mu_m d\Phi dr \\
 &\quad \text{if } A_\ell = X(I) * \sin(L*PP) \\
 &\quad \text{and } A_{\ell-1} = X(I) * \sin((L-1)*PP) \\
 &= 4 * \pi * |\mu_m| * W(M) \\
 &\quad * \left\{ \left[\frac{r^2}{2} \arcsin \left(\frac{A_\ell}{r} \right) + \frac{A_\ell}{2} \sqrt{r^2 - A_\ell^2} \right] \frac{X(I)}{X(I-1) \vee A_\ell} \right. \\
 &\quad \left. - \left[\frac{r^2}{2} \arcsin \left(\frac{A_{\ell-1}}{r} \right) + \frac{A_{\ell-1}}{2} \sqrt{r^2 - A_{\ell-1}^2} \right] \frac{X(I)}{X(I-1) \vee A_{\ell-1}} \right. \\
 &\quad \left. + \frac{\pi}{2} \left[\frac{r^2}{2} \right] \frac{X(I-1) \vee A_\ell}{X(I-1) \vee A_{\ell-1}} \right\}
 \end{aligned}$$

(See Eq. 152.)

NOTE: Since the Φ_ℓ are symmetric about $\Phi = \pi/2$, $JB(J, I, L, M) = JB(J, I, LMAX+1-L, M)$.

Glossary of Variables Used in FUNCTION JB (J, I, L, M)

<u>FORTTRAN</u> <u>Label</u>	<u>Report</u> <u>Label</u>	<u>Description</u>
A	None	Dummy variable for ARITHMETIC STATEMENT FUNCTION I1(A, R)

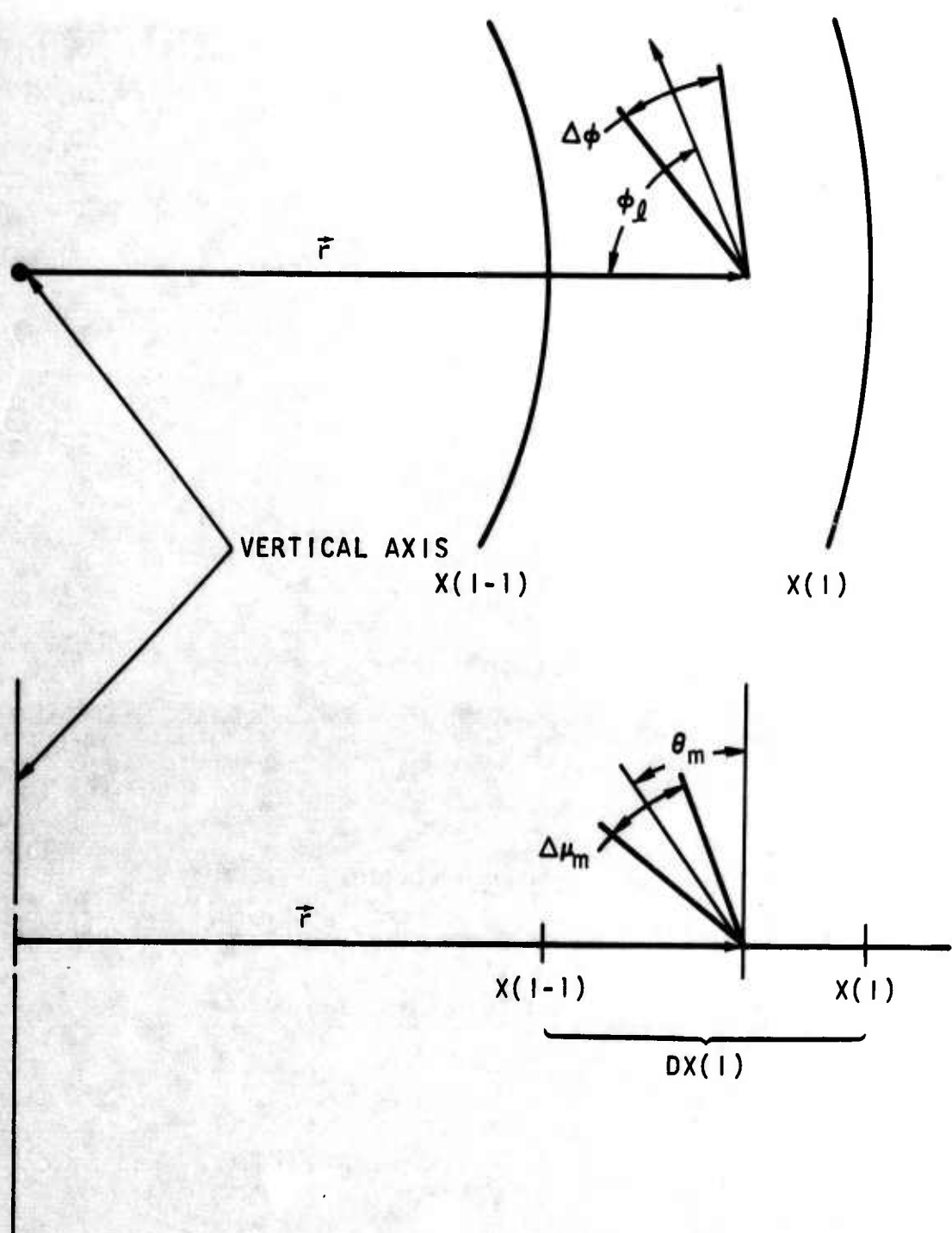


Figure 17. Fan of Directions for FUNCTION JB

AA(L)	a_l	At a given radius X(I): AA(L) = SINLP(L)*X(I), the distance of closest approach (to the vertical axis) of a horizontal ray making an azimuthal angle ($\pi - L*PP$) with a radius vector at radius X(I) (in cm)
FORPI	4π	4π
I	i	Index of the current radial hydro zone with outer radius X(I)
I1	None	ARITHMETIC STATEMENT FUNCTION
L	l	Fan index
L2	None	LMAX/2
LIM1	None	Temporary storage
LIM2	None	Temporary storage
LL	None	Index used in calculation of JB
LMAX	l_{\max}	The total number of fans being considered
M	m	Index of the current polar angle θ_m (arccos MU(M)) being considered
MU(M)	μ_m $= \cos \theta_m$	$\cos(\theta_m)$
PI4	$\pi/4$	$\pi/4$
W(M)	$\left \frac{\Delta \mu_m}{2} \right $ $= \frac{ \mu_{m+\frac{1}{2}} - \mu_{m-\frac{1}{2}} }{2}$	$ \cos[M*\Delta\theta] - \cos[(M-1)*\Delta\theta] * \frac{1}{2}$
X(I)	r_i	Outer radius of radial hydro zones (I, J), J=1, JMAX (in cm)

FUNCTION SINT (A1,B1,EPS1,FUNC)

FUNCTION SINT numerically evaluates the definite integral I,

$$I = \int_{x=A1}^{x=B1} \text{"FUNC"} dx$$

according to Simpson's method prescribed by the COMMUNICATIONS OF THE ACM ALGORITHM 182 (Ref. 7) where:

- A1 = the lower limit of integration
- B1 = the upper limit of integration
- EPS1 = the accuracy criterion in the units of the variable of integration, x.
- "FUNC" = the name of the FORTRAN FUNCTION routine which evaluates the desired integrand as a function of x.

FLOW OF CONTROL IN FUNCTION IN(Z)

FUNCTION SINT is called from SUBROUTINE OUTWRD to evaluate an η integral I of the form

$$I = \int_{Z_1}^{Z_2} [\text{INTEGRAND}] dZ$$

and FUNCTION IN is in turn called by SINT to evaluate the INTEGRAND (which is a function of Z) several times during its evaluation of I. The INTEGRAND is of the form:

$$\text{INTEGRAND} = \sin [A\sqrt{Y}]$$

where

$$A = PP*(LP-1)$$

or

$$A = \pi - \arcsin \left[X(I) - \frac{\sin (L-1)PP}{X(I-1)} \right]$$

and

 γ = a limiting angle (see GAM in glossary)Glossary of Variables Used in FUNCTION IN(Z)

<u>FORTTRAN Label</u>	<u>Report Label</u>	<u>Description</u>
ANGL1	None	<p>= $PP * (LP-1)$, the angle Φ_{l-1} measured at the inside of the current zone, or surface $X(I-1)$</p> <p>or</p> <p>= $\pi - \text{ASIN} \left[\frac{X(I) * \text{SIN} LP(L-1)}{X(I-1)} \right]$, the angle Φ_{l-1} measured at the inside of the current zone, or surface $X(I-1)$</p> <p>(in radians) (value received from SUBROUTINE OUTWRD)</p>
DELT Y	h	DY(J)
GAM	γ_0	<p>The largest azimuthal angle (≥ 0) at Z units above $Y(J-1)$ which a ray entering the inside face $X(I-1)$ of the current zone can make with the radius vector and intersect the top zone face before intersecting the outer zone face</p> <p>(in radians)</p>
HELP1	None	Temporary storage information received from SUBROUTINE OUTWRD
LMDA	$h\tau$	<p>Given J and direction θ_m:</p> <p>The horizontal distance a ray will travel between horizontal plane $Y(J-1) + Z$ and $Y(J)$; or,</p> $\left[DY(J) - Z * \text{TAN}(\theta_m) \right]$
PI	π	π
PI2	$\pi/2$	$\pi/2$
TANN	τ	$\text{TAN}(\theta_m)$

XMAX	r_i	Outer radius of the radial hydro zone currently being considered (=X(I)) (in cm) (value received from SUBROUTINE OUTWRD)
XMIN	r_{i-1}	Inner radius of the radial hydro zone currently being considered (=X(I-1)) (in cm) (value received from SUBROUTINE OUTWRD)
Z	$h-\zeta$	The distance above Y(J-1) at which the integrands for ETAIO(L, LP) and ETAIT(L, LP) are to be evaluated for FUNCTION SINT (in cm) (value received from FUNCTION SINT)

FLOW OF CONTROL IN FUNCTION BOT(Z)

FUNCTION SINT is called from SUBROUTINES INWARD and OUTWRD to evaluate η integrals and FUNCTION BOT is in turn called by SINT to evaluate the integrands (which are a function of radial distance) for these integrals several times during the course of each integration. The following table specifies each integrand type and relationship to the η calculation in SUBROUTINES INWARD or OUTWRD:

Subroutine in which integral is used	The η associated with the integral	For current fan, direction in which rays		Defining flags in this routine	Equation specifying integrand and integral
		ENTER	EXIT		
INWARD	$\eta_{B \rightarrow I}$	IN	IN	ISEND=1	Eq. (158)
INWARD	$\eta_{B \rightarrow T}$	IN	IN	ISEND=2 IAID=1	Eq. (171)
OUTWRD	$\eta_{B \rightarrow 0}$	IN	OUT	ISEND=2 IAID=4	Eq. (174)
OUTWARD	$\eta_{B \rightarrow 0}$	OUT	OUT	ISEND=2 IAID=5	Eq. (175)

Subroutine in which integral is used	The associated with the integral	For current fan, direction in which rays <u>ENTER EXIT</u>		Defining flags in this routine	Equation specifying integrand and integral
OUTWRD	$\eta_{B \rightarrow T}$	IN	OUT	ISEND=2 LAID=2	Eq. (172)
OUTWRD	$\eta_{B \rightarrow T}$	OUT	OUT	ISEND=2 LAID=3	Eq. (173)

Glossary of Variables Used in FUNCTION BOT(Z)

<u>FORTRAN Label</u>	<u>Report Label</u>	<u>Description</u>
A	w	In computing the integrands for ETABT (L, LP) and ETABO(L, LP), A= the upper azimuthal angle limit; that is, $ETA = \frac{1}{C} \int_{r_0}^{r_f} r [(A-B) \cos \phi] dr$ and B = the lower azimuthal angle limit.
AAB	None	Temporary storage
ANGL1	None	In computing the integrand for ETABI (L, LP): ANGL1 = X(I) * SINLP(LP), the distance of closest approach of a ray with azimuthal angle $\phi = LP*PP$ measured at X(I) or = X(I-1) * SINLP(L), the distance of closest approach of a ray with azimuthal angle $\phi = L*PP$ measured at X(I-1)

In computing the integrands of ETABI
(L, LP) and ETABO(L, LP):

$$\text{ANGL1} = X(I) * \text{SINLP}(LP)$$

(in cm)

ANGL2

None

= $X(I) * \text{SINLP}(LP-1)$, the distance of
closest approach of a ray with azimuthal
angle $\phi = (LP-1)*PP$ measured at $X(I)$

(in cm)

B

α

See A

GAM

γ

In computing the integrand for ETABI
(L, LP):

The largest azimuthal angle (≥ 0) which
a ray entering the bottom face of the
current zone at radius Z , $X(I-1) \leq Z \leq$
 $X(I)$, can make with the radius vector
and intersect the inside zone face before
intersecting the outside or top zone faces.

In computing the integrands for ETABT
(L, LP) and ETABO(L, LP):

Temporary storage

GAM1

γ

See GAM

GAM2

γ_0

The largest azimuthal angle ($\geq \text{GAM1}$)
that a ray entering the bottom face of
the current zone at radius Z , $X(I-1) \leq Z \leq$
 $X(I)$, can make with the radius vector
and intersect the top zone face before
intersecting the outside zone face

HELP

None

Temporary storage

HELP1

None

Temporary storage information received
from SUBROUTINES INWARD or OUTWRD

HELP2

None

Same as above

I

i

Index of the current radial hydro zone
with outer radius $X(I)$

IAID

None

Flag:

In computing the integrand for ETABT (L,LP):

- = 1 means rays enter the current zone's bottom face going inward, and exit the zone's top face going inward
- = 2 means rays enter the current zone's bottom face going inward, and exit the zone's top face going outward
- = 3 means rays enter the current zone's bottom face going outward, and exit the zone's top face going outward

In computing the integrand for ETABO (L,LP):

- = 4 means rays enter the current zone's bottom face going inward, and exit the zone's outside face going outward
- = 5 means rays enter the current zone's bottom face going outward, and exit the zone's outside face going outward

(value received from SUBROUTINES INWARD or OUTWRD)

ISEND

None

Flag:

- = 1 means integrand for ETABI (L,LP) desired
- = 2 means integrands for ETABT (L,LP) or ETABO(L,LP) desired

(value received from SUBROUTINE INWARD or OUTWARD)

LMDA	$h\tau$	<p>Given J and direction θ_m:</p> <p>The horizontal distance a ray will travel in horizontal cell layer DY(J) between intersections with surface Y(J-1) and Y(J); or,</p> $DY(J) * \tan(\theta_m)$ <p>(in cm) (value received from SUBROUTINES INWARD or OUTWRD)</p>
PI	π	π
XMAX	r_i	<p>The outer radius of the radial hydro zone currently being considered ($=X(I)$)</p> <p>(in cm) (value received from SUBROUTINES INWARD or OUTWRD)</p>
XMIN	r_{i-1}	<p>The inner radius of the radial hydro zone currently being considered ($=X(I-1)$)</p> <p>(in cm) (value received from SUBROUTINE INWARD or OUTWRD)</p>
Z	r	<p>The radial coordinate, r, at which integrands for ETAB(L, LP), ETABT(L, LP), or ETABO(L, LP) are to be evaluated for FUNCTION SINT</p> <p>(in cm) (value received from FUNCTION SINT)</p>

FLOW OF CONTROL IN FUNCTION OUT(Z)

FUNCTION SINT is called from SUBROUTINES INWARD and OUTWRD to evaluate η integrals and FUNCTION OUT is in turn called by SINT to evaluate the integrands (which are a function of axial distance) for these integrals several times during the course of each integration. The

following table specifies each integrand type and relationship to the η calculation in SUBROUTINES INWARD or OUTWRD:

SUBROUTINE in which integral is used	The η associated with the integral	For current fan, direction in which rays		Defining flags in this routine	Equation specifying integrand and integral
		ENTER	EXIT		
INWARD	$\eta_{0 \rightarrow I}$	IN	IN	ISEND=1	Eq. 157
INWARD	$\eta_{0 \rightarrow T}$	IN	IN	ISEND=2 IAID=1	Eq. 159
OUTWRD	$\eta_{0 \rightarrow T}$	IN	OUT	ISEND=2 IAID=2	Eq. 161
OUTWRD	$\eta_{0 \rightarrow 0}$	IN	OUT	ISEND=3	Eq. 162

Glossary of Variables Used in FUNCTION OUT(Z)

FORTRAN Label	Report Label	Description
A	w	<p>In computing integrands for ETAOT (L, LP) and ETAOO(L, LP), A = the upper azimuthal angle limit; that is,</p> $ETA = \frac{1}{C} \int_{Z_0}^{Z_F} [\sin(A) - \sin(B)] \vee 0. dZ$ <p>and B = the lower azimuthal angle limit (in radians)</p>
ANGL1	None	<p>In computing the integrand for ETAOI (L, LP):</p> $= PP * LP, \phi_{\ell'}$ <p>or</p> $= \frac{X(I-1)}{X(I)} * SINLP(L), \phi_{\ell} \text{ measured at } X(I)$

		<p>In computing the integrands for ETAOT (L, LP) and ETAOO(L, LP):</p> $= PP*(LP-1), \phi_{\ell'-1}$ <p>(in radians)</p>
ANGL2	None	<p>In computing the integrand for ETAOT (L, LP):</p> $= PP * LP, \phi_{\ell'}$ <p>(in radians)</p>
B	α	<p>See A</p> <p>(in radians)</p>
DELTY	h	<p>DY(J)</p> <p>(in cm)</p>
GAMMO(IG, Z)	γ	<p>IG = 1:</p> <p>= the largest azimuthal angle (≥ 0) which a ray entering Z units above Y(J-1) through the outside face of the current zone can make with the radius vector and intersect the <u>inside zone face before</u> intersecting the <u>top zone face</u></p> <p>γ_o</p> <p>IG = 2:</p> <p>= the largest azimuthal angle (≥ 0) which a ray entering Z units above Y(J-1) through the outside face of the current zone can make with the radius vector and intersect the <u>top zone face before</u> intersecting the <u>outside zone face</u></p> <p>(in radians)</p>
IAID	None	<p>Flag:</p> <p>In computing the integrand for ETAOT (L, LP):</p> <p>= 1 means considering rays entering the current zones' outer face going inward, and exiting the zone's top face going inward</p>

		<p>= 2 means considering rays entering the current zone's outer face going inward, and exiting the zone's top face going outward</p> <p>(value received from SUBROUTINES INWARD and OUTWRD)</p>
ISEND	None	<p>Flag:</p> <p>= 1 means integrand for ETAOI (L, LP) desired</p> <p>= 2 means integrand for ETAOT (L, LP) desired</p> <p>= 3 means integrand for ETAOO (L, LP) desired</p> <p>(value received from SUBROUTINES INWARD and OUTWRD)</p>
TANN	τ	$\text{TAN}(\theta_m)$
XMAX	r_i	<p>The outer radius of the radial hydro zone currently being considered (=X(I))</p> <p>(in cm)</p> <p>(value received from SUBROUTINES INWARD and OUTWRD)</p>
XMIN	r_{i-1}	<p>The inner radius of the radial hydro zone currently being considered (= X(I-1))</p> <p>(in cm)</p> <p>(value received from SUBROUTINES INWARD and OUTWRD)</p>
Z	$h - \zeta$	<p>The distance above Y(J-1) at which the integrands for ETAOI(L, LP), ETAOT (L, LP), and ETAOO(L, LP) are to be evaluated for FUNCTION SINT.</p> <p>(in cm)</p>

FLOW OF CONTROL IN FUNCTION GAMMO (IG, Z)

FUNCTION GAMMO is called by FUNCTION OUT to calculate two limiting angles that are a part of the integrands calculated in OUT.

GAMMO returns the value of one of these limiting angles for each call to the FUNCTION. (See description of GAMMO (IG, Z) in OUT glossary.)

Glossary of Variables Used in FUNCTION GAMMO (IG, Z)

FORTTRAN Label	Report Label	Description
ASN	None	$\text{Arcsin} \left(\frac{X(I-1)}{X(I)} \right)$ (in radians) (Value received from SUBROUTINES INWARDS or OUTWRD)
DELTY	h	DY(J)
HLP	None	Temporary storage
I	i	Index of the current radial hydro zone with outer radius X(I)
IG	None	Flag (See OUT Glossary)
LMDA	$\tau \zeta$	Given J and direction θ_m : The horizontal distance a ray will travel between horizontal plane Y(J-1) + Z and Y(J); or, $(DY(J) - Z) * \text{TAN}(\theta_m)$ (in cm)
LMDA2	$\tau^2 \zeta^2$	LMDA^2
TANN	τ	$\text{TAN}(\theta_m)$
XMAX	r_i	Outer radius of the radial hydro zone currently being considered (=X(I)) (in cm)
XMIN	r_{i-1}	Inner radius of the radial hydro zone currently being considered (= X(I-1)) (in cm)

Z h - ζ The distance above Y(J-1) at which GAMMO is
to be calculated for FUNCTION OUT
(in cm)
(Value received from FUNCTION OUT which
in turn receives it from FUNCTION SINT)

SECTION VII

FLOW OF CONTROL IN SUBROUTINE TRAN2 AND THE "TRANSPORT SUBROUTINE SEQUENCE" FOR A MONO- FREQUENCY, EXPLICIT TEMPERATURE PROBLEM

SUBROUTINE TRAN2 specifies the over-all flow of control through the "transportation subroutine sequence" (shown in Fig. 18). Transport calculations along characteristic ray paths supplied by SUBROUTINE DRAW are performed according to the temperatures and optical depths of the hydro zones and the η 's, or "view factors," coupling the energy passing through neighboring faces of these hydro zones.

"Transport calculations" for a given hydro zone (I,J) and a given direction of transport defines the following process:

Given the direction of transport to be in an $\begin{pmatrix} \text{inward and upward} \\ \text{outward and upward} \\ \text{inward and downward} \\ \text{outward and downward} \end{pmatrix}$

direction, the rate at which energy is leaving the zone through each fan L

on the $\begin{pmatrix} \text{inside and top} \\ \text{outside and top} \\ \text{inside and bottom} \\ \text{outside and bottom} \end{pmatrix}$ zone faces is calculated from the knowledge

of

- a. the contributions to this rate from the rate at which energy is entering the zone through all the fans LP on the

$\begin{pmatrix} \text{bottom and outside} \\ \text{inside, bottom, and outside} \\ \text{top and outside} \\ \text{inside, top, and outside} \end{pmatrix}$ zone faces assuming an isotropic,

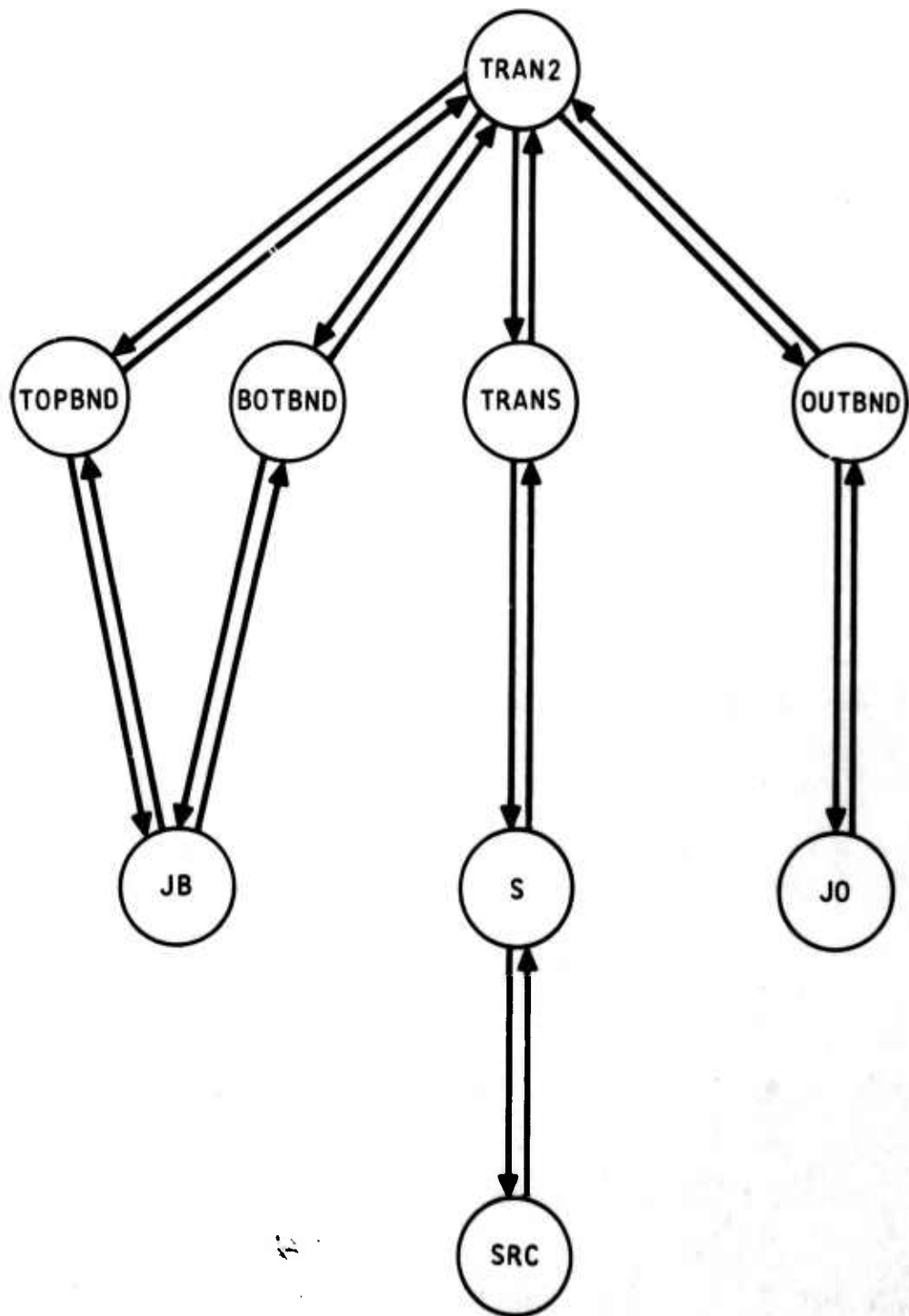


Figure 18. SUBROUTINES of the "Transport Subroutine Sequence"

homogeneous radiation field within the zone. (These contributions are represented by the term $\Sigma \eta \cdot E$ in the transport equation--see SUBROUTINE TRANS.)

and

- b. how these contributions are to be modified given a nonzero absorption coefficient and source associated with the zone.

(The modification is specified in SUBROUTINE TRANS.)

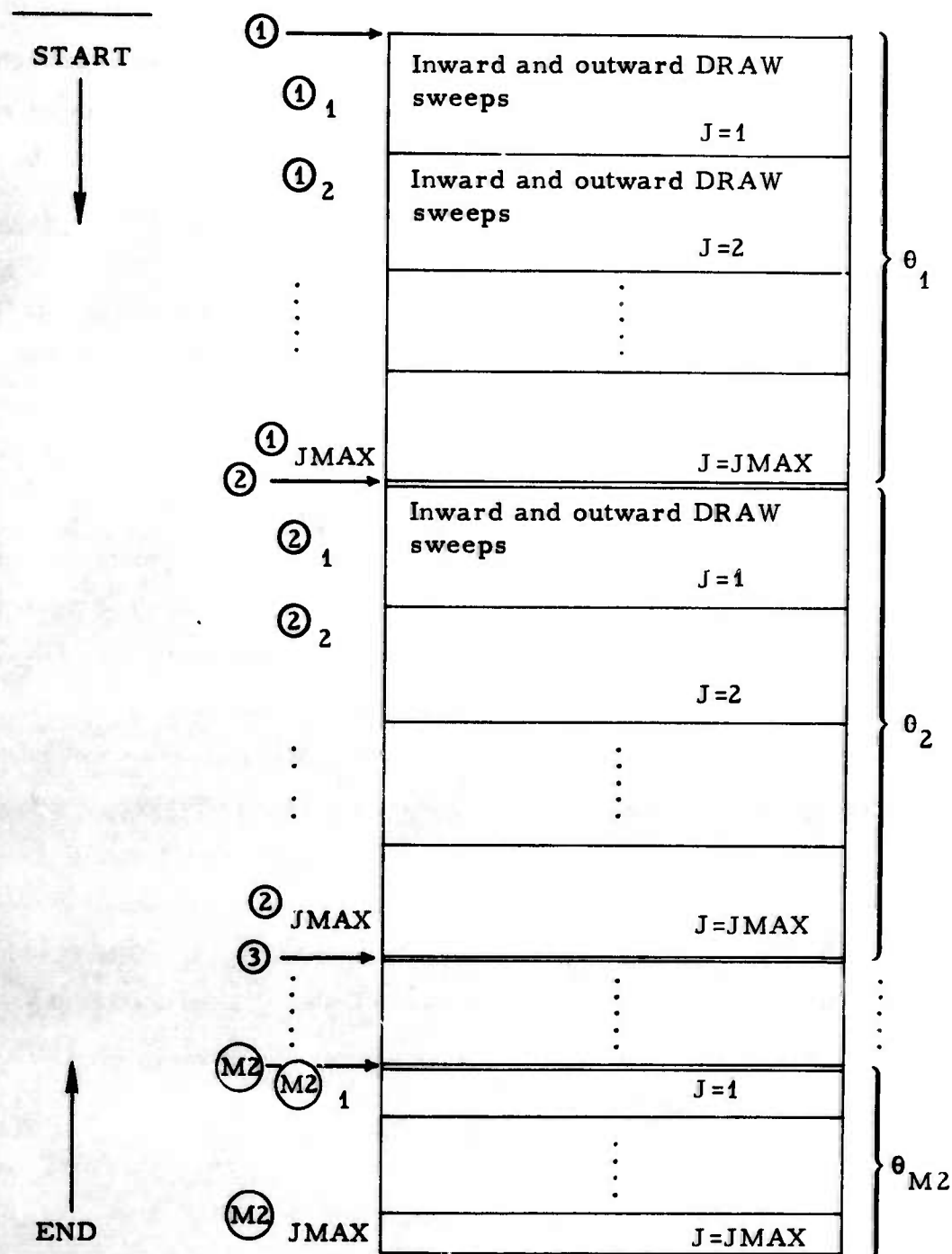
For a given polar angle θ_M transport calculations begin at the system bottom boundary and proceed (or "sweep") upward (if $\theta_M < \pi/2$) or at the system top boundary and sweep downward (if $\theta_M > \pi/2$). Sweeping upward or downward for each zone layer starts at the system outside boundary and proceeds inward, and then outward again until the system outside boundary is reached. For each θ_M , $M=1, MUMAX$, transport calculations are performed for each zone (I,J) of the entire hydro mesh.

In sweeping $\begin{pmatrix} \text{upward} \\ \text{downward} \end{pmatrix}$, the rate at which energy is leaving the $\begin{pmatrix} \text{top} \\ \text{bottom} \end{pmatrix}$ zone faces of the previous layer of zones in all fans serves as input to the $\begin{pmatrix} \text{bottom} \\ \text{top} \end{pmatrix}$ surface of the current layer of zones. When starting from the system $\begin{pmatrix} \text{bottom} \\ \text{top} \end{pmatrix}$ boundary on a $\begin{pmatrix} \text{upward} \\ \text{downward} \end{pmatrix}$ sweep, the input rates to the $\begin{pmatrix} \text{bottom} \\ \text{top} \end{pmatrix}$ zone faces are calculated from the system $\begin{pmatrix} \text{bottom} \\ \text{top} \end{pmatrix}$ boundary temperatures $\begin{pmatrix} BTHETA(I), I=1, IMAX \\ ATHETA(I), I=1, IMAX \end{pmatrix}$, and SUBROUTINES $\begin{pmatrix} \text{BOTBND} \\ \text{TOPBND} \end{pmatrix}$ perform this calculation. When sweeping $\begin{pmatrix} \text{inward} \\ \text{outward} \end{pmatrix}$, the rate at which energy is leaving the $\begin{pmatrix} \text{inside} \\ \text{outside} \end{pmatrix}$ face of the previous zone in all fans $L, \begin{pmatrix} L=1, L2 \\ L=L2P1, LMAX \end{pmatrix}$, serve as input to the fans $LP, \begin{pmatrix} LP=1, L2 \\ LP=L2P1, LMAX \end{pmatrix}$, of the $\begin{pmatrix} \text{outside} \\ \text{inside} \end{pmatrix}$ face of the current zone (except

for zones (I,J), J=1,JMAX). When starting an inward sweep through zones (IMAX,J), J=1,JMAX, from the system outside boundary, the input rates to the fans LP, LP=1, L2, are calculated from the system outside boundary temperatures (RTHETA(J), J=1,JMAX) by SUBROUTINE OUTBND.

In performing transportation calculations for a given hydro zone to calculate the rate at which energy is leaving a zone face in a fan L, the ray path of the "characteristic ray" of fan L (See SUBROUTINE DISTNC) is followed assuming that the source is a linear function of path length. The sources at the endpoints of this path are calculated by FUNCTIONS S and SRC using a bilinear interpolation on the zone-centered sources of the current and neighboring zones. If these sources are found to vary significantly (See FUNCTION TRANS) from the zone-centered source of the current zone, the source is assumed to be a constant along the ray path equal to the current zone's zone-centered source.

In sweeping upward or downward, the polar angles θ_M are taken in the order $\theta_1, \theta_{MUMAX}, \theta_2, \theta_{MUMAX-1}, \dots, \theta_{M2}, \theta_{M2+1}$ to allow for the possibility of a reflecting system top boundary (a reflecting system bottom boundary cannot be treated with the above programmed sequencing of polar angles). To perform the above sequencing, the "DRAW data" on the "DRAW data" file must be retrieved in a special manner. The "DRAW data" file is composed of blocks of logical records arranged by SUBROUTINE DRAW as shown below.



The input storage buffer array BUFF is filled and used as the transport sweeping occurs, and the "DRAW data" file illustrated above is manipulated in the following manner:

Starting the transport with the "DRAW data" file positioned at point ①.

- a. Upward sweep for θ_1 : Retrieve data from blocks ①₁, ①₂, ①₃, ..., ①_{JMAX} (each contains one or more logical records) as transport sweeping zone layers 1, 2, 3, ..., JMAX respectively.
- b. Downward sweep for θ_{MUMAX} : Backspace one block (①_{JMAX}) and retrieve data from this block as transport sweeping zone layer JMAX. Then backspace two blocks to the start of block ①_{JMAX-1} and retrieve data from this block as transport sweeping zone layer JMAX-1. ... etc.
- c. Upward sweep for θ_2 : Skip to point ②_{JMAX} and retrieve data from blocks ②₁, ②₂, ②₃, ..., ②_{JMAX} as transport sweeping zone layers 1, 2, 3, ..., JMAX respectively.

etc.

Note that the number of logical records of data per block \textcircled{N}_i is a constant for all polar angles θ_M .

For each upward or downward transport sweep associated with polar angle θ_M , the following quantities are updated:

- 1) ER(K)
- 2) EBTM, ETOP, ESIDE
- 3) W2(I)

Note each one of these quantities must be doubled since only one half of the possible fans are considered by DRAW and TRAN2 due to the symmetry inherent in a cylindrical system.

At the end of SUBROUTINE TRAN2 the ER array is checked to see if radiation transport in time DT will change the internal energy of a hydro zone by more than the allowed amount (SLUG); if it will, DT is reduced in TRAN2 to meet this restriction. The internal energy of the

system, ETH, is updated to include energy changes by radiation transport and the variables BACC, SACC, and TACC are updated.

INPUT TO "TRANSPORT SUBROUTINE SEQUENCE"

Results from "DRAW SUBROUTINE SEQUENCE" are transferred to "TRANSPORT SUBROUTINE SEQUENCE."

Conditions that Cause Termination of Execution in "TRANSPORT SUBROUTINE Sequence"

Type	Indicator	Description of Error
INPUT	S1=7.0225	MERGE \leq 0. and multifrequency problem
INPUT	S1=7.0235	IHNU < 1.
INPUT	S1=7.0250	$Q \geq 0.6$ or $Q \leq 0.4$.
INPUT	S1=7.0320	XA(K) < 0.
INPUT/ EXECUTION	S1=7.0522	An indefinite number has been created before the call to DVCHK.
INPUT	S1=7.1010	IHNU > NHNU
EXECUTION	S1=7.1068	DT < FFB, the minimum time-step control.

Results from "Transport SUBROUTINE Sequence" Transferred to HECTIC

<u>Medium</u>	<u>Variables</u>
COMMON	(a possible contribution to) DT (a contribution to) ETH ER(K), K=2, KMAX W2(I), I=1, IMAX

Glossary of Variables Used in SUBROUTINE TRAN2

FORTTRAN Label	Report Label	Description
AIX(K)	None	The total internal energy of zone K per unit mass (does not include kinetic energy). (in ergs/gm)
AMX(K)	None	The mass of zone K (in gm)
B(K)	None	In the multifrequency calculation, temporary storage
BACC	None	$\int_{t'=0}^{t'=t}$ (rate at which energy is leaving the system bottom boundary due to radiation) dt' (in ergs)
BCATAG	None	Boundary flags. A, B, and R refer to top (above), bottom, and side (right) system boundaries, respectively. If $\theta < 0$, boundary is a perfect reflector; otherwise, it is transmittive to radiation (the side cannot be a reflector).
BCBTAG	None	
BCRTAG	None	
BETA	$\frac{h\nu_1}{\theta}$	In the multifrequency calculation, the lower limit of the current frequency band
BETAP	$\frac{h\nu_2}{\theta}$	In the multifrequency calculation, the upper limit of the current frequency band
BUFF	None	\equiv CAP. PUFF(1) \equiv NOI. Input storage array for the "DRAW data" that is to be retrieved from the "DRAW data" file.
C	k_q	C is calculated in the "DRAW subroutine sequence" and used in the transport equation (SUBROUTINE TRANS) in such a way that transport calculations are exactly correct in the case of isotropic, homogeneous radiation with $s = \sigma = \text{constant}$ and I_0 , the intensity of the radiation field,

		$= \frac{s}{4\pi\sigma}$ (See Equation 177 and FUNCTION TRANS) (in cm^2 steradian) (part of "DRAW data")
D	t_q	Temporary storage for the length of the characteristic ray associated with the current zone and fan (in cm) (part of "DRAW data")
DBGPRT	None	The debug print control: $\text{DBGPRT} \geq 10^{-20}$ means debug print desired; otherwise, debug print not desired
DE	None	In calculating DT in TRAN2, the amount of energy per unit mass which a zone will lose or gain in time DT due to radiation (in ergs/gm)
DFB	None	In the multifrequency calculation, $\frac{15}{4\pi} \int \frac{\text{BETAP}}{\text{BETAe}} \frac{u^3}{e^u - 1} du$
DHNU	None	In the multifrequency calculation, the width of the current frequency band
DT	None	Time step (which can be controlled by TRAN2) (in sec)
DTEMP	None	Temporary storage
DX(I)	$r_i - r_{i-1}$	Radial zone length of zone (I, J), $I=1, \text{JMAX}$ (in cm)
DY(J)	$h=z_j - z_{j-1}$	Vertical zone length of zone (I, J) $I=1, \text{IMAX}$
E	None	The total internal energy per unit mass updated by radiation heating only (in ergs/gm)

EB(I, L)	J_q	<p>Given zone (I, J) and directions specified by M and L:</p> <p>The rate at which radiant energy is entering zone (I, J) through the bottom face (on an upward sweep) or through the top face (on a downward sweep) in fan L in direction MU(M)</p> <p>(in ergs/sec)</p>
EBTM	None	<p>The rate at which energy is entering the system through the system bottom boundary in time DT due to radiation</p> <p>(in ergs/sec)</p>
EMB1	None	In the multifrequency calculation, temporary storage
EMB2	None	Same as EMB1
ER(K)	None	<p>The rate at which zone K is gaining energy due to radiation in time DT</p> <p>(in ergs/sec)</p>
	None	In the multifrequency calculation, temporary storage
ESIDE	None	<p>The rate at which energy is entering the system through the system outside boundary in time DT due to radiation</p> <p>(in ergs/sec)</p>
ET(I, L)	J_q	<p>Given zone (I, J) and directions specified by M and L:</p> <p>The rate at which radiant energy is leaving zone (I, J) through the top face (on an upward sweep) or through the bottom face (on a downward sweep) in fan L in direction MU(M)</p> <p>(in ergs/sec)</p>
ETH	None	<p>The total energy in the system at time T</p> <p>(in ergs)</p>

ETOP	None	The rate at which energy is entering the system through the system top boundary in time DT due to radiation (in ergs/sec)
EV(I, L)	J_q	Given zone layer DY(J), zone face X(I), and directions specified by M and L: The rate at which radiant energy is passing through zone surface X(I) in fan L in direction MU(M) (in ergs/sec)
FAC	None	The axial intercept (YA) of the characteristic ray of the current fan in the radiation transport calculation equals $Y(J-1/2) + FAC * \delta y$, where $\delta y = YV(L)$ or $YH(L)$ (See INWARD, OUTWRD, or DISTNC Glossary) In the above expression for YA, $FAC = -1.$, if transporting in an upward direction or $= +1.$, if transporting in a downward direction
FACTOR	None	Temporary storage used in the calculation of ROSS(K) and PLANCK(K)
FFB	None	The minimum time-step allowed for HECTIC (in sec)
FIOUT	None	\equiv OLDTH. Saved temporarily on disk in DRAW and restored at the end of TRAN2
FOO	None	The amount of energy gain by the system in time DT due to radiation (in ergs)
GG	None	A dummy argument in the call to ES in TRAN2
HELP	None	Temporary storage

HNU	h_1	In the multifrequency calculation; the lower frequency of the current frequency band (in eV)
HNU4	None	= HNU ⁴
HNUP	h_i	In the multifrequency calculation, the upper frequency of the current frequency band (in eV)
HNUP4	None	= HNUP ⁴
I	i	In the radiation transport calculation, the index of the current radial hydro zone with outer radius X(I). Otherwise, a running index
IACT	None	Activity flag for zone K which is used in FUNCTION TRANS: IACT = 1 means zone K is active = 0 means zone K is inactive
IBUFF	None	Index of the last used entry in the current load of the input storage buffer BUFF
IHNU	None	In the multifrequency calculation, the index of the current frequency band
IMAX	i_{\max}	The number of radial hydro zones
ISEND	None	Error flag used by SUBROUTINE EDIT
ISWEEP	None	Current direction of TRAN2 sweep: = 1 means sweeping or transporting inward = 2 means sweeping or transporting outward
ITAG	None	Flag for the temperature iteration calculation: = 0 means perform temperature iteration ≠ 0 means do not perform temperature iteration
J	j	In the radiation transport calculation, index of the current axial hydro zone with upper boundary Y(J). Otherwise, a running index
JADD	None	The amount by which the value of J is incremented to give the index of the next layer of zones to be considered in the transport calculation

JMAX	j_{\max}	The number of radial hydro zones
K	None	In the radiation transport calculation, the index of the current zone. Otherwise, a running index
K1	None	Temporary storage
K2	None	Temporary storage
KDMY	None	Argument returned by the calls to FUNCTION DVCHK: = 1 means an error = 2 means no error
KFIT	None	An array constructed by HECTIC such that: JMR(KFIT(K), 2) = 1 means zone K is active JMR(KFIT(K), 2) = 0 means zone K is inactive
KMAX	None	KMAX-1 is the total number of hydro zones in the system
L		Fan index
L2	None	LMAX/2
L2P1	None	L2+1
LMAX	l_{\max}	One half the total number of fans to be considered in the radiation transport calculation
LP	l'	Fan index
M	m	Index of the current polar angle θ_m (arccos MU(M)) being considered in the radiation transport calculation
M2	None	MUMAX/2
MERGE	None	In the multifrequency calculation, a variable controlling the merging of frequency bands
MFTAG	None	Multifrequency flag: = 0 means monofrequency problem $\neq 0$ means multifrequency problem
MMM	None	A running index

MUMAX	2M	The maximum number of discrete polar directions θ_m considered in the radiation transport calculation
N	None	A running index
NB1	None	After an upward transport calculation sweep for a given $M \leq M2$, the downward sweep is performed for direction specified by M' where $M' = MUMAX+1-M$. NB1 is the number of logical records on the "DRAW data" file that must be skipped (in a backward direction) so that inward and outward transport sweeping may be performed on the zone layer JMAX
NBA	None	After sweeping inward and outward one zone layer in transporting in a downward direction, NBA = the number of records that must be skipped (in a backward direction) on the "DRAW data" file before inward and outward transport sweeping may begin on the next lower zone layer
NC	None	Integer value of cycle number
NHNU	None	In the multifrequency calculation, the total number of frequency bands
NOBUFF	None	The number of input buffer loads BUFF that will be filled for an inward or outward transport sweep along one horizontal layer of zones (J) in direction specified by M (not a function of M or J)
NOI	None	BUFF(1) The number of inward or outward DRAW zone sweeps contained in the current input buffer BUFF
NSKP	None	After sweeping in a downward direction $M(>M2)$, the number of logical records of the "DRAW data" file that must be skipped in a forward direction before an upward sweep in direction specified by $M' = MUMAX-M+2$ can begin

NVEZ	None	In the temperature iteration calculation, the temperature iteration counter
NY	None	Not used
OLDTH	None	\equiv FIOUT
P	None	\equiv PLANCK. Saved temporarily on disk in DRAW and restored at the end of TRAN2
PI	π	π
PLANCK(K)	None	The Planck mean opacity of zone K across the merged frequency band (in $\frac{1}{\text{cm}}$)
PUR	None	Not used
PUZ	None	Not used
Q	None	In the calculation of DT in TRAN2, the amount by which the internal energy per unit mass of a zone is allowed to change due to radiation transport (in ergs/gm) Also, in the multifrequency calculation, a parameter used in the merging of frequency bands
ROSS(K)	σ	The Rosseland mean absorption coefficient for zone K (in $1/\text{cm}$)
RPTAG	None	Absorption coefficient flag: $\neq 0$ means set the Planck mean absorption coefficient equal to the Rosseland mean absorption coefficient $= 0$ means keep both absorption coefficients distinct
RUR	None	Not used
RUZ	None	Not used
S1	None	Error flag for EDIT

SACC	None	$\int_{t'=0}^{t'=t} \text{(rate at which energy is leaving the system bottom boundary due to radiation)} dt'$ <p>(in ergs)</p>
SLUG	None	In the calculation of DT in TRAN2, the fraction by which the internal energy per unit mass of a zone is allowed to change due to radiation transport
SP	$4\pi S^C$	<p>The source value (the rate at which radiant energy is being generated per unit volume) assigned to the midpoint of zone K having temperature THETA(K) and absorption coefficient ROSS(K)</p> $= 4 * ROSS(K) * STEF * THETA^4(K)$ <p>(in ergs/cm³)</p>
STEF	$\frac{ac}{4}$	<p>Stefan's constant = 1.0283×10^{12}</p> <p>(in $\frac{\text{ergs}}{\text{cm}^2 \text{ sec eV}^4}$)</p>
SUM	None	<p>In the radiation transport calculation across a given zone (I,J) in a direction specified by fan L and direction M:</p> $= \sum_{LP=1, LMAX} (\text{appropriate } \eta) \cdot (\text{appropriate EV})$ $+ \sum_{LP=1, LMAX} (\text{appropriate } \eta) \cdot (\text{appropriate EB})$ $= \left(\begin{matrix} SUMO \\ \text{or} \\ SUMI \end{matrix} \right) + SUMB$ <p>(in ergs/sec)</p>
SUMB	None	See SUM
SUMO	None	See SUM
SV	None	Temporary storage
T4	None	<p>In merging frequency bands,</p> $= THETA(K)^4$ <p>(in eV⁴)</p>

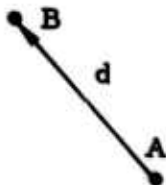
TACC	None	$\int_{t'=0}^{t'=t} \text{(rate at which energy is leaving the system top boundary due to radiation)} dt'$ (in ergs)
TAU(I)	$\pi(r_i^2 - r_{i-1}^2)$	The area of the annulus defined by radial hydro interval I (in cm^2)
TEMP(1)	None	In changing DT in TRAN2, the computed fraction by which DT is reduced to comply with the restrictive variable SLUG
THE T A(K)	θ	Temperature of zone K (in eV)
THTAMX	None	In the multifrequency calculation, the highest temperature in the system used in deciding whether to merge frequency bands (in eV)
U	None	\equiv ROSS. Saved temporarily on disk in DRAW and restored at the end of TRAN2
V	None	\equiv XA \equiv EB. Saved temporarily on disk in DRAW and restored at the end of TRAN2
VEZ	None	Not used
W2(I)	None	The rate at which energy is leaving the system top boundary adjacent to zone (I, JMAX) due to radiation, divided by TAU(I) (in $\frac{\text{ergs}}{\text{cm}^2 \text{ sec}}$)
X(I)	r_i	Outer radius of radial hydro zone I
XA	None	Temporary storage for the radial coordinate of the intersection of the current fan's characteristic ray and the "nearest zone surface." (in cm) (part of "DRAW data")

XHH	None	Radial coordinate of the midpoint of the current zone, $X(I-\frac{1}{2})+DX(I)/2$ (in cm)
Y(J)	z_j	The vertical coordinate of the upper boundary of hydro zones (I,J), $I=1, IMAX$ (in cm)
YA	None	Temporary storage for the axial coordinate of the intersection of the current fan's characteristic ray and the "nearest zone surface" (in cm) (part of "DRAW data")
YB	z_j or z_{j-1}	Temporary storage ($=Y(J)$ or $Y(J-1)$)
YH2	None	$Y(J)-DY(J)/2$, the axial coordinate of the midpoint of zones (I,J), $I=1, IMAX$ (in cm)

FLOW OF CONTROL IN FUNCTION TRANS (SIGMA, D, XA, XB, YA, YB, C, SUM, IACT, SP)

FUNCTION TRANS performs a transport calculation across the current zone of SUBROUTINE TRAN2 yielding the rate at which energy is being transported through face S of the current zone in fan L by the equation:

$$E_{S,L} = \left(\sum_{\substack{S',LP \\ S' \neq S}} \eta_{(L,LP)} \cdot E_{S',LP} \right) * e^{-\sigma d} + C \left\{ \frac{S_A}{4\pi\sigma} [1 - e^{-\sigma d}] + \frac{S_B - S_A}{4\pi\sigma^2} \left[\frac{\sigma d - 1 + e^{-\sigma d}}{d} \right] \right\}$$



where

σ = optical absorption coefficient associated with the current zone

- d = length in the current zone of the characteristic ray of the current fan L.
- S_A = the interpolated source at coordinate A, the starting point of the transport along d.
- S_B = the interpolated source at coordinate B, the end point of the transport along d.
- $\eta(L, LP)_{S' \rightarrow S}$ = "view factor" that couples the rate at which energy is entering the current zone through face S' in fan LP to the rate at which energy is leaving the current zone through face S in fan L. (See SUBROUTINES DRAW, INWARD, and OUTWRD.)
- $E_{P', LP}$ = the rate at which energy is entering the current zone through face S' in fan LP.
- $\sum_{\substack{S', LP \\ S' \rightarrow S}} \eta(L, LP)_{S' \rightarrow S} E_{S', LP}$ = the rate at which energy is leaving the current zone through face S in fan L assuming $\sigma = S_A = S_B = 0$.
- C = A constant calculated such that if $S = \sigma =$ a constant for the current zone then the transport calculation will be exact

S_A and S_B are obtained from a bilinear interpolation of the zone-centered sources of the current zone and its neighboring zones. If these interpolated values S_A and S_B and the zone-centered source S_P of the current zone differ by an amount FAC, the S_A and S_B are set equal to S_P . The source is assumed to vary linearly along the path from A to B for the transport calculation.

Glossary of Variables Used in FUNCTION TRANS (SIGMA, D, XA, XB, YA, YB, C, SUM, IACT, SP)

FORTRAN Label	Report Label	Description
A1	None	Temporary storage
A2	None	Temporary storage
A3	None	Temporary storage

A4	None	Temporary storage
A5	None	Temporary storage
C	k_q	Constant calculated such that the transport equation will yield an exact solution in the case of isotropic, homogeneous radiation with $s=\sigma$ constant and I_0 , the intensity of the radiation field, $= \frac{s}{4\pi\sigma}$ (in cm^2 steradian)
D	t_q	The length of the characteristic ray of the current fan in the current zone of SUBROUTINE TRAN2 (in cm)
FAC	f	Magnitude criterion for source term interpolation: if $\frac{\max(\text{SA}, \text{SB}, \text{SP}) - \min(\text{SA}, \text{SB}, \text{SP})}{\min(\text{SA}, \text{SB}, \text{SP})} < \text{FAC}$, normal source interpolation occurs; otherwise, the source is considered constant over the current zone
HELP	None	Temporary storage
SA	$4\pi S_i^i$	The source is considered to vary linearly from SA at point (XA, YA) to SB at point (XB, YB) $\left(\text{in } \frac{\text{ergs}}{\text{cm}^3 \text{ sec}} \right)$
SB	$4\pi S^0$	
SP	$4\pi S^c$	The source assigned to the midpoint of the current zone which has absorption coefficient SIGMA $\left(\text{in } \frac{\text{ergs}}{\text{cm}^3 \text{ sec}} \right)$
SIGMA		The Rosseland mean absorption coefficient for the current zone. (in $\frac{1}{\text{cm}}$)
SUM	None	In the radiation transport calculation across the current zone: $= \sum_{\text{LP}=1, \text{LMAX}} (\text{appropriate } \eta) \cdot (\text{appropriate EV})$

$$+ \sum_{LP=1, LMAX} (\text{appropriate } \eta) \cdot (\text{appropriate EB})$$

(in ergs/sec
(see TRAN2 glossary)

XA, XB,
YA, YB

None

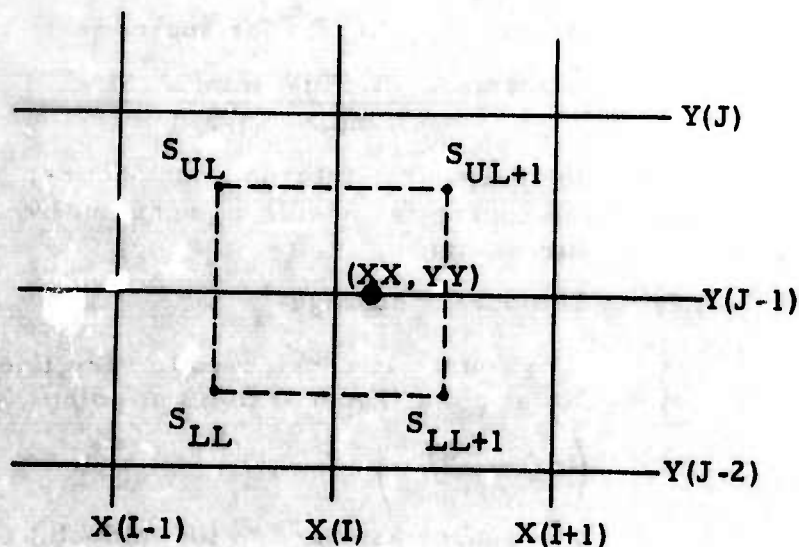
The radiation transport calculation is performed along the characteristic ray path from point (XA, YA) to point (XB, YB)

Note: (XB, YB) - (XA, YA) = D

(in cm)

FLOW OF CONTROL IN FUNCTION S(XX, YY)

Given coordinates (XX, YY), FUNCTION S calculates the source at this point with a bilinear interpolation on the zone-centered sources of the neighboring zones.



If point (XX, YY) is such that

$$X(1) - \frac{DX(1)}{2} \leq XX \leq X(IMAX) - \frac{DX(IMAX)}{2}$$

and

$$Y(1) - \frac{DY(1)}{2} \leq YY \leq Y(JMAX) - \frac{DY(JMAX)}{2}$$

an imaginary "source zone" is created that contains the point (XX, YY) and has associated with its vertices the zone-centered sources (S_{UL} , S_{UL+1} , S_{LL} , S_{LL+1}) of the neighboring zones which are obtained from

FUNCTION SRC. A bilinear interpolation is performed on the vertices of the "source zone" to obtain the source at point (XX, YY).

If point (XX, YY) lies outside the above limits, interpolation for the source at this point is performed according to the following algorithm:

- a. If $XX > X(IMAX) - \frac{DX(IMAX)}{2}$ and $[YY < Y(JMAX) - \frac{DY(JMAX)}{2}$
or $YY > Y(1) - \frac{DY(1)}{2}]$, a linear interpolation is performed
between S_{UL} and S_{LL} on YY.
- b. If $XX < X(1) - \frac{DX(1)}{2}$ and $[YY < Y(JMAX) - \frac{DY(JMAX)}{2}$
or $YY > Y(1) - \frac{DY(1)}{2}]$, a linear interpolation is performed
between S_{UL+1} and S_{LL+1} on YY.
- c. If $YY > Y(JMAX) - \frac{DY(JMAX)}{2}$ and $[XX < X(IMAX) - \frac{DX(IMAX)}{2}$
or $XX > X(1) - \frac{DX(1)}{2}]$, a linear interpolation is performed
between S_{LL} and S_{LL+1} on XX.
- d. If $YY < Y(1) - \frac{DY(1)}{2}$ and $[XX < X(IMAX) - \frac{DX(IMAX)}{2}$
or $XX > X(1) - \frac{DX(1)}{2}]$, a linear interpolation is performed
between S_{UL} and S_{UL+1} on XX.
- e. Otherwise, the source at (XX, YY) is set equal to the nearest
zone-centered source.

Glossary of Variables Used in FUNCTION S (XX, YY)

FORTTRAN Label	Report Label	Description
ALL	None	Temporary storage
ALR	None	Temporary storage

AUL	None	Temporary storage
AUR	None	Temporary storage
DELT	None	DXP * DYP (in cm ²)
DX(I)	None	Radial zone length of hydro zones (I,J), J=1,JMAX (in cm)
DXP	None	Radial "source zone" length (in cm)
DY(J)	h	Vertical length of hydro zones (I,J), J=1,JMAX (in cm)
DYP	None	Vertical "source zone" length (in cm)
I	i	Index of current radial hydro zone with outer radius X(I)
ILFT		The inner face of the "source zone" is contained within the vertical layer of zones (ILFT,J), J=1,JMAX
IMAX	i _{max}	The number of radial hydro zones
J	j	Index of the current axial hydro zone with upper boundary Y(J)
JLOW	None	The lower face of the "source zone" is contained within the horizontal layer of zones (I,JLOW), I=1,IMAX
JMAX	j _{max}	The number of axial hydro zones
KLL	None	Index of the hydro zone containing the lower left corner of the "source zone"
KUL	None	Index of the hydro zone containing the upper left corner of the "source zone"
X(I)	r _i	Outer radius of radial hydro zone I (in cm)
XLFT	None	Inner radius of the "source zone" (in cm)

XRGT	None	Outer radius of the "source zone" (in cm)
XX	None	The radius at which the source term is to be calculated (in cm)
Y(J)	z_j	The vertical coordinate of the upper boundary of hydro zones (I, J), I=1, IMAX (in cm)
YDN	None	Axial coordinate of lower face of the "source zone" (in cm)
YUP	None	Axial coordinate of the upper face of the "source zone" (in cm)
YY	None	The axial coordinate at which the source term is to be calculated (in cm)

FLOW OF CONTROL IN FUNCTION SRC (KKK)

Given hydro zone KKK, the FUNCTION SRC calculates the source associated with the midpoint of the zone. If the zone is considered "inactive" by HECTIC, then the source is zero; otherwise,

$$S_{KKK} = \sigma_{KKK} \text{ ac } \theta_{KKK}^4$$

or

$$S_{KKK} = \sigma_{KKK} \text{ ac } \theta_{KKK}^4 \int_{\frac{h\nu_1}{\theta_{KKK}}}^{\frac{h\nu_2}{\theta_{KKK}}} \frac{15}{\pi^4} \frac{u^3 du}{e^u - 1}$$

if the problem is monofrequency or multifrequency, respectively, where

- $\frac{ac}{4}$ = Stefan's constant
- ν_1, ν_2 = the frequency limits
- h = Planck's constant
- σ_{KKK} = Rosseland mean absorption coefficient of zone (KKK)
- θ_{KKK} = Temperature of zone (KKK)
- S_{KKK} = zone-centered source for zone (KKK)

Glossary of Variables Used in FUNCTION SRC (KKK)

FORTTRAN Label	Report Label	Description
DFB	None	In the multifrequency calculation, $\frac{15}{\pi^4} \int_A^B \frac{u^3}{e^u - 1} du$
HNU	$h\nu_1$	In the multifrequency calculation, the lower frequency of the current frequency band (in eV)
HNUP	$h\nu_2$	In the multifrequency calculation, the upper frequency of the current frequency band (in eV)
KFIT	None	See SUBROUTINE TRAN2
KKK	None	Index of the zone in which the source term is to be calculated
MFTAG	None	Multifrequency flag: = 0 means monofrequency frequency problem ≠ 0 means multifrequency problem
ROSS(KKK)	σ	The Rosseland mean absorption coefficient for zone KKK (in $\frac{1}{cm}$)

STEF	$\frac{ac}{4}$	Stefan's constant = 1.0283×10^{12} (in $\frac{\text{ergs}}{\text{cm}^2 \text{ sec eV}^4}$)
THETA(KKK) 0		Temperature of zone KKK (in eV)

FLOW OF CONTROL IN SUBROUTINE BOTBND

For a given polar angle $\theta_m(\arccos MU(M))$ and a system bottom boundary temperature array B_θ (BTHTETA(I), I=1,IMAX), SUBROUTINE BOTBND calculates the rate at which energy is entering each zone of the layer of zones (I,1), I=1,IMAX, in each fan L, L=1,LMAX, for a mono-frequency HECTIC problem.

The system bottom boundary is considered to be a black-body surface divided into IMAX annuli, each annulus I having area $\pi (X(I)^2 - X(I-1)^2)$ and an associated black-body temperature BTHTETA(I). The intensity of radiation, BINT, emitted into the system through each annulus I of this black-body surface integrated over all frequencies is given by

$$BINT = \frac{ac}{4\pi} B_\theta^4$$

$$\text{where } \frac{ac}{4} = \text{Stefan's constant}$$

The rate at which energy is entering zone (I,1) which borders annulus I in fan L is given by

$$EB(I,L) = (BINT) * (\text{the rate at which energy is entering zone (I,1) through its bottom face in fan L in the case of isotropic, homogeneous radiation } \sigma = S = 0 \text{ with the intensity of the radiation field} = 1)$$

Glossary of Variables Used in SUBROUTINE BOTBND

FORTTRAN Label	Report Label	Description
BINT	None	Temporary storage
BTHETA(I)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (I, 1) (in eV)
EB(I, L)	J_q	Given zone (I, J) and directions specified by M and L: The rate at which radiant energy is entering zone (I, J) through its bottom face in fan L and direction MU(M) (in ergs/sec)
I	i	A running index referring to the radial hydro zones with outer radius X(I)
IMAX	i_{max}	The number of radial hydro zones
JB(1, I, L, M)	k_q	The FUNCTION JB returns a value which is the rate at which energy is entering zone (I, 1) through the face Y(0) in fan L in direction MU(M) in the case of isotropic, homogeneous radiation ($\sigma=S=0$) with I_0 , the intensity of the radiation field, = 1. (in ergs/sec)
L	ℓ	A running index referring to the fan L
LMAX	ℓ_{max}	The total number of fans to be considered
M	m	Index of the current polar angle $\theta_m(\arccos MU(M))$ being considered
PI	π	π
STEF	$\frac{ac}{4}$	Stefan's constant = 1.0283×10^{12} $\left(\text{in } \frac{\text{ergs}}{\text{cm}^2 \text{ sec eV}^4} \right)$

FLOW OF CONTROL IN SUBROUTINE OUTBND

For a given polar angle θ_m ($\arccos MU(M)$) and a particular section $DY(J)$ of the system outside boundary with boundary temperature array R_θ ($RTHETA(J)$, $J=1, JMAX$), SUBROUTINE OUTBND calculates the rate at which energy is entering zone (IMAX, J) in each fan L, $L=1, L2$, for a monofrequency HECTIC problem.

The system outside boundary is considered to be a black-body surface divided into JMAX surfaces, each surface $DY(J)$ having area $2*\pi*X(IMAX)*DY(J)$ and an associated black-body temperature $RTHETA(J)$. The intensity of radiation, BINT, emitted into the system through surface $DY(J)$ of this black-body surface integrated over all frequencies is given by

$$BINT = \frac{ac}{4\pi} R_\theta^4$$

where

$$\frac{ac}{4} = \text{Stefan's constant}$$

The rate at which energy is entering in fan L of zone (IMAX, J) with outside face $DY(J)$ is given by

$$EV(I, L) = (BINT) * \begin{array}{l} \text{(the rate at which energy is entering} \\ \text{zone (IMAX, J) through its outside} \\ \text{face in fan L in the case of isotropic,} \\ \text{homogeneous radiation } \sigma=S=0 \text{ with the} \\ \text{intensity of the radiation field} = 1). \end{array}$$

Glossary of Variables Used in SUBROUTINE OUTBND

FORTTRAN Label	Report Label	Description
BINT	None	Temporary storage
EV(IMAX, L)	J_q	Given zone layer $DY(J)$, surface $X(IMAX)$, and directions specified by M and L: The rate at which radiant energy is entering surface $DY(J)$ in fan L in direction $MU(M)$ (in ergs/sec)

IMAX	i_{\max}	The number of radial hydro zones
J	j	Index of the current layer of hydro zones DY(J)
JO(J, IMAX, L, M)	k_q	The FUNCTION JO returns a value which is the rate at which energy is entering zone (IMAX, J) through face X(IMAX) in fan L in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma=S=0$ with I_0 , the intensity of the radiation field, = 1 (in ergs/sec)
L	l	A running index referring to the fan L
L2	None	LMAX/2
M	m	Index of the current polar angle θ_m (arccos MU(M)) being considered
PI	π	π
RTHETA(J)	θ_{out}	The temperature of the black-body system boundary adjacent to zone (IMAX, J) (in eV)
STEF	$\frac{ac}{4}$	Stefan's constant = 1.0283×10^{12} $\left(\text{in } \frac{\text{ergs}}{\text{cm}^2 \text{ sec eV}^4}\right)$

FLOW OF CONTROL IN SUBROUTINE TOPBND

For a given polar angle θ_m (arccos MU(M)) and a top (above) system boundary temperature array A_θ (ATHETA(I), I=1, IMAX), SUBROUTINE TOPBND calculates the rate at which energy is entering each zone of the layer of zones (I, JMAX), I=1, IMAX, in each fan L, L=1, LMAX for a monofrequency HECTIC problem.

The system top boundary is considered to be a black-body surface divided into IMAX annuli, each annulus I have area $\pi(XI)^2 - X(I-1)^2$ and an associated black-body temperature ATHETA(I). The intensity of radiation,

BINT, emitted into the system through each annulus I of this black-body surface integrated over all frequencies is given by:

$$BINT = \frac{ac}{4\pi} A_{\theta}^4$$

where

$$\frac{ac}{4} = \text{Stefan's constant}$$

The rate at which energy is entering zone (I,JMAX) in fan ℓ through annulus I is given by

$$EB(I, L) = (BINT) * (\text{the rate at which energy is entering zone (I,JMAX) through its top face in fan L in the case of isotropic, homogeneous radiation } \sigma = S=0 \text{ with the intensity of the radiation field} = 1.)$$

Glossary of Variables Used in SUBROUTINE TOPBND

FORTTRAN Label	Report Label	Description
ATHETA(I)	θ_{out}	Temperature of the black-body system boundary adjacent to zone (I,JMAX) (in eV)
BCATAG	None	Boundary flag referring to the above (top) system boundary < 0 means boundary is a perfect reflector > 0 means boundary is transmittive to radiation
BINT	None	Temporary storage
EB(I, L)	J_q	Given zone (I,J) and directions specified by M and L: The rate at which radiant energy is entering zone (I,J) through its top face in fan L in direction MU(M) (in ergs/sec)
I	i	A running index referring to the radial hydro zones with outer radius X(I)
IMAX	i_{max}	The number of radial hydro zones

JB(JMAX,I, L,M)	k_q	The FUNCTION JB returns a value which is the rate at which energy is entering zone (I,JMAX) through the face Y(JMAX) in fan L in direction MU(M) in the case of isotropic, homogeneous radiation $\sigma=S=0$ with I_0 , the intensity of the radiation field, = 1 (in ergs/sec)
IMAX	i_{\max}	The number of axial hydro zones
L	ℓ	A running index referring to fan L
LMAX	ℓ_{\max}	The total number of fans to be considered given θ_m
M	m	The index of the current polar angle θ_m (arccos MU(M)) being considered
PI	π	π
STEF	$\frac{ac}{4}$	Stefan's constant = 1.0283×10^{12} $\left(\text{in } \frac{\text{ergs}}{\text{cm}^2 \text{ sec eV}^4} \right)$

SECTION VIII

NONEQUILIBRIUM DIFFUSION IN HECTICSTATEMENT OF THE PROBLEM

The basis for the nonequilibrium diffusion approximation (see Section IV of Ref. 4 and Ref. 8) is the assumption that the second and higher directional moments of the radiant intensity are negligible, so that

$$I(\vec{\Omega}) = \frac{c}{4\pi} E + 3\vec{F} \cdot \vec{\Omega} \quad (181)$$

where E is the density and \vec{F} is the flux of radiant energy. When this assumption is made, Eq. (102) can be written as two equations (Eq. (52) of Ref. 4):

$$\begin{aligned} \nabla \cdot \vec{F} + c \sigma E &= s \\ \frac{c}{3} \nabla E + \sigma \vec{F} &= 0 \end{aligned} \quad (182)$$

where

$$s = \int S d\Omega \quad (183)$$

in which the direction variable $\vec{\Omega}$ no longer appears. The flux \vec{F} may also be eliminated, leaving just one second-order equation

$$-\frac{c}{3\sigma} \nabla^2 E + c \sigma E = s \quad (184)$$

for the density E .

The program TDRAD is a subroutine designed to be used in HECTIC for calculating radiative transfer using Eq. (184). To obtain numerical estimates of the energy density E , TDRAD solves the following set of finite-difference-equations (Eq. (56) of Ref. 4):

$$B_{i,j}^L (E_{i,j} - E_{i-1,j}) + B_{i,j}^R (E_{i,j} - E_{i+1,j}) + B_{i,j}^A (E_{i,j} - E_{i,j+1}) \\ + B_{i,j}^B (E_{i,j} - E_{i,j-1}) + A_{i,j} E_{i,j} = D_{i,j} \quad (185)$$

where $E_{i,j}$ is the density of radiant energy in zone $Z_{i,j}$ and

$$B_{i,j}^L = \frac{4 r_{i-1}}{3(r_i^2 - r_{i-1}^2) [\sigma_{i,j}(r_i - r_{i-1}) + \sigma_{i-1,j}(r_{i-1} - r_{i-2})]} \\ B_{i,j}^R = \frac{4 r_i}{3(r_i^2 - r_{i-1}^2) [\sigma_{i,j}(r_i - r_{i-1}) + \sigma_{i+1,j}(r_{i+1} - r_i)]} \\ B_{i,j}^A = \frac{2}{3(z_j - z_{j-1}) [\sigma_{i,j}(z_j - z_{j-1}) + \sigma_{i,j+1}(z_{j+1} - z_j)]} \\ B_{i,j}^B = \frac{2}{3(z_j - z_{j-1}) [\sigma_{i,j}(z_i - z_{j-1}) + \sigma_{i,j-1}(z_{i-1} - z_{j-2})]} \\ A_{i,j} = \sigma_{i,j} \\ D_{i,j} = S_{i,j} \quad (186)$$

Let k be a single index that enumerates zones. That is, let

$$k = (j - 1)i_{\max} + i$$

(Traditionally, the FORTRAN subscript K for zone $Z_{i,j}$ is $k + 1$ in OIL and its derivatives). Then, Eq. (185) can be conveniently written

$$\sum_{m'} B_{k,m} x_m = s_k \quad (187)$$

or, in matrix notation,

$$B x = s \quad (188)$$

where

$$\left. \begin{aligned} x_k &= E_{i,j} \\ s_k &= D_k \\ B_{k,k-1} &= -B_{i,j}^L \\ B_{k,k+1} &= -B_{i,j}^R \\ B_{k,k+I} &= -B_{i,j}^A \\ B_{k,k-I} &= -B_{i,j}^B \\ B_{k,k} &= A_{i,j} + B_{i,j}^L + B_{i,j}^R + B_{i,j}^A + B_{i,j}^B \\ B_{k,m} &= 0 \quad \text{otherwise} \end{aligned} \right\} \quad (189)$$

where $\gamma = i_{\max}$.

A method of successive approximations due to Oliphant (Ref. 9) is used in TDRAD to solve Eq. (188). It was soon discovered that the rate of convergence of Oliphant's method is often very low, especially when the zones of the HECTIC mesh are optically thin. Intensive effort has been exerted in the search for efficient methods of solutions of systems of the form of Eq. (185), especially in the fields of thermal and neutron diffusion (Ref. 10).

The most successful procedures have been successive over-relaxation, or SOR, and alternating-direction-implicit, or ADI. Accordingly, tests were conducted to compare the performance of these procedures with Oliphant's to see which would be most suitable for application to problems of the type that TDRAD is designed to treat.

FORTTRAN expressions for the above-mentioned variables are listed below, where

$$K = k + 1 = (j - 1) i_{\max} + i + 1$$

$$B_{i,j}^L = XA(K)$$

$$B_{i,j}^R = XC(K)$$

$$B_{i,j}^B = XALP(K)$$

$$B_{i,j}^A = XGAM(K)$$

$$D_{i,j} = XD(K)$$

$$A_{i,j} = -XA(K) - XB(K) - XC(K) - XALP(K) - XGAM(K)$$

$$E_{i,j} = ERAD(K)$$

The appendix contains listings of the FORTRAN source statements for a version of TDRAD that uses the ADI method. The programming of Oliphant's method and SOR is included in supplementary listings.

COMMON FEATURES OF THE METHODS

All three methods investigated are based on the idea of using an approximate inverse of the matrix B to obtain a correction to an estimate of the solution of Eq. (188). The methods differ mainly in their procedures for obtaining the approximate inverse.

To be more explicit, suppose C is an approximate inverse of B in the

sense that

$$CB = I - E \quad (190)$$

where I is the identity and E is small. Suppose further that x^l is the l^{th} in a series of estimates of the solution of $Bx = s$.

The correction to x^l is taken to be $C(s - Bx^l)$, so the $l + 1^{\text{st}}$ estimate becomes

$$x^{l+1} = x^l + C(s - Bx^l) \quad (191)$$

Let $By = s$ and let the error in x^l be $e^l = y - x^l$. Then $s - Bx^l = s - B(y - e^l) = Be^l$, so

$$\begin{aligned} e^{l+1} &= y - x^{l+1} \\ &= y - [x^l + C(s - Bx^l)] \\ &= y - (x^l + CBe^l) \\ &= e^l - CBe^l \\ &= (I - CB) e^l \\ &= Ee^l \end{aligned}$$

It follows that $e^l = E^l e^0$, where $E^0 = I$ and $E^{l+1} = EE^l$, the $l+1^{\text{st}}$ power of E . Therefore, $x^l \rightarrow y$ provided that E is so small that the moduli of its eigenvalues are all less than unity.

Another common feature of the various methods is the use of convergence accelerating parameters. What is done amounts to setting

$$x^{l+1} = x^l + \omega C(s - Bx^l) \quad (192)$$

where ω is an adjustable scalar parameter. It is often true that methods of successive approximation of the type described by Eq. (191) underestimate the size of the correction, and using Eq. (192) with $\omega > 1$, but not too large, will improve the convergence rate. Actually, Eq. (192) is not used exactly as written in any of the three methods that were tested, as will be seen in the following detailed discussion of the methods.

OLIPHANT'S METHOD

Oliphant's method (Ref. 9) is, briefly, to modify the matrix B in such a way that it becomes easy to invert with the expectation that if the modification is not too great, then the inverse of the new matrix will closely approximate that of B.

First, however, let U be the upper part of B: i.e., let

$$U_{k,m} = \begin{cases} B_{k,m} & \text{if } m > k \\ 0 & \text{otherwise} \end{cases} \quad (193)$$

and set

$$C = B + \lambda U \quad (194)$$

The next step is to find a small matrix H such that C + H may be expressed as a product of easily determined triangular factors. Thus, let

$$C + H = WV \quad (195)$$

where

$$\begin{aligned} W_{k,l} &= 0 & \text{for } l > k \\ V_{l,m} &= 0 & \text{for } l > m \end{aligned} \quad (196)$$

The particular solution of Eq. (195) that is used is

$$W_{k,k-1} = B_{k,k-1}$$

$$W_{k,k-1} = B_{k,k-1}$$

$$W_{k,k} = B_{k,k} - W_{k,k-1} V_{k-1,k} - W_{k,k-1} V_{k-1,k}$$

$$W_{k,m} = 0 \quad \text{for } m = k, k-1, \text{ or } k-1$$

$$\begin{aligned}
V_{k,k+I} &= (\lambda + 1) B_{k,k+I} / W_{k,k} \\
V_{k,k+1} &= (\lambda + 1) B_{k,k+1} / W_{k,k} \\
V_{k,k} &= 1 \\
V_{k,m} &= 0 \quad \text{for } m \neq k, k+1 \text{ or } k+I \\
H_{k,k+I-1} &= W_{k,k-1} V_{k-1, k+I-1} \\
H_{k,k-I+1} &= W_{k,k-I} V_{k-I, k-I+1} \\
H_{k,m} &= 0 \quad \text{for } m \neq k-I+1 \text{ or } k+I-1
\end{aligned} \tag{197}$$

where $I = i_{\max}$.

Note that $W_{k,m}$ and $V_{k,m}$ are expressed in terms of B and $W_{k',m'}$ and $V_{k',m'}$ for $k' < k$ or $k' = k$ and $m' < m$, so that there is no problem of finding a sequence in which to carry out the indicated operations. The verification of Eq. (197) is accomplished by direct substitution in Eq. (195). Adding $(\lambda U + H)x$ to both sides of Eq. (188) yields

$$(B + \lambda U + H)x = s + (\lambda U + H)x$$

or

$$W V x = s + (\lambda U + H)x \tag{198}$$

so that

$$x = V^{-1} W^{-1} [s + (\lambda U + H)x] \tag{199}$$

If x^l is a trial solution, one revises it by substituting it for x in the RHS of Eq. (199). Thus,

$$x^{l+1} = V^{-1} W^{-1} [s + (U + F : l)] \tag{200}$$

The matrix $V^{-1}W^{-1}$ plays the role of approximate inverse of B. To see the analogy between Eq. (191) and Eq. (200), observe that

$$\begin{aligned} x^{l+1} &= x^l + V^{-1}W^{-1} \left[s + (\lambda U + H) x^l - W V x^l \right] \\ &= x^l + V^{-1}W^{-1} (s - Bx^l) \end{aligned}$$

The numerical work is rather simple. To see what is involved, rewrite Eq. (199) in three parts as follows:

$$\left. \begin{aligned} y &= s + (\lambda U + H) x^l \\ Wz &= y \\ V x^{l+1} &= z \end{aligned} \right\} \quad (201)$$

Writing these equations out in full yields

$$\begin{aligned} y_k &= s_k + \lambda B_{k,k+1} x_{k+1}^l + \lambda B_{k,k+I} x_{k+I}^l \\ &\quad + H_{k,k+I-1} x_{k+I-1}^l + H_{k,k-I+1} x_{k-I+1}^l \end{aligned} \quad (202)$$

$$z_k = (y_k - W_{k,k-1} z_{k-1} - W_{k,k-I} z_{k-I}) / W_{k,k} \quad (203)$$

$$x_k^{l+1} = z_k - V_{k,k+1} x_{k+1}^{l+1} - V_{k,k+I} x_{k+I}^{l+1} \quad (204)$$

Equations (202) and (203) are evaluated simultaneously in order of increasing k , so the y 's need not actually be stored. Equation (205) is then evaluated in reverse order so that x_{k+1}^{l+1} and x_{k+I}^{l+1} will be available for calculating x_k^{l+1} . Oliphant's acceleration procedure amounts to replacing Eq. (204) by

$$x_k^{l+1} = \omega (z_k - V_{k,k+1} x_{k+1}^{l+1} - V_{k,k+I} x_{k+I}^{l+1}) + (1 - \omega) x_k^l \quad (205)$$

Usually, $-1 \leq \lambda \leq 1$ and $0.5 \leq \omega \leq 1.5$.

FORTTRAN expressions for the above-mentioned variables are supplied in the following list (because of the transience of FORTTRAN assignments, "x = A and y = A" does not imply "x = y"):

$$k = K-1$$

$$I = \text{IMAX}$$

$$B_{k,k} = \text{XB}(K)$$

$$B_{k,k-1} = \text{XA}(K)$$

$$B_{k,k+1} = \text{XC}(K)$$

$$B_{k,k-I} = \text{XALP}(K)$$

$$B_{k,k+I} = \text{XGAM}(K)$$

$$x_k^l = \text{ERAD}(K)$$

$$s_k = -\text{XD}(K)$$

$$\lambda = \text{CMXK} - 1.0$$

$$U_{k,k+1} = \text{XC}(K)$$

$$U_{k,k+I} = \text{XGAM}(K)$$

$$W_{k,k} = \text{XB}(K)$$

$$W_{k,k-1} = \text{XA}(K)$$

$$W_{k,k-I} = \text{XALP}(K)$$

$$V_{k,k+I} = \text{XGAM}(K)$$

$$V_{k,k+1} = \text{XC}(K)$$

$$H_{k,k+I-1} = XA(K)*XGAM(K-1)$$

$$H_{k,k-I+1} = XALP(K)*XC(K-IMAX)$$

$$y_k = XH$$

$$z_k = XG(K)$$

$$x_k^{\ell+1} = ENEW$$

ALTERNATING-DIRECTION-IMPLICIT METHOD

The alternating-direction-implicit, or ADI, method for solving elliptic systems of difference equations is usually attributed to Peaceman and Rachford (Ref. 11). Their procedure is, briefly, to solve Eq. (185) alternately by rows and by columns. The variant that was chosen for the comparison test operates in the following way.

Let $E_{i,j}^{\ell}$ be the ℓ^{th} estimate of $E_{i,j}$ and let ω be a parameter. The procedure is first to solve

$$\begin{aligned} B_{i,j}^L (E_{i,j}^{\ell+\frac{1}{2}} - E_{i-1,j}^{\ell}) + B_{i,j}^R (E_{i,j}^{\ell+\frac{1}{2}} - E_{i+1,j}^{\ell}) + (1 + \omega) A_{i,j} E_{i,j}^{\ell+\frac{1}{2}} \\ = D_{i,j} + \omega E_{i,j} - B_{i,j}^A (E_{i,j}^{\ell} - E_{i,j+1}^{\ell}) - B_{i,j}^B (E_{i,j}^{\ell} - E_{i,j-1}^{\ell}) \end{aligned} \quad (206)$$

for $E_{i,j}^{\ell+\frac{1}{2}}$ and then to solve

$$\begin{aligned} B_{i,j}^A (E_{i,j}^{\ell+1} - E_{i,j+1}^{\ell+1}) + B_{i,j}^B (E_{i,j}^{\ell+1} - E_{i,j-1}^{\ell+1}) + (1 + \omega) A_{i,j} E_{i,j}^{\ell+1} \\ = D_{i,j} + \omega E_{i,j}^{\ell+\frac{1}{2}} - B_{i,j}^L (E_{i,j}^{\ell+\frac{1}{2}} - E_{i-1,j}^{\ell+\frac{1}{2}}) - B_{i,j}^R (E_{i,j}^{\ell+\frac{1}{2}} - E_{i+1,j}^{\ell+\frac{1}{2}}) \end{aligned} \quad (207)$$

The parameter ω usually falls in the range $0 \leq \omega \leq 20$.

The system of equations represented by Eq. (206) may be reduced to j_{\max} independent subsystems, one for each value of j . Each of the subsystems has the form

$$a_i x_{i-1} + b_i x_i + c_i x_{i+1} = d_i \quad (208)$$

for $1 \leq i \leq i_{\max}$. To solve Eq. (210), look for a solution of the form

$$x_i = \gamma_i x_{i+1} + \delta_i \quad (209)$$

Using Eq. (209) to eliminate x_{i+1} and x_{i-1} from Eq. (208) yields

$$a_i(\gamma_{i-1} x_i + \delta_{i-1}) + b_i x_i + c_i \frac{x_i - \delta_i}{\gamma_i} = d_i$$

or

$$\left(a_i \gamma_{i-1} + b_i + \frac{c_i}{\gamma_i} \right) x_i = d_i + \frac{c_i \delta_i}{\gamma_i} - a_i \delta_{i-1} \quad (210)$$

Setting both sides of Eq. (210) to zero gives

$$\gamma_i = \frac{-c_i}{a_i \gamma_{i-1} + b_i} \quad (211)$$

$$\delta_i = -\frac{1}{c_i} (d_i - a_i \delta_{i-1})$$

$$= \frac{d_i - a_i \delta_{i-1}}{a_i \gamma_{i-1} + b_i} \quad (212)$$

Since $a_1 = 0$, the values of γ_0 and δ_0 are irrelevant. Equations (211) and (212) may be solved recursively for γ_i and δ_i , $i = 1, 2, \dots, I = i_{\max}$. Since $C_I = 0$, $\gamma_I = 0$. Hence, Eq. (209) may be solved for x_i , $i = I, I-1, \dots, 2, 1$, starting with $x_I = \delta_I$.

The system of equations represented by Eq. (207) is of the same type, and the same procedure may be applied to it.

In the FORTRAN program, γ and δ are stored in arrays XGI and XDI, respectively. E^l is in ERAD and $E^{l+\frac{1}{2}}$ is in EH.

SUCCESSIVE OVERRELAXATION METHOD

The method of successive overrelaxation (SOR) was probably known to Gauss, but Southwell (Ref. 12) deserves credit for reviving it. In its most primitive form, the method reduces to solving Eq. (185) for $E_{i,j}$ using whatever is the best available estimate for $E_{i,j\pm 1}$ and $E_{i\pm 1,j}$, then extrapolating the result by some constant amount, and then proceeding to some new (i,j) . In the present context, SOR refers to the process of successively solving Eq. (185) for various subsets of the set of all the E 's, using the latest estimates of the values of the remaining ones.

To be more specific, let $i_{1,n}$ and $i_{2,n}$ be given for $1 \leq n \leq N$, where $1 \leq i_{1,n} \leq i_{2,n} \leq i_{\max}$. One iteration of the procedure consists of solving N sets of equations, where set n consists of Eq. (185) for $i_{1,n} \leq i \leq i_{2,n}$ and $1 \leq j \leq j_{\max}$. Of course, if $i_{1,1} = 1$ and $i_{2,1} = i_{\max}$, then the first step of the first iteration would produce the exact solution. Usually, the number of columns

$$\Delta i_n = i_{2,n} - i_{1,n} + 1 \quad (213)$$

lies between 1 and 5.

The procedure for carrying out step n of an iteration is as follows. First, the residuals $r_{i,j}$ of Eq. (185) are calculated,

$$r_{i,j} = D_{i,j} - A_{i,j} F_{i,j} - B_{i,j}^L (E_{i,j} - E_{i-1,j}) - \dots, \quad (214)$$

for $1 \leq j \leq j_{\max}$ and $i_{1,n} \leq i \leq i_{2,n}$, using the latest values for the E 's. Then corrections $\epsilon_{i,j}$ are calculated using Eq. (185) with $\epsilon_{i,j}$ in place of $E_{i,j}$ and $r_{i,j}$ in place of $D_{i,j}$. Equation (185) then reads

$$\begin{aligned}
& B_{i,j}^L (\epsilon_{i,j} - \epsilon_{i-1,j}) + B_{i,j}^R (\epsilon_{i,j} - \epsilon_{i+1,j}) + B_{i,j}^A (\epsilon_{i,j} - \epsilon_{i,j+1}) \\
& + B_{i,j}^B (\epsilon_{i,j} - \epsilon_{i,j-1}) + A_{i,j} \epsilon_{i,j} = r_{i,j}
\end{aligned} \tag{215}$$

where again, $i_{1,n} \leq i \leq i_{2,n}$. If i is outside that range, then $\epsilon_{i,j} = 0$, so there are not an excessive number of unknowns. The ϵ 's having been found, the E 's are updated by

$$E_{i,j}^{\text{new}} = E_{i,j}^{\text{old}} + \omega \epsilon_{i,j} \tag{216}$$

where ω is an acceleration parameter. The method used for solving Eq. (215) is Gaussian elimination. Rewrite Eq. (215) in the form

$$\sum b_{k,m} e_m = d_k \tag{217}$$

where

$$\left. \begin{aligned}
e_k &= \epsilon_{i,j} \\
d_k &= r_{i,j} \\
b_{k,k} &= A_{i,j} + B_{i,j}^L + B_{i,j}^R + B_{i,j}^A + B_{i,j}^B \\
b_{k,k-1} &= -B_{i,j}^L \quad (i \neq i_{1,n}) \\
b_{k,k+1} &= -B_{i,j}^R \quad (i \neq i_{2,n}) \\
b_{k,k-1} &= -B_{i,j}^B \\
b_{k,k+1} &= -B_{i,j}^A
\end{aligned} \right\} \tag{218}$$

and

$$\begin{aligned} k &= (j - 1) t + i + 1 - i_{1,n} \\ t &= \Delta i_n \end{aligned} \quad (219)$$

As in Oliphant's method, $(b_{k,m})$ is expressed as a product of two triangular factors $(w_{k,m})$ and $(v_{k,m})$. Here, the factors are to have the properties

$$\begin{aligned} w_{k,m} &= 0 \quad \text{unless } k - t \leq m \leq k \\ w_{k,k} &= 1 \\ v_{k,m} &= 0 \quad \text{unless } k \leq m \leq k + t \end{aligned} \quad (220)$$

so that

$$b_{k,m} = \sum_{l=k-t}^m w_{k,l} v_{l,m} \quad \text{for } m < k \quad (221)$$

$$b_{k,m} = \sum_{l=m-t}^k w_{k,l} v_{l,m} + v_{k,m} \quad \text{for } k \leq m \quad (222)$$

The solution of Eq. (222) is straightforward. Suppose that $w_{l,m}$ and $v_{l,m}$ are known for all m when $l < k$. Then Eq. (221) can be used to find $w_{k,m}$,

$$w_{k,m} = \frac{b_{k,m} - \sum_{l=k-t}^{m-1} w_{k,l} v_{l,m}}{v_{m,m}} \quad (223)$$

for m running from $k - t$ up to $k - 1$; and then $v_{k,m}$ can be found from Eq. (223):

$$v_{k,m} = b_{k,m} - \sum_{l=m-t}^{k-1} w_{k,l} v_{l,m} \quad (224)$$

The w 's and v 's are found once for all iterations. There are

$$\sum_{n=1}^N (2t_n + 1) t_{n, j_{\max}}$$

of them -- too many to hold internally for interesting sizes of t_n . With the w 's and v 's known, the unknown e 's are found by solving $w v e = b e = r$ in two steps: $w f = r$, then $v e = f$. Thus,

$$f_k = r_k - \sum_{l=k-t}^{k-1} w_{k,l} f_l \quad (225)$$

which may be evaluated for increasing k , and then

$$e_k = \frac{f_k - \sum_{l=k+1}^{k+t} v_{k,l} e_l}{v_{k,k}} \quad (226)$$

which is evaluated by starting with $k = k_{\max} = j_{\max} t$ and working back down.

The matrices w and v for all steps are packed in the array EF in the FORTRAN program. The matrix elements for step L are assigned to $(EF(KF), KF=LB, LC)$, where $LB=LBASE(L)$ and $LC=LBASE(L+1)-1$.

$w_{k,m}$, for $m < k$, and $v_{k,m}$, for $k \leq m$, are stored in $EF(LB+KF)$, where $KF = (2i+1)(k-1) + m$. The vectors d , e , and f are stored, respectively, in arrays RES , EP , and RES .

THE TESTING PROCEDURE

The three methods described above were exercised on a number of problems. Three of the problems were used for detailed tests reported in the next section.

To describe the tests themselves, it is convenient to use the matrix notation $Bx = s$ of Eq. (188) to designate the system of equations to be solved, in which case the successive estimates of the solution are given by Eq. (191), where the matrix C depends on one or two parameters as well as the type of test. Let

$$d^l = x^{l+1} - x^l = C(s - Bx^l) \quad (227)$$

be the difference of two successive iterates. Then one can use

$$\delta_l = \left(\sum_k (d_k^l)^2 \right)^{\frac{1}{2}} \quad (228)$$

as a measure of the change on iteration .

The ratio

$$\rho_l = \delta_l / \delta_{l+1} \quad (229)$$

is a measure of the rate of convergence of the iteration. But

$$\begin{aligned} d^l &= x^{l+1} - x^l \\ &= y - e^{l+1} - (y - e^l) \\ &= E e^{l-1} - E e^l \\ &= E d^{l-1} = E^l d^0 \end{aligned}$$

in the notation used earlier. Consequently, the sequence of difference vectors d^l can be expected to converge to an eigenvector of the error matrix E , and ρ_l will generally converge to the reciprocal of the modulus of its largest eigenvalue. Asymptotically, at least, $1/\log \rho$ is the number of iterations required to get an e -fold reduction of the error; and if T is the computing time for an iteration, then

*See the section entitled "Common Features of the Method."

$$\phi = \frac{\log \rho}{T} \quad (230)$$

is the rate of convergence in e-folds per unit of time.

TEST RESULTS

The various methods described above were exercised on a number (~10) of test problems. They all performed well on problems containing no zones less than 10 mfp in diameter, and they all converged slowly, if at all, on problems containing many zones of diameter ~0.1 mfp. Three problems of the latter type were used as a basis for a detailed comparison.

Problem 1: $i_{\max} = j_{\max} = 20$
 $r_i - r_{i-1} = 0.1$, for all i ; $z_j - z_{j-1} = 0.1$, for all j ; $\sigma_{i,j} = 1$,
 for i and j .

Problem 2 is the same as Problem 1 except that $\sigma_{i,j} = 0.1$ for all i and j .

Problem 3: $i_{\max} = j_{\max} = 20$.

$$r_i - r_{i-1} = \begin{cases} 0.01 & \text{for } i \leq 10 \\ 0.1 & \text{for } 10 < i \end{cases}$$

$$z_j - z_{j-1} = \begin{cases} 0.01 & \text{for } j \leq 10 \\ 0.1 & \text{for } 10 < j \end{cases}$$

$$\sigma_{i,j} = 0.1, \quad \text{for all } i \text{ and } j$$

The test results are summarized in tables I, II, and III. The parameters that were varied are those described earlier in the sections entitled "Oliphant's Method," "Alternating-Direction-Implicit Method," and "Successive Overrelaxation Method." ADI appears not to perform as well as the other two on easy problems but did quite well on the hardest problem (3). It also seemed to be fairly insensitive to the parameter ω so long as

it fell in the range $5 \leq \omega \leq 10$. Oliphant's method performed very well on easy problems but could not be made to work on Problem 3. (Detailed checking of intermediate results were made to ensure the accuracy of the test of Oliphant's method applied to Problem 3.) The value of ϕ in ADI could be increased by storing rather than regenerating the array of γ 's, which does not change after the first iteration, but Oliphant's method would still converge more rapidly on problems like Problem 1. SOR performs best of all on all problems relative to the criterion of largest ϕ . The value of ϕ is based on the time required to compute an iteration but does not account for the preliminary time spent in calculating the approximate inverse, which is very high for high Δi . Consequently, SOR would not perform well on practical problems that converge in a few (~ 10) iterations. Another difficulty with SOR is its large storage requirement, which would make it necessary to use I/O files on large problems.

It is not entirely clear what conclusions can be drawn from the test results. It can be said that Oliphant's method is about as good as any on problems to which TDRAD can practically be applied. ADI is perhaps a better choice because of its greater reliability and insensitivity to the value of ω . SOR, while it works well on hard problems, requires the choice of the sequences $i_{1,n}$ and $i_{2,n}$ and is slow on easy problems because of the setup time it requires.

One fact stands out: none of the methods perform well on problems that involve streaming through thin regions. It may well be that in two or three dimensions, discrete ordinate methods are more efficient than nonequilibrium diffusion when streaming is important.

TABLE I
TEST OF ADI

Problem	ω	ρ	ϕ
1	0	1.026	0.313
	1	1.057	0.663
	3	1.122	1.30
	5	1.191	2.03
	7.5	1.28	2.9
	10	1.21	2.3
2	1	1.019	0.23
	3	1.022	0.24
	5	1.021	0.25
	7.5	1.022	0.26
	10	1.019	0.23
	12	1.017	0.19
3	1	0.962	-0.45
	3	1.026	0.31
	5	1.043	0.50
	7.5	1.072	0.83
	10	1.138	1.54
	12	1.119	1.34

TABLE II
TEST OF OLIPHANT'S METHOD

Problem	λ	ω	ρ	ϕ			
1	0.2	0.9	1.09	1.73			
		1.0	1.13	2.45			
		1.1	1.22	3.85			
		1.2	1.25	4.37			
		1.3	0.98	-0.35			
	0.3	0.7	1.06	1.14			
		0.8	1.08	1.61			
		0.9	1.13	2.43			
		1.0	1.22	4.02			
		1.1	1.32	5.64			
	0.4	1.2	1.00	0.94			
		0.7	1.07	1.46			
		0.8	1.12	2.25			
		0.9	1.21	3.56			
		1.0	1.38	6.38			
	0.5	1.1	0.96	-0.09			
		0.6	1.06	1.25			
		0.7	1.10	1.95			
		0.8	1.18	3.32			
		0.9	1.34	5.91			
2	0.2	1.0	1.01	0.16			
		1.0	1.01	0.25			
		0.9	1.01	0.26			
		1.0	1.03	0.56			
		1.1	0.82	-3.9			
	0.3	1.0	0.43	-17			
		0.4	0.8	1.009	0.17		
			1.0	1.009	0.18		
			1.2	1.008	0.18		
			-0.8	0.8	1.009	0.17	
	1.0			1.008	0.17		
	1.2	1.007		0.15			
	-0.6	1.0		1.007	0.14		
		-0.4		1.0	1.006	0.13	
			-0.2	1.0	1.006	0.12	
				0	1.0	1.004	0.08

TABLE III
TEST OF SOR

Problem	N	Δi	ω	ρ	ϕ
1	6	5	1.0	1.48	2.7
			1.1	1.67	3.5
			1.2	2.04	4.9
			1.25	2.52	6.4
2	6	5	1.2	1.043	0.29
			1.3	1.053	0.35
			1.4	1.069	0.45
			1.5	1.092	0.57
	3	10	1.2	1.24	0.94
			1.3	1.30	1.2
			1.4	1.41	1.6
			1.5	1.79	2.7
			1.6	1	0
	6	5	1.2	1.039	0.26
			1.3	1.049	0.34
			1.5	1.085	0.58
			1.8	1	0
3	3	10	1.2	1.27	1.12
			1.3	1.33	1.3
			1.5	1.59	2.2
			1.8	1	0

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APPENDIX

LONG2, SHORT2, AND NONEQUILIBRIUM
DIFFUSION ROUTINES

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LONG2 ROUTINES

Subroutine DRAW
Subroutine TRAN2

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Subroutine DRAW

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C      DRAW
C      SUBROUTINE DRAW
C      COMPILED 16NOV66,WBL
C      29SEPT67,PRW
C
C      *****
C      H E C T I C   C O M M O N
C      *****
C
COMMON AMN, AMHILN, AMN, ATOM, BCATAG, BCRTAG, BCLTAG,
1UCRTAG, CAPIN, CAPS, CDUT, CMXK, CMXON, CNDE, CO,
2CUE, CSTOP, CV, CYCLE, DBGPRT, DGREY, DHNU, DIANTP,
3UMIN, UNN, UT, DTBGR, DTC, DTH, DTNA, DTO,
4UTH, DTH1, DTH2, DTUF, DTVF, DUMPT7, ECK, EFRAC,
5EII, ERNCHT, ETH, EZEHO, FDTG, FFA, FFB, FIFT,
6FUDGE, GV, HCB, MCP, MH, MNU, MNUP, MNUS,
7HVB, I, IC, IGOTO, IH, IMNU, IMAX, IMPTAG,
8ISEND, ISR, ITAG, ITRMAX, IU, IV, I1, I2,
9J, JC, JH, JMAX, JMAXA, JU, JV, JS,
COMMON K, KMAXA, KMN1, KMN2, KP, L, M,
1MENGE, MFTAG, MZ, N, NC, NHNU, NK, NPC,
2NH, NRM, NT, NTAPE, NY, NI, N2, N3,
3NH, ODDC, PABOVE, PBLO, PI, PIDTS, PRINTL, PRINTS,
4PHOB, PRK, P, HADE, HC, RFT, RPTAG, RR,
5SCUR, SCH, SCHE, SCYCLE, SGNL, SIG, SIGH, SLUG,
6SN, SPHOB, SVMAX, SVS, S1, S2, S3, S4,
7SS, T, TAUDTS, THICK, TMAX, TRAD, URR, UT,
8UU, UVMAX, VABOVE, VAPE, VBLO, VEL, VS, WSA,
9XMAX,
COMMON AIX(1200), AMX(1200), ALAMH(1200),
1ALAMV(1200), B(1200), CAP(1200), ER(1200),
2FIOUT(1200), K(1200), P(1200), SHLQ(1200),
3THETA(1200), U(1200), V(1200), XALP(1200),
4XB(1200), SOLID(400), PL(200), Y(101),
5UY(100), GAMC(100), YAMC(100), UL(100),
6W2(100), X(53), DX(52), TAU(52),
7BETAB(50), TEMP(12), HEAD(12)
C
C      *****
C      DIMENSION FLEFT(100), PH(100), SIGC(100)
C      EQUIVALENCE (FLEFT, UL), (PL, PR)
C      EQUIVALENCE (PL(101), SIGC)
C      EQUIVALENCE (S20, BLANK(12)), (ZP126, BLANK(13)), (ZP136, BLANK(14)),
C      ( ), (KMX, BLANK(15)), (dT, BLANK(16)), (ANUMBR, BLANK(19)), (LOCZ,
C      , BLANK(22)), (H, BLANK(25)), (NA, BLANK(26))
C      EQUIVALENCE (BACC, BLANK(47)), (SACC, BLANK(48)), (TACC, BLANK(49))
C
C      REAL MENGE
C
C      DIMENSION OLDTH(1), PLANCK(1), ROSS(1), XA(1)
C      DIMENSION RX(50), HPT(100), KXA(1), NER(1)
C      EQUIVALENCE (OLDTH, FIOUT), (PLANCK, P), (ROSS, U), (ER, NER)
C      EQUIVALENCE (KXA, XA, V)
C      COMMON /BLACK/BETAB(50), BETAA(50)
C      COMMON /RADUAT/CLAMDA, ATHETA(50), RTHETA(50), BTHETA(50)

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C   ATHETA, RTHETA, BTHETA CONTAIN BOUNDARY TEMPERATURES (IN EV) FOR THE
C   ABOVE, RIGHT, AND BOTTOM BOUNDARIES RESPECTIVELY
C
    IITOT = 0
    REWIND 4
    WRITE(4) FIOUT, P, U, V
    REWIND 4
    KMAX = KMAXA - 1
    REWIND 3
    DO 10 K = 2, KMAX
10  XALP(K) = 0.
    DO 20 I = 1, IMAX
    BETAA(I) = 0.
20  BETAB(I) = 0.
    DO 25 J = 1, JMAX
25  BETAR(J) = 0.
    SHU = 1.0
    READ (5,960) CLAMUA
    READ (5,960) (ATHETA(I), I = 1, IMAX)
    READ (5,960) (RTHETA(J), J = 1, JMAX)
    READ (5,960) (BTHETA(I), I = 1, IMAX)
    WRITE (6,9003) CLAMUA
    WRITE (6,9004) (ATHETA(I), I = 1, IMAX)
    WRITE (6,9005) (RTHETA(J), J = 1, JMAX)
    WRITE (6,9006) (BTHETA(I), I = 1, IMAX)
    READ 960, ANG6
    NANG = ANG6
    WRITE (6,104J) NANG
30  IF (X(0).LT.1.E-10) GO TO 40
    S1 = 5.0025
    GO TO 955
40  NHUP = 0
    IGWD = 2
C
C   G R I D   L O O P
C
    DO 910 IANG = 1, NANG
    READ 960, EMU, DMU, DZ1, DR1
    WRITE (6,1060) EMU, DMU, DZ1, DR1
    IF (EMU .LT. .99995) GO TO 45
    S1 = 5.0021
    GO TO 955
C   THETA=0 PROHIBITED AS AN INPUT CASE
45  ETA = SQRT(1.0-EMU*EMU)
    SMU = SMU - DMU
    IF (SMU .GE. 0.) GO TO 47
    S1 = 5.0022
    GO TO 955
47  COTAN = EMU/ETA
    DMF = 1.0/ETA
    DZF = 1.0/EMU
    IRX = X(IMAX)/DR1 + 0.8
C   0.0 IS AN ALEXANDERISM
    IF (IRX .LE. 100) GO TO 48
    S1 = 5.0023

```

DRAW 510
DRAW 520
DRAW 530
DRAW 570
DRAW 8054

DRA8056
DRA8057
DRA8058
DRA8061
DRA8062
DRA8063

DRA8073
DRA8074

DRA8076
DRA8077

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      GO TO 955
48  HPT(1) = 0.0
      UJ 50 IIN=2,IRX
      50 HPT(IIN)=RPT(IIN-1)+UR1
      ZXTNA=X(IMAX)*COTAN
C   FIRST SET UP LIKE NEITHER TOP NOR BOTTOM IS REFLECTING
      JZX=(Y(JMAX)-Y(0)+ZXTNA)/UZ1+0.5
      ZSTRT=Y(JMAX)+(0.5-FL0AT(JZX))*UZ1
      ZEND=Y(JMAX)+ZXTNA
      IF(BCBTAG.LT. 0.0) ZSTRT=Y(0)+0.5*UZ1
      IF(BCATAG.LT. 0.0) ZEND=Y(JMAX)
      IF(BCBTAG.GT. -1.E-20 .OR. BCATAG.GT. -1.E-20) GO TO 95
      S1 = 5.0024
      GO TO 955
95  WT = UJU * UZ1 * DR1 * ETA
C
C   C O M B   L O O P
C
      DO 900 IHD=1,IRX
      IF(IHD.GT. 2) GO TO 100
      IF(IHD.EQ. 1) WT = 0.5*WT
      IF(IHD.EQ. 2) WT = 2.0*WT
100  CONTINUE
      PARIM=HPT(IHD)
      IQ=0
      DO 290 I=1,IMAX
      QW=X(I)**2-PARIM**2
      IF(QW.LT.1.E-10) GO TO 290
      IQ=IQ+1
      RX(IQ)=SQRT(QW)
290  CONTINUE
      IF(IQ.EQ.0) GO TO 900
      INSTRT=IMAX-IQ+1
C
C   R A Y   L O O P
C
      ZCNTH=ZSTRT-UZ1
      DO 890 IZHAY=1,1000
      ZCNTH=ZCNTR+UZ1
      IF(ZCNTH.GT. ZEND) GO TO 900
      ITOP=U
      KXA(1)=0
      KXA(201)=0
C   B E G I N   U P W A R D   T R A C E
      IUP=1
      IUPUP=1
      IF(ZCNTR.LT.Y(0)+1.E-5.AND.BCBTAG.LT.0.) GO TO 890
      IF(ZCNTH.LT.Y(0)) GO TO 300
      GO TO 340
C   CENTER OF RAY IS BELOW THE SYSTEM - - INITIALIZE FOR UPWARD TRACE
300  SOELZ=Y(0)-ZCNTH
      IF(COTAN.LT.1.E-10) GO TO 730
      SOELX=SOELZ/COTAN
      IF(SOELX.GT.RX(IQ)) GO TO 730
      IQ=IQ-1

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DHA8159
 DHA8160
 DHA8161
 DHA8162
 DHA8163
 DHA8164
 DHA8165
 DHA8166
 DHA8167
 DHA8168

DHA8171
 DHA8172
 DHA8173

DHA8176
 DHA8177
 DHA8178
 DHA8179
 DHA8180

DHA8183
 DHA8184
 DHA8185
 DHA8186
 DHA8187

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310 IF (100.EQ.0.OR.SDELX.GT.RX(100)) GO TO 320	DHA8188
100=100-1	DHA8189
IF (100.GT.0) GO TO 310	DHA8190
1TH=1HSTRT	DHA8191
GO TO 330	DHA8192
320 1TH=100+1HSTRT	DHA8193
330 JTH=1	DHA8194
2TH=Y(10)	DHA8195
BETAB(1TH) = BETAB(1TH) + WT	
XTH=SDELX	DHA8197
GO TO 380	DHA8198
340 IF (ZCNTH.LT.Y(JMAX)) GO TO 350	DHA8199
ITOP=1	DHA8200
GO TO 640	DHA8201
C CENTER OF RAY IS IN SYSTEM - - INITIALIZE FOR UPWARD TRACE	
350 DO 370 J=1,JMAX	DHA8203
IF (ZCNTH.LT.Y(J)) GO TO 360	DHA8204
GO TO 370	DHA8205
360 ZTH=ZCNTH	DHA8206
XTH=0.	DHA8207
1TH=1HSTRT	DHA8208
JTH=J	DHA8209
GO TO 380	DHA8210
370 CONTINUE	DHA8211
S1=5.0075	DHA8212
GO TO 955	
C G E N E R A L C A S E O F I N C R E A S I N G Z	
380 DELZ=Y(JTH)-ZTH	DHA8215
ICOR = 0	
KTH=1TH+1+IMAX*(JTH-1)	DHA8216
IXX=1TH+1HSTRT+1	DHA8217
IF (IXX.GT.0) GO TO 390	DHA8218
S1=5.008	DHA8219
GO TO 955	
390 DELX=RX(1XX)-XTH	DHA8221
400 IF (DELZ-DELX*COTAN) 410,430,420	DHA8227
C RAY HITS TOP OF CELL	
410 DELS=DELZ*DEF	DHA8229
ZTR=Y(JTH)	DHA8230
JTH=JTH+1	DHA8232
XTH=XTH+DELZ/COTAN	DHA8233
GO TO 440	DHA8234
C RAY HITS SIDE OF CELL	
420 DELS=DELX*DXF	DHA8236
XTH=RX(1XX)	DHA8237
ZTR=ZTH+DELX*COTAN	DHA8238
1TH=1TH+1	DHA8240
GO TO 440	DHA8241
C RAY HITS CORNER OF CELL	
430 DELS=DELX*DXF	DHA8243
XTH=RX(1XX)	DHA8244
ZTR=Y(JTH)	DHA8245
JTH=JTH+1	DHA8246
1TH=1TH+1	DHA8248
ICOR = 1	DHA8249

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C STORE THE PATH LENGTH AND CELL INDEX	
440 KXA(IUP)=KTH	DRA0251
XB(IUP)=DELS	DRA0252
IUP=IUP+1	DRA0253
C CHECK TO SEE IF, AND WHENE, RAY HITS SYSTEM BOUNDARY	DHA0255
IF (ITH.LE.IMAX) GO TO 460	
C RAY HITS SIDE BOUNDARY	
IF (ICOR.EQ.0) GO TO 446	
BETAH(JTR-1) = BETAH(JTH-1) + WT	
GO TO 449	
446 BETAH(JTR) = BETAH(JTH) + WT	
449 IF (IUP.LE.200) GO TO 450	
NXIT=2	DHA0257
GO TO 640	DHA0258
450 NTH=2	DHA0259
GO TO 640	DHA0260
460 IF (JTH.LE.JMAX) GO TO 380	DHA0261
C RAY HITS TOP BOUNDARY	
IF (BCATAG.LT.0.) GO TO 480	DHA0262
BETAH(1TH) = BETAH(1TR) + WT	
IF (IUP.LE.200) GO TO 470	DHA0263
NXIT=3	DHA0264
GO TO 640	DHA0265
470 NTH=3	DHA0266
GO TO 640	DHA0267
C RAY HITS TRANSMITTIVE TOP BOUNDARY	
480 IF (IUP.LE.200) GO TO 490	DRA0268
S1=5.0105	DRA0269
GO TO 955	
C RAY HITS REFLECTIVE TOP BOUNDARY -- INITIALIZE FOR DOWNWARD TRACE	
C GO DOWN IN Z.	DHA0272
490 ZTH=Y(JMAX)	DHA0273
JTH=JMAX	DHA0274
C GENERAL CASE OF DECREASING Z	
ICOR = 0	
500 IXX=ITH-INSTRT+1	DHA0276
MZCRS=1	DHA0277
IF (IXX.GT.0) GO TO 510	DHA0278
S1=5.0112	DHA0279
GO TO 955	
510 DELX=HX(IXX)-XTH	DHA0281
DELZ=ZTH-Y(JTH-1)	DHA0282
KTH=ITH+1+IMAX+(JTH-1)	DHA0283
520 IF (DELZ-DELX*COTAN) 530,550,540	DHA0291
C RAY HITS BOTTOM OF CELL	
530 DELS=DELZ+DXF	DHA0293
ZTH=Y(JTH-1)	DHA0294
JTR=JTH-1	DHA0295
XTH=XTH+DELZ/COTAN	DHA0296
MZCRS=2	DRA0297
GO TO 560	DHA0298
C RAY HITS SIDE OF CELL	
540 DELS=DELX+DXF	DRA0300
XTR=RX(IXX)	DRA0301
ZTR=ZTR-DELX*COTAN	DRA0302

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ITH=ITH+1	DHA8304
GO TO 560	DHA8305
C RAY HITS CORNER OF CELL	
550 UELS=UELX*DXF	DHA8307
XTH=X(IXX)	DHA8308
ZTH=Y(JTH-1)	DHA8309
ALAMV(KTH)=ALAMV(KTH)+1.	DHA8310
JTH=JTH-1	DHA8311
ITH=1/N+1	DHA8312
ICOR = 1	
MZCRS=2	DHA8313
C STORE THE PATH LENGTH AND CELL INDEX	
560 KXA(IUP)=KTH	DHA8315
X(IUP)=UELS	DHA8316
IUP=IUP+1	DHA8317
C CHECK TO SEE IF, AND WHERE, RAY HITS SYSTEM BOUNDARY	
IF (ITH.LE.IMAX) GO TO 580	DHA8319
C RAY HITS SIDE BOUNDARY	
IF (ICOR.EQ. 0) GO TO 566	
BETAR(JTH+1) = BETAR(JTH+1) + WT	
GO TO 569	
566 BETAR(JTH) = BETAR(JTH) + WT	
569 IF (IUP.LE. 200) GO TO 570	
NXIT=2	DHA8321
GO TO 640	DHA8322
570 NTER=2	DHA8323
GO TO 640	DHA8324
580 IF (JTH.LT.1) GO TO 600	DHA8325
GO TO (500+590), MZCRS	DHA8326
590 M=KTH-IMAX	DHA8327
GO TO 500	DHA8329
C RAY HITS BOTTOM BOUNDARY	
600 IF (BCBTAG.LT. 0.) GO TO 620	
C RAY HITS TRANSMITTIVE BOTTOM BOUNDARY	
BETAB(ITH) = BETAB(ITH) + WT	
IF (IUP.LE.200) GO TO 610	DHA8332
NXIT=1	DHA8333
GO TO 640	DHA8334
610 NTER=1	DHA8335
GO TO 640	DHA8336
620 IF (IUP.GT.200) GO TO 630	DHA8337
S1=5.0135	DHA8338
GO TO 955	
C RAY HITS REFLECTIVE BOTTOM BOUNDARY - - INITIALIZE FOR UPWARD TRACE	
630 ZTH=Y(0)	DHA8341
JTH=1	DHA8342
GO TO 380	DHA8343
C BEGIN DOWNWARD TRACE	
640 IF (IUP.GT.200) GO TO 730	DHA8345
IUPUP=IUP-1	DHA8346
IUP=201	DHA8347
IF (ZCNTH.GT.Y(JMAX)-1.E-5.AND.BCATAG.LT.0.) GO TO 890	DHA8348
IF (IUP.GT.0) GO TO 650	DHA8349
GO TO 640	DHA8350
C CENTER OF RAY IS AT TOP OR OUT OF SYSTEM AND NO UPWARD TRACE	

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C	WAS PERFORMED	DHA8352
650	SUELZ=ZCNTH-Y(JMAX)	DHA8353
	IF (COTAN.LT.1.E-10) GO TO 730	DHA8354
	SUELA=SUELZ/COTAN	DHA8355
	IF (SUELX.GT.RX(10)) GO TO 730	DHA8356
	I00=I0-1	DHA8357
660	IF (SUELX.GT.RX(I00)) GO TO 670	DHA8358
	I00=I00-1	DHA8359
	IF (I00.GT.0) GO TO 660	DHA8360
	ITR=INSTNT	DHA8361
	GO TO 680	DHA8362
670	ITH=I00+INSTRT	DHA8363
680	JTH=JMAX	DHA8364
	ZTH=Y(JMAX)	DHA8365
	XTR=SUELX	DHA8366
	KTH=ITH+1+IMAX*(JTH-1)	
	BETAA(ITH) = BETAA(ITR) + WT	
	GO TO 500	DHA8368
690	IF (ZCNTH.GT.Y(U)) GO TO 700	DHA8369
	GO TO 730	DHA8370
C	CENTER OF RAY IS INSIDE SYSTEM AND AN UPWARD TRACE WAS PERFORMED -	
C	- INITIALIZE FOR DOWNWARD TRACE	
700	DO 720 J=1,JMAX	DHA8372
	JH=JMAX+1-J	DHA8373
	IF (ZCNTH.GT.Y(JH-1)) GO TO 710	DHA8374
	GO TO 720	DHA8375
710	ZTH=ZCNTH	DHA8376
	XTH=0.	DHA8377
	ITR=INSTNT	DHA8378
	JTH=JH	DHA8379
	GO TO 500	DHA8380
720	CONTINUE	DHA8381
	S1=5.0160	DHA8382
	GO TO 955	
C	TRACE OF A RAY IS COMPLETE	
C	(EXCEPT THETA=0 TRACE)	
C		
C	MERGE STARTING SEGMENTS IF IN SAME CELL	
730	IF (KXA(1).LE.0.AND.KXA(201).LE.0) GO TO 890	DHA8386
	IF (KXA(1).EQ.KXA(201)) GO TO 740	DHA8387
	IDNMN=201	DHA8388
	GO TO 750	DHA8389
740	IDNMN=202	DHA8390
	X0(1)=XB(1)+XB(201)	DHA8391
C	CONTRIBUTE TO SUMMED PATH LENGTHS FOR EACH CELL	
750	L1=1	DHA8393
	L2=IUPUP	DHA8394
	L3=1	DHA8395
	IUP=IUP-1	DHA8396
760	DO 770 I1=L1,L2	DHA8397
	K=KXA(I1)	DHA8398
770	XALP(K) = XALP(K) + X0(I1)*WT	
	GO TO (780,790), L3	DHA8400

780 IF (IUNMN.GT.IUP) GO TO 790	DHA8401
L1=IUNMN	DHA8402
L2=IUP	DHA8403
L3=2	DHA8404
GO TO 760	DHA8405
C TURN THE UPWARD DATA AROUND	
790 IUPM=IUPUP/2	DHA8407
IF (IUPM.LE.0) GO TO 810	DHA8408
DO 800 I=1,IUPM	DHA8409
IB=IUPUP+1-I	DHA8410
KX1=KXA(I)	DHA8411
XX2=XB(I)	DHA8412
KXA(I)=KXA(IB)	DHA8413
XB(I)=XB(IB)	DHA8414
KXA(IB)=KX1	DHA8415
800 XB(IB)=XX2	DHA8416
810 IF (IUPUP.LT.1.AND.IUP.LT.IDNMN) GO TO 890	DHA8417
IF (IUP.LT.IDNMN) GO TO 830	DHA8418
C MOVE DOWNWARD DATA SNUG NEXT TO UPWARD DATA	
DO 820 I=IDNMN,IUP	DHA8420
IN=IUPUP+1-IDNMN+I	DHA8421
KXA(IN)=KXA(I)	DHA8422
820 XB(IN)=XB(I)	DHA8423
II2=IUPUP+IUP-IDNMN+1	DHA8424
ITOT=II2	DHA8425
IF (IUPUP.LT.1) NTER=3	DHA8426
GO TO 840	DHA8427
830 II2=IUPUP	DHA8428
IF (ZCNTR .LT. T(0)) NXIT =1	
ITOT=IUPUP	DHA8430
840 NPOINT = 1	
GO TO 2359	
C DEBUG EDIT	
880 IF (ABS(UBGPRT) .LT. 1.E-20) GO TO 890	
WRITE(6,970)EMU,IRD,IZHAY,NTER,NXIT,WT	
WRITE (6,980) (KXA(I),XB(I),I=1,II2)	DHA8453
890 CONTINUE	DHA8454
C END OF RAY LOOP	
900 CONTINUE	DHA8455
C END OF COMB LOOP	
WRITE (6,9001) ITOT	
910 CONTINUE	DHA8456
C END OF GRID LOOP	
C	
C THETA = 0 TRACE	
C	
2349 WTP = SMU * 0.5	
WRITE (6,9007) SMU	
DO 2351 J=1,JMAX	
XB(J) = UB(J)	
KXA(J) = 1 + IMAX * (J-1)	
C I = 0 IN ABOVE STATEMENT	
DO 2350 I=1,IMAX	
K = 1 + 1 + IMAX * (J-1)	
WT = WTP * TAU(I)	

```

C  NOTE -- FINAL SUM OF LENGTHS FOR REFLECTING AND NONREFLECTING CASE
C  SHOULD BE EQUAL IN ACCORDANCE WITH NO. OF PASSES IN TRAN2
2350 XALP(K) = XALP(K) + XB(J)*WT
2351 CONTINUE
      IF (BCATAG.LT.0.) GO TO 2352
      IF (BCBTAG.LT.0.) GO TO 2354
C  NO REFLECTION
      ITOT = JMAX
      NTER = 1
      NXIT = 3
      DO 2350 I=1,IMAX
        WT = WTP * TAU(I)
        BETAA(I) = BETAA(I) + WT
2350  BETAB(I) = BETAB(I) + WT
      GO TO 2357
C  TOP REFLECTING
2352 IEND = JMAX - 1
      XB(JMAX) = XB(JMAX) + XB(JMAX)
      DO 2353 J=1,IEND
        XB(JMAX+J) = XB(JMAX-J)
2353  KXA(JMAX+J) = KXA(JMAX-J)
      ITOT = JMAX + IEND
      NTER = 1
      NXIT = -1
      DO 2351 I=1,IMAX
        WT = WTP * TAU(I)
2351  BETAB(I) = BETAB(I) + WT * 2.
      GO TO 2357
C  BOTTOM REFLECTING
2354 IEND = JMAX - 1
      DO 2355 J=1,IEND
        XB(JMAX+J) = XB(J+1)
2355  KXA(JMAX+J) = KXA(J+1)
      XB(JMAX) = XB(1) + XB(1)
      KXA(JMAX) = 1
C  I=0 IN ABOVE STATEMENT
      DO 2356 J=1,IEND
        XB(JMAX-J) = XB(JMAX+J)
2356  KXA(JMAX-J) = KXA(JMAX+J)
      ITOT = JMAX + IEND
      NTER = 3
      NXIT = -1
      DO 2352 I=1,IMAX
        WT = WTP * TAU(I)
2352  BETAA(I) = BETAA(I) + WT * 2.
2357 I12 = ITOT
C  COMPUTE KXA(I) AND PACK RAY DATA FOR ITH COMB (WHICH CONTAINS ONLY
C  ONE RAY) IN THE THETA=0 CASE
      IM = 0
2370 IM = IM + 1
      DO 2358 J=1,ITOT
2358  KXA(J) = KXA(J) + 1
      WT = WTP * TAU(IM)
      NPOINT = 2
      GO TO 2359

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C  DEBUG EDIT
911  IF (AUS(DBGPRT).LT.1.E-20) GO TO 912
      WRITE (6,9008) IM,NTEN,NKIT,WT
      WRITE (6,980) (KXA(I),XB(I),I=1,I12)
912  IF (IM.LT.IMAX) GO TO 2370
      WRITE (6,9001) IITOT
C  WRITE END RECORD ON UATA TAPE
      NEH(IGWU+1)=1
      NEH(IGWU+2)=1
      NEH(IGWU+3)=1
      ER(IGWU+4) = -10.
      NEH(IGWU+5) = -10
      ER(IGWU+6) = -10.
      NEH(2)=NMGP+1
      IGWD = IGWD + 6
      NEH(1)=IGWD
      WRITE (3) (ER(I),I=1,IGWD)
      END FILE 3
      REWIND 3
C  CALCULATE EFFECTIVE AREAS (XALP)
      IWARN=1
      DO 950 K=2,KMAX
      I=MOD(K-2,IMAX)+1
      J=(K-2)/IMAX+1
      IF (XALP(K).LT.1.E-10) GO TO 920
      XALP(K) = TAU(I) * UY(J) * 2. * PI/XALP(K)
C  = EFFECTIVE AREA OF CELL K
C  = 4*PI*VOLUME OF CELL / SUM(LENGTH * WEIGHT OF RAY)
C  NOTE -- XALP WAS HALF THE SUM OF LENGTHS SINCE EACH RAY TRACE IN
C  DNAM ACCOUNTS FOR TWO RAYS IN SYSTEM
      GO TO 950
920  XALP(K)=1.E20
      GO TO (930,940), IWARN
930  IWARN=2
      WRITE (6,1050)
940  WRITE (6,980) K
950  CONTINUE
C  FINIALIZE VARIABLES AND PRINT DEBUG IF CALLED FOR
      READ(4) FIOUT,P,U,V
      REWIND 4
      IF (AUS(DBGPRT) .LT. 1.E-20) GO TO 953
      WRITE (6,990)
      WRITE (6,1000) (XALP(K),K=2,KMAX)
953  CONTINUE
      RETURN
886  S1 = 5.1111
955  ISENG =2
      READ(4) FIOUT,P,U,V
      REWIND 4
      CALL EDIT
C
C  PACK RAY DATA
C
2359  NWD = I12 + I12 + 4
      IITOT = IITOT + ITOT

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DRA8458
DRA8459
DRA8460

DRA8463

DRA8465

DRA8467
DRA8468

DRA8470
DRA8471
DRA8472
DRA8473
DRA8474

DRA8476
DRA8477

DRA8479
DRA8480
DRA8481
DRA8482

DRA8487
DRA8488

DRA8495

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      IF (IGWD + NND .GT. 1023) GO TO 870
850  NER(IGWD+1)=ITOT
      NER(IGWD+2)=ENTER
      NER(IGWD+3)=EXIT
      ER(IGWD+4) = WT
      DO 860 I=1,112
      ITWD = IGWD + I + I + 3
      NER(ITWD)=KXA(I)
860  ER(ITWD+1)=XB(I)
      IGWD=IGWD+NND
      NNGP=NNGP+1
      GO TO 885
870  NER(1)=IGWD
      NER(2)=NNGP
      WRITE (3) (ER(I),I=1,IGWD)
      IGWD=2
      NNGP=0
      GO TO 850
885  GO TO (880,911), NPOINT
C
C
C
960  FORMAT (PE12.5)
970  FORMAT (/6H MU = ,F10.4,6X,8H R INDEX,15,6X,8H Z INDEX,15,6X,9H BC
      1 ENTER,15,6X,8H BC EXIT,15,4X,3H WT,1PE12.5)
980  FORMAT (6(15,1PE15.8))
990  FORMAT (6H1AREAS)
1000 FORMAT (1P8E15.8)
1010 FORMAT (50HICENSUS OF RAYS CROSSING THE OUTER SYSTEM BOUNDARY)
1020 FORMAT (40HICENSUS OF RAYS CROSSING THE TOP SYSTEM BOUNDARY)
1030 FORMAT (51HICENSUS OF RAYS CROSSING THE BOTTOM SYSTEM BOUNDARY)
1040 FORMAT (11H0//15HGRID INPUT FOR ,12,24H GRIDS AND SEGMENT COUNT/)
1050 FORMAT (42H WARNING. NO RAYS GO THROUGH THESE ZONES )
1060 FORMAT (1H ,5X,4H MU = ,1PE12.5,13H DELTAMU = ,E12.5,10H DZ1=
      1,E12.5,10H DR1= ,E12.5)
1070 FORMAT (13H TOTAL CROSSINGS FOR OUTER =,F10.2)
1080 FORMAT (13H TOTAL CROSSINGS FOR BOTTOM =,F10.2)
1090 FORMAT (13H1TOTAL NUMBER OF RAYS EMPLOYED =,F10.2)
1100 FORMAT (13H TOTAL CROSSINGS FOR TOP =,F10.2)
9001 FORMAT (1H ,5X,47HAFTER COMPLETION OF ABOVE GRID, TOTAL SEGMENTS=,
      .110//)
9003 FORMAT (11H0,10X,8HCLAMDA= ,F12.2)
9004 FORMAT (11H0,10X,9HATHETA(1)/(1P7E15.5))
9005 FORMAT (11H0,10X,9HATHETA(1)/(1P7E15.5))
9006 FORMAT (11H0,10X,9HATHETA(1)/(1P7E15.5))
9007 FORMAT (1H ,5X,8H MU = 1.,13X,8HDELTAMU = ,E12.5)
9008 FORMAT (/12H MU = 1.,7X,7H INDEX,15,7X,12H2 INDEX -1,6X,9H
      1BC ENTER,15,6X,8H BC EXIT,15,4X,3H WT,1PE12.5)
      END

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DRA8434
 DRA8435
 DRA8436
 DRA8437
 DRA8439
 DRA8440
 DRA8441
 DRA8442
 DRA8445
 DRA8446
 DRA8448
 DRA8449
 DRA8450

DRA8496

DRA8500
 DRA8501
 DRA8502

DRA8507

DRA8508

Subroutine TRAN2

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C      TRAN2
C      SUBROUTINE TRAN2
C      COMPILED 16NOV66, bbl
C
C      29SEPT67, PRW
C
C      H E C T I C   C O M M O N
C      *****
C      COMMON AIN, AMBIEN, ANN, ATUM, BCATAG, BCBTAG, BCLTAG,
1BCRTAG, CAPIN, CAPS, CDUT, CMXK, CMXOM, CNDE, CO,
2COE, CSTOP, CV, CYCLE, DBGPRT, DGREY, DHNU, DIANTP,
3DMIN, DNN, DT, DTBUGR, DTC, DTH, DTNA, DTG,
4DTR, DTH1, DTH2, DTUF, DTVF, DUMPT7, ECK, EFRAC,
5E11, ERRCRT, ETH, EZERO, FDTG, FFA, FFB, FIFT,
6FUDGE, EV, HCB, MCP, HM, HNU, HNUF, HNUS,
7HVB, I, IC, IGOTO, IH, IHNU, IMAX, IMPTAG,
8ISENU, ISR, ITAG, ITRMAX, IU, IV, I1, I2,
9J, JC, JH, JMAX, JMAXA, JU, JV, J5,
COMMON K, KMAX, KMNI, KMNI2, KP, L, M,
1MERGE, MFTAG, MZ, N, NC, NHNU, NK, NPC,
2NR, NRM, NT, NTAPE, NY, NI, N2, N3,
3N4, ODDC, PAUOVE, PBLO, PI, PIOTS, PRINTL, PRINTS,
4PROB, PHR, PH, RADE, RC, RFT, RPTAG, RR,
5SCDR, SCR, SCHE, SCYCLE, SGNL, SIG, SIGH, SLUG,
6SH, SPROB, SVMAX, SVS, S1, S2, S3, S4,
7SS, T, TAUOTS, THICK, TMAX, TRAD, URR, UT,
8UU, UVMAX, VABOVE, VAPE, VBLO, VEL, WS, WSA,
9XMAX,
COMMON
1ALAMV(1200), AIX(1200), AMX(1200), ALAMH(1200),
2FIOUT(1200), B(1200), CAP(1200), ER(1200),
3JHETA(1200), KF1T(1200), P(1200), SMLQ(1200),
4XB(1200), SOLID(400), PL(200), XALP(1200),
5DY(100), GAMC(100), YAMC(100), Y(101),
6W2(100), X(53), DX(52), UL(100),
7BETAB(50), TEMP(12), HEAD(12), TAU(52),
C
C      *****
C      DIMENSION FLEFT(100), PR(100), SIGC(100)
C      EQUIVALENCE (FLEFT,UL), (PL,PR)
C      EQUIVALENCE (PL(101),SIGC)
C      EQUIVALENCE (S20,BLANK(12)), (ZP126,BLANK(13)), (ZP136,BLANK(14)),
1(KMX,BLANK(15)), (WT,BLANK(16)), (ANUMBR,BLANK(19)), (LOC2,BLANK(22))
2(H,BLANK(25)), (NA,BLANK(26))
C      EQUIVALENCE (BACC,BLANK(47)), (SACC,BLANK(48)), (TACC,BLANK(49))
C
C      REAL MERGE
C
C      DIMENSION PUR(50), PUZ(100), RUR(50), RUZ(100)
C      DIMENSION OLDTH(1), PLANCK(1), ROSS(1), XA(1)
C      DIMENSION KXA(1)
C      DIMENSION KXB(1)
C      EQUIVALENCE (OLDTH,FIOUT), (PLANCK,P), (ROSS,U),
1(KXA,XA,V), (KXB,XB)
COMMON /BLACK/ BETAR(50), BETAA(50)

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TRN1000
TRN1001

TRN1003
TRN1004
TRN1005
TRN1006

HECCOM22
HECCOM23
HECCOM24
HECCOM25
HECCOM26
HECCOM27
HECCOM28

TRN1034
TRN1035
TRN1036
TRN1037

TRN1041
TRN1042
TRN1043
TRN1044
TRN1045

TRN1047
TRN1048
TRN1050


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      DIMENSION PA(50),PK(50),PB(50)
      COMMON /RADUAT/ CLAMDA,ATHETA(50),RTHETA(50),BTHETA(50)
C   ATHETA,RTHETA,BTHETA CONTAIN BOUNDARY TEMPERATURES (IN EV) FOR THE
C   ABOVE,RIGHT,AND BOTTOM BOUNDARIES RESPECTIVELY
C   *****
C   XALP, ALAMV, ALAMH, AND BETAB INVIOLEATE -- FORMED IN DRAW
C   NY(IN LINDLEY COMMON) IS TEMPERATURE ITERATION INDEX USED IN KAPPA
      FOO=0.0
      ETOP=0.0
      ESIDE=0.0
      EBTM=0.0
      KMAX=KMAXA-1
      IF (MFTAG.EQ.0) DHNU=1.
      NVEZ=1
      NY=NVLZ
      IF (ITAG.EQ.0) NVEZ=2
      VEZ=I,VEZ
      CALL DVCHK (KDMY)
      THTAMX=.025
C   CALCULATE GEOMETRY FACTORS AND FIND HIGHEST TEMPERATURE
      DO 10 I=1,IMAX
      PUR(I)=PI/TAU(I)
10  KUR(I+1)=4.0E10*(X(I)/(X(I+1)-X(I-1)))
      DO 20 J=1,JMAX
      PUZ(J)=1./UY(J)
20  KUZ(J+1)=2.0E10/(Y(J+1)-Y(J-1))
      DO 30 K=2,KMAX
      IF (JMK(KFIT(K),2).NE.1) GO TO 30
      IF (THETA(K).LT.THTAMX) GO TO 30
      THTAMX=THETA(K)
30  CONTINUE
      THTAMX=AMAX1(THTAMX,BCLTAG,BCRTAG,BCBTAG,BCATAG)
C   RECENTHY POINT FOR SECOND TEMPERATURE ITERATION
40  DO 50 K=2,KMAX
      XA(K)=0.
      XB(K)=0.
50  ER(K)=0.0
C
C
C   *****
C   DO FREQUENCY BOOKKEEPING
C   *****
C   SET UP MAX FREQ BOUNDARY
      HNUP=1.0E6
      HNUP4=1.0E24
      IF (MFTAG.EQ.0) GO TO 180
60  IHNU=IHNU+1
      CALL KAPPA
      HNUS=HNUP**4
      DHNU=HNUP-HNU
C

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      THN1051
      THN1052
      THN1053
      THN1054
      THN1055
      THN1056
      THN1057
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      THN1060
      THN1061
      THN1062
      THN1063
      THN1064
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      THN1079
      THN1080
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      THN1095
      THN1096
      THN1097
      THN1098
      THN1099
      THN1100
      THN1101

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C	MERGE GROUPS WITH HNU MORE THAN (MERGE) TIMES LARGEST THETA	TRN1102
C	IF (MERGE.GT.0.0) GO TO 70	TRN1103
	S1=7.0225	TRN1104
	GO TO 965	TRN1105
70	IF (THTAMX-HNU/MERGE) 100.80.80	TRN1107
80	IF (HNU-1) 90.210.120	TRN1108
90	S1=7.0235	TRN1109
	GO TO 965	
C		TRN1111
C	REJECT TAPE IF MORE THAN HALF OF GROUPS MERGE	TRN1112
C		TRN1113
100	IF (HNU+HNU-NHNU) 120.110.110	TRN1114
110	Q=AMOD(MERGE,1.)	TRN1115
	IF (Q.GT.0.4.AND.Q.LT.0.6) GO TO 120	TRN1116
	S1=7.0250	TRN1117
	GO TO 965	
120	DO 130 K=2,KMAX	TRN1119
	IF (JMK(KFIT(K)-2).NE.1) GO TO 130	TRN1120
	T4=THETA(K)**4	TRN1121
	BETA=HNU/THETA(K)	TRN1122
	BETAP=HNUP/THETA(K)	TRN1123
	DFB=PLNKUT(BETA*BETAP)	TRN1124
	IF (DFB.LT.1.E-10) GO TO 130	TRN1125
	TEMP(1)=DFB*T4	TRN1126
	EMB1=EXP(-BETA)	TRN1127
	EMB2=EXP(-BETAP)	TRN1128
	TEMP(2)=DFB+0.0384974/T4*(HNU4/(1.0-EMB1)*EMB1-HNUP4/(1.0-EMB2)*EMB2)	TRN1129
	182)	TRN1130
C		TRN1131
C	FORM NUMERATORS AND DENOMINATORS OF MERGED KAPPAS	TRN1132
C		TRN1133
	XA(K)=XA(K)+TEMP(1)	TRN1134
	XB(K)=XB(K)+TEMP(2)	TRN1135
	B(K)=B(K)+PLANCK(K)+TEMP(1)	TRN1136
	ER(K)=ER(K)+TEMP(2)/ROSS(K)	TRN1137
130	CONTINUE	TRN1138
	HNUP=HNU	TRN1139
	HNUP4=HNU4	TRN1140
	IF (THTAMX-HNU/MERGE) 60.140.140	TRN1141
C		TRN1142
C	FORM MERGED KAPPAS	TRN1143
C		TRN1144
140	DO 170 K=2,KMAX	TRN1145
	IF (XA(K)) 150.170.160	TRN1146
150	S1=7.0320	TRN1147
	GO TO 965	
160	ROSS(K)=XB(K)/ER(K)	TRN1149
	PLANCK(K)=B(K)/XA(K)	TRN1150
	ER(K)=0.	TRN1151
170	CONTINUE	TRN1152
	HNUP=1.0E6	TRN1153
	HNUP4=1.0E24	TRN1154
	OHNU=HNUP-HNU	TRN1155
	GO TO 230	TRN1156

C		TRN1157
C	MONOFREQUENCY CALCULATION	TRN1158
C		TRN1159
	180 MNNU=1	TRN1160
	IMNU=1	TRN1161
	DO 190 K=2,KMAX	TRN1162
	IF (JMN(KFIT(K),2).NE.1) GO TO 190	TRN1163
	B(K) = 3.2732E+11 * THETA(K)*THETA(K)*THETA(K)*THETA(K)	
C	3.2732E+11 = 1.0283E+12/PI = STEFANS CONSTANT (EV)/PI	
C	B = 5/(4*PI*H055)	
	190 CONTINUE	TRN1165
	MNU=.001	TRN1166
	GO TO 230	TRN1167
C	*****	TRN1168
C	BEGIN FREQUENCY LOOP	TRN1169
C	*****	TRN1170
C	*****	TRN1171
C	*****	TRN1172
	200 IMNU=IMNU+1	TRN1173
	CALL KAPPA	TRN1174
	DMNU=MNU*HNU	TRN1175
	MNU4=MNU**4	TRN1176
C		TRN1177
C	TYPICAL GROUP CALCULATION OF SOURCES	TRN1178
C		TRN1179
	210 DO 220 K=2,KMAX	TRN1180
	IF (JMN(KFIT(K),2).NE.1) GO TO 220	TRN1181
	DFB=PLNKUT(MNU/THETA(K),MNU/THETA(K))	TRN1182
	B(K) = DFB * 3.2732E+11 * THETA(K)*THETA(K)*THETA(K)*THETA(K)	
	220 CONTINUE	TRN1184
C		TRN1185
C	SET BLACKBODY CONDITIONS	TRN1186
C		TRN1187
	230 IF (BCATAG .LT. 1.E-20) GO TO 240	
	DO 235 I=1,IMAX	
	IF (ATHETA(I) .GT. 1.E-5) GO TO 234	
	PA(I) = 0.	
	GO TO 235	
	234 PA(I) = PLNKUT(MNU/ATHETA(I),MNU/ATHETA(I))	
	PA(I) = PA(I)*TAU(I)/BETAA(I)*1.0283E12 * ATHETA(I)*ATHETA(I)	
	I	
	235 CONTINUE	
	240 IF (BCHTAG .LT. 1.E-20) GO TO 250	
	DO 245 J=1,JMAX	
	IF (RTHETA(J) .GT. 1.E-5) GO TO 244	
	PHI(J) = 0.	
	GO TO 245	
	244 PHI(J) = PLNKUT(MNU/RTHETA(J),MNU/RTHETA(J))	
	PHI(J) = PHI(J)*PI*2.*X(IMAX)*DY(J)/BETAR(J) * 1.0283E12	
	1 * RTHETA(J)*RTHETA(J)*RTHETA(J)*RTHETA(J)	
	245 CONTINUE	
	250 IF (BCBTAG .LT. 1.E-20) GO TO 270	
	DO 255 I=1,IMAX	
	IF (BTHETA(I) .GT. 1.E-5) GO TO 254	
	PB(I) = 0.	

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      GO TO 255
254  P5(I) = PLNKUT(HNU/BTHETA(I),HNU/BTHETA(I))
      P5(I) = P5(I)+TAU(I)/BETAB(I)+1.0283E12 * BTHETA(I)+BTHETA(I)
      * BTHETA(I)+BTHETA(I)
1
255  CONTINUE
270  CALL DVCHK (KDMY)
      GO TO (280,290), KDMY
280  S1=7.0522
      GO TO 965
C
C      FORM NOSSLAND AND PLANCK OPTICAL DEPTHS
C
C      DOUBLE ON STORAGE FOR ABSORPTION COEFFICIENTS, MU, AND LAMBDA
290  DO 300 I=1,IMAX
      W2(I)=0.
      K=1
      M=K+1
      DO 300 J=1,JMAX
      FACTOR=1.
      IF (ABS(HPTAG),67,1,1-20) FACTOR=NOSS(K)/PLANCK(K)
      NOSS(K)=AMAX1(NOSS(K)+AMX(K)/(TAU(I)+DY(J)),1.E-20)
      PLANCK(K)=AMAX1(PLANCK(K)+FACTOR*AMX(K)/(TAU(I)+DY(J)),1.E-20)
C      USE ONLY ROSS FOR NOW. ALAMH AND ALAMV USED FOR RAY COUNT
C      USING PLANCK AS ENERGY FLUX STORAGE IN RAY LOOP FOR DEBUG AND W2
      M=M+1
      K=K+1
300  CONTINUE
C      *****
C      TWO-DIMENSIONAL TRANSPORT SECTION *****
C
C      *****
      DO 306 J=1,JMAX
      DO 305 I=1,IMAX
      K = I + 1 + IMAX*(J-1)
      ALAMV(K) = 0.
      IF ( CLAMUA/ROSS(K) .LT. AMIN1(DY(J),DX(I)) .AND.
1      ( JMR(KFIT(K),2).EQ.1) ) ALAMV(K)=1.
305  CONTINUE
306  CONTINUE
C      ALAMV(K)=1. IMPLIES CELL K IS OPTICALLY THICK AND ACTIVE
      NEWIND 3
C      KXB, XB IS -BUFFER- STORAGE
C      KXA, XA USED FOR BOTH CELL INDEX AND PATH LENGTH(TRAN2 ONLY)
C
C      RAY - TAPE RECORD LOOP (NGRP RAYS)
C
310  READ (3) NGWD,(XB(I),I=2,NGWD)
      NGRP=KXB(2)
      IGWD=2
      IF (IAUS(NGRP).LE.210) GO TO 320
      S1=7.0541
      GO TO 965
C
C      RAY LOOP

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 THN1207

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 THN1230
 THN1231

THN1232
 THN1233
 THN1234

TRN1236
 TRN1237
 TRN1238
 TRN1239

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C 320 UU 820 IGHP=1,NGHP
      ITOT=KXB(IGWU+1)
      NTER=KXB(IGWU+2)
      NAIT=KXB(IGWU+3)
      WT = XB(IGWU+4)
      I1 = IGWU + 4
      UU 330 I=1,ITOT
      KXA(I)=KXB(I1+1)
      XA(I+400)=XB(I1+2)
      I1=I1+2
      IGWU=I1
C 330 TEST FOR -ALL WAYS PROCESSED- SIGNAL
      IF (ITOT.EQ.1.AND.KXA(1).LT.0) GO TO 830
C 340 BYPASS IF WAY TRAVERSES TOTALLY INACTIVE REGION
      IACT=1
      UU 340 I=1,ITOT
      KI=KXA(I)
      IF (JMN(KFIT(KT),2).NE.1) GO TO 340
      IACT=2
      340 CONTINUE
      GO TO (720,350), IACT
C 350 GO FORWARD
      IF (NTER.EQ.1.AND.NTER.LE.3) GO TO 360
      S1=7.0550
      GO TO 365
      360 KSTRT=KXA(1)
      IPASS=1
C 370 APPLY BOUNDARY CONDITION TO START TRACE OF RAY
      UU TO (370,400,420), NTER
      370 IF (BCHTAG) 380,440,390
      380 S1=7.0556
      GO TO 365
      390 I=MOD(KSTRT-2,IMAX)+1
      AJ = PB(1) * WT
      GO TO 450
      400 IF (BCHTAG) 380,440,410
      410 J=(KSTRT-2)/IMAX+1
      AJ = PB(J) * WT
      GO TO 450
      420 IF (BCHTAG) 380,440,430
      430 I = MOD(KSTRT-2,IMAX) + 1
      AJ = PA(1) * WT
      GO TO 450
      440 AJ=0.
C 450 BEGIN LOOP OVER ZONES PASSED THROUGH BY RAY
C 460 FOR DEBUG ONLY, USE PLANK ARRAT TO STORE AJ
C 470 XALP IS EFFECTIVE AREA OF RAY IN ZONE
      PLANK(1) = AJ
      KOLD = 1
      ALAMP(KOLD) = 0.
C 480 OUTSIDE OF SYSTEM IS CONSIDERED OPTICALLY THIN
C 490
C 500 S E G M E N T   L O O P
C 510

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 THN1289


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DO 540 IUY=1,ITOT
GO TO (490,500), IPASS
490 ITR=IUY
GO TO 510
500 ITH=ITOT+1-IUY
510 K=KXA(ITH)
520 IF (ALAMV(KOLD) .GT. 0.5) GO TO 525
C
C PREVIOUS ZONE WAS THIN
C
IF (ALAMV(K) .LT. 0.5) GO TO 528
C PRESENT ZONE IS THICK
ER(K) = ER(K) + AJ
AJN = AJ
IF (IUY.LT.ITOT) GO TO 530
AJN = U(K)*XALP(K)*WT
ER(K) = ER(K) - AJN
GO TO 530
C
C PREVIOUS ZONE WAS THICK
C
525 IF (ALAMV(K) .LT. 0.5) GO TO 527
C PRESENT ZONE IS THICK
IF (IUY .LT. ITOT) GO TO 529
AJN = U(K)*XALP(K)*WT
ER(K) = ER(K) - AJN
GO TO 530
C PRESENT ZONE IS THIN
527 AJ = U(KOLD)*XALP(KOLD)*WT
ER(KOLD) = ER(KOLD) - AJ
528 IF (JMK(KFIT(K),2).EQ.0) GO TO 529
C NORMAL TRANSPORT CALCULATION
SIGS = MUSS(K)*XA(ITH+400)
ESS = EXP(-SIGS)
AJN = AJ + ESS + XALP(K) * U(K) * (1.-ESS) * WT
ER(K) = ER(K) + AJ - AJN
C
LN = ERGS/SEC XALP = EFFECTIVE AREA
GO TO 530
C PRESENT ZONE IS INACTIVE OR THICK
529 AJN = AJ
530 PLANCK(IUY+1) = AJN
KOLD = K
540 AJ=AJN
C IF PLANCK DEBUG IS DESIRED IT SHOULD BE ADDED HERE
C ON SECOND PASS ORDER OF PLANCK PRINT WILL BE SAME AS TRACE DIRECTION
C F I N I S H S E G M E N T L O O P
C H A N G E E T H B Y E N E R G Y E N T E R I N G A N D L E A V I N G M E S H
550 F00 = F00 + (PLANCK(1) - PLANCK(ITOT+1))
C EVALUATE CONTRIBUTION TO FLUX ACROSS TOP BOUNDARY
IF (BCATAG.LT.0.) GO TO 610
IF (INTER.NE.3) GO TO 580
I=MOD(KXA(1)-2,IMAX)+1
C IN ERGS/CM**2/SEC
GO TO (560,570), IPASS
560 W2(I)=W2(I)-PLANCK(I)/TAU(I)

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TRN1297
 TRN1298
 TRN1299
 TRN1301
 TRN1302
 TRN1304

TRN1317

TRN1322

TRN1324

TRN1325

TRN1326

TRN1327

TRN1329

GO TO 580	TRN1331
570 W2(1)=W2(1)+PLANCK(1)*TOT+1/TAU(1)	TRN1333
580 IF (NEXIT.LE.3) GO TO 610	TRN1334
I=MOD(KXA(TOTI-2,IMAX)+1	
IN EGUS/CM**2/SEC	
GO TO (590,600), IPASS	TRN1336
590 W2(1)=W2(1)+PLANCK(1)*TOT+1/TAU(1)	TRN1338
GO TO 610	
600 W2(1)=W2(1)-PLANCK(1)/TAU(1)	TRN1340
610 GO TO (620,720), IPASS	
620 CONTINUE	
630 KSTMT=KXA(TOTI	TRN1345
GO TO (640,650,660), NTER	TRN1346
640 EUTM=EUTM+PLANCK(1)	TRN1347
GO TO 670	TRN1348
650 ESIDE=ESIDE+PLANCK(1)	TRN1349
GO TO 670	TRN1350
660 ETOP=ETOP+PLANCK(1)	TRN1351
670 IF (NEXIT.GT.0) GO TO (680,690,700), NEXIT	
GO TO (680,690,700), NTER	
C THICKEN CASE AND REFLECTION -- RAY ENTERS AND EXITS THRU BOTTOM	
680 EUTM=EUTM-PLANCK(1)*TOT+1	TRN1353
GO TO 710	TRN1354
690 ESIDE=ESIDE-PLANCK(1)*TOT+1	TRN1355
GO TO 710	TRN1356
700 EUTM=EUTM-PLANCK(1)*TOT+1	TRN1357
710 IPASS=2	TRN1358
C IF THETA=0 AND REFLECTIVE BOUNDARY, ONLY MAKE ONE PASS ALONG RAYS	
IF (NEXIT.LE.0) GO TO 800	TRN1359
GO TO (370,400,420), NEXIT	TRN1360
C BOTH PASSES COMPLETE -- RETURN FOR NEXT RAY	TRN1361
720 GO TO (730,740,750), NTER	
730 EUTM = EUTM - PLANCK(1)*TOT+1	TRN1363
GO TO 760	
740 ESIDE = ESIDE - PLANCK(1)*TOT+1	TRN1365
GO TO 760	
750 ETOP = ETOP - PLANCK(1)*TOT+1	TRN1367
760 GO TO (770,780,790), NEXIT	
770 EUTM = EUTM + PLANCK(1)	TRN1369
GO TO 800	
780 ESIDE = ESIDE + PLANCK(1)	TRN1371
GO TO 800	
790 ETOP = ETOP + PLANCK(1)	TRN1373
800 CALL DVCHK (KDMT)	TRN1374
GO TO (810,820), KDMY	TRN1375
810 S1=7.0810	
GO TO 965	TRN1377
820 CONTINUE	
C END RAY LOOP	TRN1378
GO TO 310	
C ALL RAYS THACED -- TRANSPORT COMPLETE	
C START OF THICK-THICK EXPLICIT DIFFUSION CALCULATION	
C	


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C 8/3PI=8.37758 16/3PI*2=52.63789 THIS EXPLICIT DIFFUSION
C CALCULATION IS ONLY VALID FOR MONOFREQUENCY CASE
030 DO 825 I=1,IMAX
    KPP = I+1
    J = 0
    AV = TAU(I) * 8.37758
    AM = X(I) * 52.63789
021 K = KPP
    KPP = K + IMAX
    J = J + 1
    IF (J.GT.JMAX) GO TO 825
    IF (ALAMV(K) .LT. 0.5) GO TO 821
C CELL K IS THICK AND ACTIVE
    IF ( (1.6E.IMAX) .OH. (ALAMV(K+1).LT.0.5) ) GO TO 822
C CALCULATE NET DIFFUSION IN HORIZONTAL DIRECTION BETWEEN CELLS K,K+1
    DERM = DY(J)*(B(K)-B(K+1))/(ROSS(K)*DX(I) + ROSS(K+1)*DX(I+1))*AM
    ER(K) = ER(K) - DERM
    ER(K+1) = ER(K+1) + DERM
022 IF ( (J.6E.JMAX) .OH. (ALAMV(KPP).LT.0.5) ) GO TO 821
C CALCULATE NET DIFFUSION IN VERTICAL DIRECTION BETWEEN CELLS K,KPP
    DERY = (B(K)-B(KPP)) / (ROSS(K)*DY(J) + ROSS(KPP)*DY(J+1)) * AV
    EK(K) = EK(K) - DERY
    EK(KPP) = EK(KPP) + DERY
    GO TO 821
025 CONTINUE
C FINISH OF THICK-THICK EXPLICIT DIFFUSION CALCULATION
C
C
    NEWIND 3
    MWUP=MNU
    MWUP4=MNU4
    IF (IMNU-MWNU) 200,850,840
040 S1=7.1010
    GO TO 965
050 CONTINUE
070 WRITE (6,970) NC
    IF (AUSIDBGPRT) .LT. 1.E-20) GO TO 880
    WRITE (6,980) (EK(K),K=2,KMAX)
    ESUM = 0.
    DO 331 K=2,KMAX
    ESUM = ESUM + EK(K)
    ALAMV(K) = B(K) * 4. * PI * ROSS(K)
331 ALAMH(K)=ER(K)*(TAU(I)*DY(J)*ROSS(K))+4.*PI*B(K)
    WRITE (6,9010) (ALAMH(K),K=2,KMAX),ESUM
    WRITE (6,9010) (ROSS(K),K=2,KMAX)
    WRITE (6,9010) (B(K),K=2,KMAX)
    WRITE (6,9010) (ALAMV(K),K=2,KMAX)
C
C ADVANCE FREQ, STORE EMERGENT FLUX, TEST FOR COMPLETION OF GROUPS
C
C *****
C END FREQUENCY LOOP
C *****

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TRN1381
 TRN1382
 TRN1383
 TRN1384

 TRN1390
 TRN1391

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 TRN1393
 TRN1394
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 TRN1397
 TRN1398
 TRN1399

C	ITERATE ON TEMPERATURE	TRN1400
880	GO TO (890,910), NVEZ	TRN1401
890	NVEZ=2	TRN1402
	VEZ=NVEZ	TRN1403
	NY=NVEZ	TRN1404
	NU=0	TRN1405
	DO 900 K=2,KMAX	TRN1406
C	WORK, SOURCE TERMS OMITTED	TRN1407
	IF (JMR(KFIT(K),2).NE.1) GO TO 900	TRN1408
	OLDTH(K)=THETA(K)	TRN1409
	E=ALX(K)+ER(K)*DT/AMX(K)	TRN1410
	SV=TAU(1)*DY(J)/AMX(K)	
	CALL ES (SV,E,TEMP(1),TEMP(2),66)	TRN1412
	THETA(K)=0.5*(THETA(K)+TEMP(1))	TRN1413
900	CONTINUE	TRN1414
	IF (MTAG.EQ.0) CALL KAPPA	TRN1415
	GO TO 40	TRN1416
910	IF (ITAG.EQ.0) GO TO 930	TRN1417
	DO 920 K=2,KMAX	TRN1418
920	THETA(K)=OLDTH(K)	TRN1419
C	H A N G E I N T E R N A L E N E R G I E S	TRN1420
930	TEMP(1)=1.	TRN1421
	DO 940 K=2,KMAX	TRN1422
	IF (JMR(KFIT(K),2).NE.1) GO TO 940	TRN1423
	DE=ER(K)*DT/AMX(K)	TRN1424
	Q=SLUG*ALX(K)	TRN1425
	IF (ABS(DE).LT.ABS(Q)) GO TO 940	
	TEMP(1)=AMIN1(Q/ABS(DE),TEMP(1))	TRN1427
940	CONTINUE	TRN1428
	U(TEMP)=DT*TEMP(1)	
	IF (DTLEMP*.GT.FFB) GO TO 950	
	WRITE (6,970) NC	TRN1431
	WRITE (6,980) (ER(K),K=2,KMAX)	TRN1432
	WRITE (6,990) DE,K*ALX(K),AMX(K)	TRN1433
	S1=7.1068	TRN1434
	GO TO 965	
950	T=T-DT+UTEMP	
	DT=DTEMP	
	FOO=FOO*DT	
	ETH = ETH + FOO	
	BACC = BACC + EBTM*DT	
	SACC = SACC + ESIDE*DT	
	TACC = TACC + ETOP*DT	
C	FOO,ETH,BACC,SACC,TACC ARE IN ERG	TRN1447
	WRITE (6,1000) NC,DT	
	WRITE (6,1010) BACC,SACC,TACC,FOO	
	READ(4) FIOUT,P,U,V	
	RETURN	TRN1451
967	S1 = 7.1111	
965	READ(4) FIOUT,P,U,V	
	ISEND = 2	
	CALL EDIT	TRN1452
C	970 FORMAT (52H)INTERNAL ENERGY GAIN PER UNIT TIME DUE TO RADIATION,OH	TRN1454
	1 CYCLE =15)	

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980 FORMAT (1X,1P10E11.4) TRN1455
990 FORMAT (5H DE =1PE11.4,5H K =110,6H AIX =E11.4,6H AMX =E11.4) TRN1456
1000 FORMAT (19H0 FLUX FOR CYCLE I4.7H DT =1PE10.3) TRN1457
1010 FORMAT (9H BOTTOM =1PE11.4,6H SIDE =E11.4,7H TOP =E11.4,7H FOO TRN1458
1=E11.4) TRN1459
9010 FORMAT (1H0/(1X,1P10E11.4))
END TRN1460
```

SHORT2 ROUTINES

DRAW Subroutine Sequence

Subroutine DRAW

Subroutine INWARD

Subroutine OUTWRD

Subroutine DISTNC

Function JO(J, I, L, M)

Function JB(J, I, L, M)

Function SINT(A1, B1, EPS1, FUNC)

Function IN(Z)

Function BOT(Z)

Function OUT(Z)

Function GAMMO(IG, Z)

Transport Subroutine Sequence

Subroutine TRAN2

Function TRANS(SIGMA, D, XA, XB, YA, YB,
C, SUM, IACT, SP)

Function S(XX, YY)

Function SRC(KKK)

Subroutine BOTBND

Subroutine OUTBND

Subroutine TOPBND

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RAD4
RAD5
RAD6
RAD7
RAD8
RAD9
RAD10
RAD11
RAD12
RAD13
RAD14
RAD15
RAD16
RAD17

PP,
XIN,

LMAD,
TOLER,

HLP2,
TANW,

LMDA,
MUMAX,
TOLR,
CHKSUM

WRITE(6) FIOUT,CAP,P,U,V
REWINJ 4
REWIND 4
REWIND 3
KMAX = KMAX+1
STEP = 1.0283E+12
PI2 = PI/2.
FURPI = 4.*PI
PI4 = PI/4.
READ (5,9500)
WRITE (6,9510)
LMAX,MUMAX,TOLER,CHKSUM
LMAX,MUMAX,TOLER,CHKSUM
(LMATA(I),I=1,IMAX)
(RTHETA(J),J=1,JMAX)
(RTHETA(I),I=1,IMAX)
(RTHETA(J),J=1,JMAX)
(RTHETA(I),I=1,IMAX)
(RTHETA(J),J=1,JMAX)
PP = PI/FLOAT(LMAX)
L2 = LMAX/2
L2PI = L2 + 1
UO 270 L=LMAX
PHI(L) = PP * (FLOAT(L) - .5)
COSPHI(L) = COS(PHI(L))
DO 275 L=0,LMAX
275 SINPL(L) = SIN(PP*FLOAT(L))

```



```

SIMP2 = SIN(PP/2.)
M2 = MUMA/2
C MU=1 AND MU=0 ARE PROHIBITED CASES
C REQUIREMENT --
C CHOICE OF MVS MUST BE SYMMETRIC ABOUT 0.
DO 200 L=1,LMAX
  MU(L) = COSPHI(L)
200 MU(L) = SIN(PI(L))
DO 600 N=1,LMAX
  W(N) = ABS( COS(FLOAT(N)*PP) - COS(FLOAT(N-1)*PP) ) * .5
C
C
KK2 = 1
N00 = 0
DO 3000 M=1,M2
  DO 2000 J=1,JMAX
    NOBUFF = 1
  C ZERO OUT SUMCHECKS
  DO 650 K=1,LMAX
    DO 650 N=1,LMAX
      SUMB(N,K) = 0.
650 SUBJ(N,K) = 0.
  C ICHECK = 0
  C GOING INWARD
  I = IMAX
  I = I-1
  IF (I .GT. 1) GO TO 700
  C GOING OUTWARD
  DO 1000 I=1,IMAX
    CALL OUTWU
  1000 CONTINUE
  C SUMCHECK
  IF (ICHECK .EQ. 0) WRITE (6,9054) M,J
  I00 = 0
  I00 = 0
  DO 1500 N=1,IMAX
    DO 1500 K=1,LMAX
      IF (ABS(SUMB(N,K)-1.) .LT. CHKSUM) GO TO 1450
      WRITE (6,7050) M,J,N,K,SUMB(N,K)
      I00 = 1
1450 CONTINUE
      IF (ABS(SUMB(N,K)-1.) .LT. CHKSUM .OR. K.GT.L2) GO TO 1499
      WRITE (6,9051) M,J,N,K,SUMB(N,K)
      I00 = 1
1499 CONTINUE
1500 CONTINUE
      IF (I00 .EQ. 0) WRITE (6,9052) M,J
      IF (I00 .EQ. 0) WRITE (6,9053) M,J
      4000 WRITE (3) (BUFF2(I,M),I,M=1,1023)
      KK2 = 1
      N00 = 0
2000 CONTINUE
3000 CONTINUE

```

```

C
WRITE (3) (BUFF2(ILM),ILM=1,1023)
HEAD (4) FIOU,CAP,P,U,V
REWIND 4
WRITE (6,9876) NOBUFF
RETURN

C
9050 FORMAT (7H0FOR ME,12,3H JE,12,3H I=,12,4H LP=,12,12H SUMB(I,LP)=,1
      .PE19,6)
9051 FORMAT (7H0FOR ME,12,3H JE,12,3H I=,12,3H I=,12,4H LP=,12,12H SUMO(I,LP)=,1
      .PE19,6)
9052 FORMAT (7H0FOR ME,12,3H JE,12,29H CHECKSUMS FOR BOTTOM ETAS=1.)
9053 FORMAT (7H0FOR ME,12,3H JE,12,30H CHECKSUMS FOR OUTSIDE ETAS=1.)
9054 FORMAT (7H0FOR ME,12,3H JE,12,29H CHECKSUMS FOR INSIDE ETAS=1.)
9500 FORMAT (216,2E12.5)
9510 FORMAT (12H1 DRAW INPUT/5X,5HLMAX=,16/5X,6HNUMAX=,16/5X,6HTOLER=,1
      .PE19,6/5X,7HCHKSUM=,1PE19,6///)
9600 FORMAT (6E12.5)
9700 FORMAT (1H0,10X,9HATHETA(I)/(1PBE19,6))
9701 FORMAT (1H0,10X,9HATHETA(J)/(1PBE19,6))
9702 FORMAT (1H0,10X,9HATHETA(I)/(1PBE19,6))
9876 FORMAT (31H1 NUMBER OF BUFFERS PER LAYER =,15///)
      END

```

243

RAD1
RAD2
RAD
RAD3
RAD4
RAD5
RAD6
RAD7
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RAD9
RAD10
RAD11
RAD12
RAD13
RAD14
RAD15
RAD16
RAD17

COMMON
1L2,
2ISNEEP,
COMMON
1UELTY,
2P12,
3JMAX,
COMMON
1
2
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4
COMMON
COMMON
COMMON
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COMMON
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COMMON
REAL
MU , LNOA

/RAD1/
L2P1,
LLO4,
COMMON
/RAD2/
FORP1,
P14,
STEF
/RAD3/
COSPHI(6),
CH(6),
DUM2(1),
W(6),
/INEX/
/LNUH/
/DJ5/
X(6),
/SUNCH/
/BNOT/
MU , LNOA

IAID,
LMAX,
LHI
ANG1,
HELP1,
SIMP2,
DUM1(1),
CVIN(6),
MU(6),
SINLP(6),
CV(6)
M,L,LP
LCHECK,
UV(6),
YH(6)
KK2,
XV(6),
SUMB(50,6),
RTMETHA(50),
ATPETA(50)

ISEND,
ISENDO,
CKSUM,
PP,
XMIN,
AA(16),
CVOUT(6),
PHI(6),
SMU(6),
DH(6),
SUMO(50,6),
ATPETA(50)

ICHECK,
RMAX,
ASH,
LNOA,
TOLER,
RTMETHA(50),
ATPETA(50)

KMAX,
M2,
ANG2,
HELP2,
TANN,
DUM1(1),
CVIN(6),
MU(6),
SINLP(6),
CV(6)
KK2,
XV(6),
SUMB(50,6),
RTMETHA(50),
ATPETA(50)

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XMAX = X(I)
XMIN = X(I-1)
ASN = ASIN(X(I-1)/X(I))
LIM2 = DY(I) - (X(I)-X(I-1))/TAN(ASN)
IF (LIM2 .LT. 0.) GO TO 1080
L = 1
LP = 1
Y = 0.
1055 EPS = TOLER * CV(LP) / (X(I)-X(I-1))*SMU(M)
1060 GP = Y
IF (X(I)-SINLP(L)) .GT. X(I-1)+SINLP(L) GO TO 1070
C LP CANNOT FILL L
  ANGL1 = PP * FLOAT(LP)
  Q = SIN(0.0+LIM2*EPS*OUT)
  ETA01(L,LP) = (Q-GP) / CV(LP)
  * FORPI * X(I) * W(M) * SMU(M)
  LP = LP + 1
  GO TO 1055
C LP FILLS L
1070 CONTINUE
  ANGL1 = ASIN( X(I-1)+SINLP(L)/X(I) )
  Q = SIN(0.0+LIM2*EPS*OUT)
  ETA01(L,LP) = (Q-GP) / CV(LP)
  * FORPI * X(I) * W(M) * SMU(M)
  L = L + 1
  IF (L .LE. L2) GO TO 1060
  GO TO 1100
1080 DO 1090 L=1,L2
  DO 1090 LP=1,L2
1090 ETA01(L,LP) = 0.
1100 CONTINUE
C
IF (UGPRT .LT. 1.E-7) GO TO 1120
WRITE (6,9003) M,J,I
DO 1110 L=1,L2
  WRITE (6,9001) L,(ETA01(L,LP),LP=1,L2)
1110 CONTINUE
1120 CONTINUE
C COMPUTE U/LI,XV(LI),YV(LI)
  LLOW = L2P1
  LHI = LMAX
  ISENU = 1
  CALL DISTNC
1130 KK1 = 0
  DO 2000 L=1,L2
  DO 1995 LP=1,L2
    RUFF1(KK1+LP) = ETA01(L,LP)
    KK1 = KK1 + L2
  DO 1996 LP=1,L2
    RUFF1(KK1+LP) = ETA01(L,LP)
    KK1 = KK1 + L2
  RUFF1(KK1+1) = DV(LI)
  RUFF1(KK1+2) = XV(LI)
  RUFF1(KK1+3) = YV(LI)
  RUFF1(KK1+4) = CVIN(LI)

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      KK1 = KK1 + 4
      2000 CONTINUE
C
C   COMPUTE AND WRITE E T A B T
C   E T A B T (LP) = FRACTION OF ENERGY IN FAN LP PASSING INTO CELL
C   (I,J) THRU THE BOTTOM OF THE CELL THAT PASSES OUT CELL (I,J)
C   THRU THE TOP IN FAN L (LP MEASURED AT X(I))
      731 ISEMU = 2
      LMDA = DY(J) * (SMU(M)/MU(M))
      XMIN = X(I-1)
      XMAX = X(I)
      AID = DY(J) * DY(J) * (SMU(M)/MU(M)) * (SMU(M)/MU(M))
      HELP1 = AID - X(I-1) * X(I-1)
      HELP2 = AID - X(I) * X(I)
      DO 800 LP=1,L2
      EPS = TOLER * CH(LP) / (FORPL * W(M) * ABS(MU(M)))
      ANGL1 = X(I) * SIN(LP)
      ANGL2 = X(I) * SIN(LP-1)
C   ENTER INWARD, EXIT INWARD
      LIM1 = AMAX1( ANGL1, AMIN1( X(I), SORT(LMDA * LMDA * X(I-1) * X(I-1))) )
      IF 1 LIM1 .LT. X(I) , GO TO 733
      L = LP
      E T A B T (LP) = 0.
      GO TO 800
      733 I AID = 1
      L = LP
      E T A B T (LP) = SINT(LIM1, X(I), EPS * DOT) / CH(LP)
      * FOW1 * W(M) * ABS(MU(M))
      800 CONTINUE
C
      IF (LUGPHT .LT. 1.E-20) GO TO 900
      WRITE (6,9004) M,J,1
      DO 900 L=1,L2
      WRITE (6,9001) L, (E T A B T (L,LP), LP=1,L2)
      900 CONTINUE
      900 CONTINUE
C
C   COMPUTE AND WRITE E T A B T
C   E T A B T (LP) = FRACTION OF ENERGY IN FAN LP PASSING INTO CELL
C   (I,J) THRU THE OUTSIDE OF THE CELL THAT PASSES OUT CELL (I,J) THRU
C   THE TOP IN FAN L
      ISEMU = 2
      UELTY = DY(J)
      TANN = SMU(M)/MU(M)
      XMAX = X(I)
      XMIN = X(I-1)
      ASH = ASIN(X(I-1)/X(I))
      LIM1 = AMAX1( 0. , DY(J) * SORT(X(I) * X(I) - X(I-1) * X(I-1)) / TANN )
      DO 1200 LP=1,L2
      EPS = TOLER * CV(LP) / (FORPL * X(I) * W(M) * SMU(M))
      ANGL1 = PP * FLOAT(LP-1)
      ANGL2 = PP * FLOAT(LP)
      LIM2 = DY(J) - X(I) * COS(FLOAT(LP) * PP) / TANN

```



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C CASE1 -- EXIT INWARD
  IAD = 1
  L = LP
  ETAOT(L,LP) = SINT(LINI*DY(J),EPS,OUT) / CV(L,LP)
  * FUMPI * X(1) * BIN * SUMIN
1200 CONTINUE
C
  IF (UMGHT .LT. 1.E-20) GO TO 1300
  WRITE (6,9005) M,J,I
  DO 1290 L=1,L2
  WRITE (6,9001) L,(ETAOT(L,LP),LP=1,L2)
1290 CONTINUE
1300 CONTINUE
C COMPUTE UH(L),XH(L),YH(L)
  LLOW = L2P1
  LHI = LMAX
  ISENGU = 2
  CALL DISTAC
  DO 1970 LP=1,L2
  DO 2500 L=1,L2
  X1 = KK1 + L2
  1997 BUFF1(KK1+LP) = ETABT(L,LP)
  1998 BUFF1(KK1+LP) = ETAOT(L,LP)
  X1 = KK1 + L2
  BUFF1(KK1+1) = UH(L)
  BUFF1(KK1+2) = XH(L)
  BUFF1(KK1+3) = YH(L)
  BUFF1(KK1+4) = CH(L)
  X1 = KK1 + 4
2500 CONTINUE
  IF (KK1 .GT. 1023) STOP
  IF (KK2+KK1 .LT. 1023) GO TO 2505
  WRITE (5) (BUFF2(ILM),ILM=1,1023)
  HOSUFF = HOSUFF + 1
  KK2 = 1
  N00 = 0
2505 DO 2510 KK=1,KK1
2510 BUFF2(KK2+KK) = BUFF1(KK)
  KK2 = KK2 + KK1
  N00 = N00 + 1
2600 CONTINUE
  DO 2002 LP=1,L2
  DO 2002 L=1,L2
  SUM8(I,LP) = SUM8(I,LP) + ETABT(L,LP) + ETABT(L,LP)
  2002 SUM8(I,LP) = SUM8(I,LP) + ETAOT(L,LP) + ETAOT(L,LP)
  NCTUNH
C
9000 FORMAT (1H0/20H ETABT(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
9001 FORMAT (3H L=12/16X,1P=1H,6)
9003 FORMAT (1H0/20H ETAOT(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
9004 FORMAT (1H0/20H ETABT(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
9005 FORMAT (1H0/20H ETAOT(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
END

```

Subroutine OUTWRD

[illegible]

```

C
C
C
      REAL MERGE
      COMMON /RAD1/ L2P1, LMAX, IAIU, KMAX, ICHECK, ISEND,
      215HELP, LLOW, LHI, ADEL2, ASN, CVOUT(6),
      COMMON /RAD2/ ANG1, HELP1, LMDA, PP,
      1UELTV, FUPP1, SINP2, TANN, XMIN,
      2P12, SINAI, STEF, DUM1(1),
      COMMON /RAD3/ COSPHI(6),
      1 CH(6), MU(6),
      2 DUM2(1), SINLP(6),
      3 W(6), CV(6)
      COMMON /TICKX/ M,L,LP LCHECK, KK2, NOBUFF
      COMMON /DKUMH/ DV(6), TV(6), DH(6),
      COMMON /DIS/ YH(6)
      1 XH(6), SUMB(50.6), RTHETA(50), ATHETA(50)
      COMMON /SUNGCK/ BTHETA(50)
      COMMON /UNDT/
      REAL MU, LMDA
C
C
      EQUIVALENCE (FIOUT,OLDTH), (U,ROSS),
      (CAP1(1),BUFF1(1),BUFF1(1),NOI), (PI(1),PLANCK(1),BUFF2(1),NOO)
      . DIMENSION OLDTH(1), ROSS(1), BUFF1(1), BUFF(1), PLANCK(1),
      . BUFF2(1)
C
      EQUIVALENCE (V(1),XA(1),EB(1))
      DIMENSION XA(1)
      DIMENSION
      . ETAIO(6.6), ETAIV(6.6), EV(50.6), ET(50.6),
      . ETABO(6.6), ETAOI(6.6), ETAOT(6.6), ETABO(6.6),
      . ETABO(6.6), ETABO(6.6), ETABO(6.6), ETABO(6.6)
C
C
      REAL JO, JB
      REAL LIMI, LIM2
      EXTERNAL BOT
      EXTERNAL OUT
      EXTERNAL IN
      REAL IN
C
C
      C COMPUTE SPECIFIED CONSTANTS
      DO 90 L=0,LMAX
      90 AA(L) = SINLP(L) * A(1)
      DO 50 LP=2,1,LMAX
      50 CVIN(LP) = JO(JP1-1,LP,M)
      DO 60 LP=1,LMAX
      60 CVOUT(LP) = JO(J,1,LP,M)
      DO 70 LP=1,LMAX
      70 CH(LP) = JH(JP1-1,LP,M)

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C      ISHELP = 2
C      ZERO ETAS
DO 500 LP=1,LMAX
  ETAIL(LP) = 0.
  ETATIL(LP) = 0.
  ETABUL(LP) = 0.
  ETABOIL(LP) = 0.
  ETABOIL(LP) = 0.
  ETATIL(LP) = 0.
  ETATIL(LP) = 0.
500 CONTINUE
C      COMPUTE AND WRITE ETAIO, ETATIT
C      ETABOIL(LP) = FRACTION OF (ENERGY IN FAN LP PASSING INTO CELL
C      (I,J) THRU THE INSIDE OF THE CELL) THAT PASSES OUT CELL (I,J) THRU
C      THE OUTSIDE IN FAN L
C      ETATIL(LP) = FRACTION FROM INSIDE CELL (I,J) TO TOP
C      SUM OF ABOVE TWO FRACTIONS = 1.
  IF (I.EQ. 1) GO TO 2000
  DELTY = DY(J)
  TANN = SMU(M)/NU(H)
  XMAX = X(I)
  XMIN = X(I-1)
  HELP1 = X(I-1)*X(I-1) - X(I)*X(I)
  LP = LMAX
  EPS = TOLER + CVIN(LP) / (FORPI*(1-1)*W(M)*SMU(M))
  DO 600 N = 1, EPS
    L = LMAX
    H = 0.
    Q = 0.
    2020 QP = 0
    IF (X(I)*SIMP(L-1) .LT. X(I-1)*SIMP(LP-1)) GO TO 2030
    C      LP DOES NOT FILL L
    XI = SIMP(LP-1)
    ANML1 = PP * FLOAT(LP-1)
    LU = 0
    60 TO 2040
    C      LP AT LEAST FILLS L
    2030 XI = X(I) * SIMP(L-1)/X(I-1)
    ANML1 = P1 - ASIN(XI)
    LD = 1
    2040 R = DY(J) * XI
    IF (Q-LT.(2.*EPS)) .OW. DX(I)*(2.*EPS).GT.DY(J)*TANN )
      GO TO 2050
    C      OUTSIDE WALL RECEIVES SOME RADIATION FROM INSIDE
    Q = SIMT(0.,DY(J),EPS,IN)
    Q0 = Q - QP
    ETABOIL(LP) = Q0 / CVIN(LP)
    * FORPI * X(I-1) * W(M) * SMU(M)
    60 TO 2060
    2050 Q0 = 0.
    2060 ETATIL(LP) = (R - RP - Q0) / CVIN(LP)
    * FORPI * X(I-1) * W(M) * SMU(M)

```

```

LP = LP - 1 + LU
IF ( LP .EQ. L2 ) GO TO 2100
EPS = TOLER * CVIN(LP) / (FONP1 * X(I-1) * W(N) * SMU(N))
L = L - LU
GO TO 2020
2000 DO 2090 L=L2P1, LMAX
      DO 2090 LP=L2P1, LMAX
      ETAIO(L, LP) = 0.
2090 ETAIT(L, LP) = 0.
2100 CONTINUE
C
      IF (UW * PNT .LT. 1.E-20) GO TO 2150
      WRITE (6, 80003) M, J, I
      DO 2110 L=L2P1, LMAX
      WRITE (6, 80001) L, (ETAIO(L, LP), LP=1, LMAX)
2110 CONTINUE
      WRITE (6, 80004) M, J, I
      DO 2120 L=L2P1, LMAX
      WRITE (6, 80001) L, (ETAIT(L, LP), LP=1, LMAX)
2120 CONTINUE
2150 CONTINUE
C
C C H E C K S U M   E T A I O ,   E T A I T
C
      IF (1 .EQ. 1) GO TO 2210
      DO 2200 LP=L2P1, LMAX
      SUM00 = 0.
      SUMT1 = 0.
      SUM = 0.
      DO 2190 L=L2P1, LMAX
      SUM00 = SUM00 + ETAIO(L, LP)
      SUMT1 = SUMT1 + ETAIT(L, LP)
2190 CONTINUE
      SUM = SUM00 + SUMT1
      IF (ABS(SUM-1.0) .LT. CHKSUM) GO TO 2200
      ICECK = 1
      WRITE (6, 90017) M, J, I
      WRITE (6, 90018) LP, SUMT1, SUM00 SUM
2200 CONTINUE
C
2210 ISEMU = 2
      LMDA = DY(J) * (SMU(N)/MU(N))
      XMIN = X(I-1)
      XMAX = X(I)
      AID = DY(J) * DY(J) * (SMU(N)/MU(N)) * (SMU(N)/MU(N))
      HELP1 = AID - X(I-1) * X(I-1)
      HELP2 = AID - X(I) * X(I)
      DO 950 LP=1, LMAX
      EPS = TOLER * CH(LP) / (FONP1 * W(N) * ABS(MU(N)))

```

```

      ANGL1 = X(I) * SIN(LP)
      ANGL2 = X(I) * SIN(LP-1)
      IF (LP .GT. L2) GO TO 925
      C ENTER INWARD, EXIT OUTWARD
      LIM1 = ANGL1( ANGL2 , X(I-1) )
      LAID = 4
      L = LMAX + 1 - LP
      ETABO(L,LP) = SINT(LIM1,X(I),EPS,BOT) / CH(LP)
      * * FORPI * N(M) * ABS(MU(M))
      GO TO 949
      C ENTER OUTWARD, EXIT OUTWARD
      925 LIM1 = ANGL1( ANGL1 , X(I-1) )
      LAID = 5
      L = LP
      ETABO(L,LP) = SINT(LIM1,X(I),EPS,BOT) / CH(LP)
      * * FORPI * N(M) * ABS(MU(M))
      949 CONTINUE
      950 CONTINUE
      C
      IF (DUMPRT .LT. 1.E-20) GO TO 970
      WRITE (6,8005) M,J,I
      DO 960 L=L2P1,LMAX
      WRITE (6,8001) L,ETABO(L,LP),LP=1,LMAX)
      960 CONTINUE
      970 CONTINUE
      C
      C COMPUTE AND WRITE ETABO
      C ETABO(L,LP) = FRACTION OF ENERGY IN FAN LP PASSING INTO CELL
      C (I,J) THRU THE OUTSIDE OF THE CELL THAT PASSES OUT CELL (I,J) THRU
      C THE OUTSIDE IN FAN L
      ISEND = 3
      DELTY = DY(J)
      TANN = SMU(M)/MU(M)
      XMAX = X(I)
      XMIN = X(I-1)
      ASN = ASIN(X(I-1)/X(I))
      DO 1300 LP=1,L2
      ANGL1 = PP * FLOAT(LP-1)
      ANGL2 = PP * FLOAT(LP)
      L = LMAX + 1 - LP
      LIM2 = DY(J) - 2.*X(I)*COS(PP*FLOAT(LP))/TANN
      IF (LIM2 .LT. 0.) GO TO 1230
      EPS = TOLER * CVOUT(LP) / (FORPI*X(I)*N(M)*SMU(M))
      ETABO(L,LP) = SINT(LIM2,EPS,OUT) / CVOUT(LP)
      * * FORPI * X(I) * N(M) * SMU(M)
      GO TO 1240
      1230 ETABO(L,LP) = 0.
      1240 CONTINUE
      1300 CONTINUE
      C
      IF (DUMPRT .LT. 1.E-20) GO TO 1320
      WRITE (6,8006) M,J,I
      DO 1310 L=L2P1,LMAX
      WRITE (6,8001) L,ETABO(L,LP),LP=1,L2)

```



```

1310 CONTINUE
1320 CONTINUE
C COMPUTE DVIL, XVIL, YVIL
LLOW = 1
LHI = L2
LXNUM = 1
CALL DISTHC
KK1 = 0
DO 3000 L=L2PI, LMAX
DO 2990 LP=L2PI, LMAX
2990 BUFF1(KK1+LP-L2) = ETAIO(L, LP)
KK1 = KK1 + L2
DO 2991 LP=L2PI, LMAX
2991 BUFF1(KK1+LP) = ETABO(L, LP)
KK1 = KK1 + LMAX
DO 2992 LP=L2, L2
2992 BUFF1(KK1+LP) = ETAOO(L, LP)
KK1 = KK1 + L2
BUFF1(KK1+1) = DVIL
BUFF1(KK1+2) = XVIL
BUFF1(KK1+3) = YVIL
BUFF1(KK1+4) = CVOUT(L)
KK1 = KK1 + 4
3000 CONTINUE
C
C COMPUTE AND WRITE ETABT (EXITING OUTWARD)
C ETABT(L, LP) = FRACTION OF (ENERGY IN FAN LP PASSING INTO CELL
C (I,J) THRU THE BOTTOM OF THE CELL) THAT PASSES OUT CELL (I,J)
C THRU THE TOP IN FAN L (LP MEASURED AT X(I))
LMDA = 2
LMDA = DY(J) * (SMU(M)/MU(M))
XMIN = X(I-1)
XMAX = X(I)
AID = DY(J)*DY(J)*(SMU(M)/MU(M))*(SMU(M)/MU(M))
HELP1 = AID - X(I-1)*X(I-1)
HELP2 = AID - X(I)*X(I)
DO 800 LP=L2, LMAX
EPS = TOLER * CH(LP) / (FORPI * W(M) * ABS(MU(M)))
ANGL1 = X(I) * SIN(LP/LP)
ANGL2 = X(I) * SIN(LP/LP-1)
IF (LP * 61. L2) GO TO 740
C ENTER INWARD, EXIT OUTWARD
735 LIM1 = AMAX1( X(I-1) , ANGL2 )
LAI2 = 2
L = LMAX + 1 - LP
ETABT(L, LP) = SINT(LIM1, X(I), EPS, BOT) / CH(LP)
* FOMPI * W(M) * ABS(MU(M))
GO TO 800
740 CONTINUE
C ENTER OUTWARD, EXIT OUTWARD
LIM1 = AMAX1( X(I-1) , ANGL1 )
LAI2 = 3
L = LP
ETABT(L, LP) = SINT(LIM1, X(I), EPS, BOT) / CH(LP)

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      * FORPI = W(N) * ABS(MU(N))
800 CONTINUE
C
      IF (DUPRT .LT. 1.E-20) GO TO 900
      WRITE (6,8000) M,J,I
      DO 830 L=L2P1,LMAX
      WRITE (6,8001) L,ETABT(L,LP),LP=L1,LMAX)
830 CONTINUE
900 CONTINUE
C
C COMPUTE AND WHITE FJ AOT (EXITING OUTWARD)
C ETAOT(L,LP) = FRACTION OF ENERGY IN FAN LP PASSING INTO CELL
C (I,J) THRU THE OUTSIDE OF THE CELL THAT PASSES OUT CELL (I,J) THRU
C THE TOP IN FAN L
C CASE2 -- EXIT OUTWARD
      ISEMU = 2
      IAD = 2
      UELTY = DY(J)
      TANN = SMU(N)/MU(N)
      XMAX = X(I)
      XMIN = X(I-1)
      ASN = ASIN(X(I-1)/X(I))
      DO 1200 LP=L1,L2
      L = LMAX + 1 - LP
      LIM2 = DY(J) - X(I)*COS(FLOAT(LP)*PP)/TANN
      IF (LIM2 .LT. 0.) GO TO 1130
      ANGL1 = PP * FLOAT(LP-1)
      ANGL2 = PP * FLOAT(LP)
      EPS = TOLER * CVOUT(LP) / (FORPI=X(I)*W(N)*SMU(N))
      LIM1 = ANGL1 + EPS
      LIM2 = ANGL2 - EPS
      ETAOT(L,LP) = SINT(LIM1-LIM2-EPS-OUT) / CVOUT(LP)
      * FORPI = X(I) * W(N) * SMU(N)
      GO TO 1200
1130 ETAOT(L,LP) = 0.
1200 CONTINUE
C
      IF (DUPRT .LT. 1.E-20) GO TO 1220
      WRITE (6,8002) M,J,I
      DO 1210 L=L2P1,LMAX
      WRITE (6,8001) L,ETAOT(L,LP),LP=L1,L2)
1210 CONTINUE
1220 CONTINUE
C COMPUTE DH(L),XH(L),YH(L)
      LLOW = 1
      LHI = L2
      ISEMU = 2
      CALL DISTNC
      DO 3000 L=L2P1,LMAX
      DO 3001 LP=L2P1,LMAX
      XH1 = XH1 + L2
      DO 3005 LP=L1,LMAX
      YH1(YH1+LP) = ETABT(L,LP)
      XH1 = XH1 + LMAX
3000
3005

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```

DO 3003 LP=1,L2
  WUFF1(KK1+LP) = ETAT(L,LP)
  KK1 = KK1 + L2
  WUFF1(KK1+1) = DM(L)
  WUFF1(KK1+2) = VM(L)
  WUFF1(KK1+3) = TH(L)
  WUFF1(KK1+4) = CH(L)
  KK1 = KK1 + 4
3000 CONTINUE
  IF (KK1 .GT. 1023) STOP
  IF (KK2+KK1 .LT. 1023) GO TO 3505
  WRITE (3) (WUFF2(ILM),ILM=1,1023)
  NBUUFF = NBUUFF + 1
  KK2 = 1
  N00 = 0
  DO 3510 KKK=1,KK1
  3505 DO 3510 KKK=1,KK1
  3510 WUFF2(KK2+KKK) = WUFF1(KKK)
  KK2 = KK2 + KK1
  N00 = N00 + 1
3600 CONTINUE
  DO 3100 LP=1,LMAX
  3100 L=L2P1,LMAX
  SUM0(1,LP) = SUM0(1,LP) + ETAB0(L,LP) + ETAB1(L,LP)
  SUM0(1,LP) = SUM0(1,LP) + ETAB0(L,LP) + ETAB1(L,LP)
  RETURN
C
8000 FORMAT (1H0/40H ETAT(L,LP) (SWEEPING OUTWARD) FOR M=12.3H J=1
      .2.3H I=12/60X,2HLP)
8001 FORMAT (3H I=12/6X,1P8E14.6)
8002 FORMAT (1H0/40H ETAT(L,LP) (SWEEPING OUTWARD) FOR M=12.3H J=1
      .2.3H I=12/60X,2HLP)
8003 FORMAT (1H0/20H ETAB0(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
8004 FORMAT (1H0/20H ETAB1(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
8005 FORMAT (1H0/20H ETAB0(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
8006 FORMAT (1H0/20H ETAB0(L,LP) FOR M=12.3H J=12.3H I=12/60X,2HLP)
9017 FORMAT (2H0SUM OF INSIDE E:AS FOR M=12.4H, J=12.4H, I=12.12H
      .OVER L FOR 15X,4MSUM,11X,4MSUM,11X,3MSUM/)
9018 FORMAT (9X,3HLP=12,1P3E15.6)
9019 FORMAT (29H0CHECKSUMS FOR INSIDE ETAS=1.)
      END

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2PI2,  PI4,  SINP2,  TANN,  TOLER,  XMIN,  RAD5
JMAX,  STEF,  /RAD3,  DUM1(1),  AA(16),  RAD6
COMMON  COSPHI(6),  CVIN(6),  CVOUT(6),  RAD7
1  CH(6),  MU(6),  PHI(6),  RAD8
2  DUM2(1),  SIMLP(6),  SMU(6),  RAD9
3  M(6),  CV(6),  RAD10
COMMON  /INDEX,  M,L,LP,  NOBLE'S,  RAD11
COMMON  /URUMM,  LCHECK,  KK2,  YV(6),  RAD12
COMMON  /UIS,  DV(6),  XV(6),  DH(6),  RAD13
COMMON  XM(6),  YH(6),  RAD14
1  /SURCHK,  SUMB(50,6),  SUMO(50,6),  RAD15
COMMON  /BNDT,  BTHETA(50),  RTHETA(50),  ATHTA(50),  RAD16
REAL  MU,  LMUA,  RAD17

EQUIVALENCE (FLOUT,OLDTH), (U,ROSS),
. (CAP1),BUFF1(1),BUFF(1),NO1), (P1),PLANCK(1),BUFF2(1),NOO)
. DIMENSION OLDTH(1), ROSS(1), BUFF1(1), BUFF(1), PLANCK(1),
. BUFF2(1)

EQUIVALENCE (V(1),XA(1),EB(1))
DIMENSION XA(1)
DIMENSION EB(50,6), EV(50,6), ET(50,6),
. ETAIO(6,6), ETAII(6,6), ETABI(6,6), ETABT(6,6),
. ETABO(6,6), ETAOI(6,6), ETAOT(6,6), ETAOO(6,6)

NOTE -- IN COMMON, THE FOLLOWING VARIABLES HAVE CHANGED NAMES IN
THIS ROUTINE
LLOW GOES TO LIM1
LHI GOES TO LIM2
I = 11
GO TO (5,55), ISENUU

DV(L) = L,LE,L2(SWEEPING INWARD)
CHARACTERISTIC LENGTH OF FAN (L,M) IN CELL (I,J) MEASURED FROM
X(I-1),Y(J-5) IN DIRECTION PHI(LMAX+1-L),(PI-ACOS(MU(M)))
TO THE NEAREST CELL FACE
L,GT,L2(SWEEPING OUTWARD)
DV(L) = L,GT,L2(SWEEPING OUTWARD)
CHARACTERISTIC LENGTH OF FAN (L,M) IN CELL (I,J) MEASURED FROM
X(I),Y(J-5) IN DIRECTION PHI(LMAX+1-L),(PI-ACOS(MU(M))) TO
THE NEAREST CELL FACE
ALL L
XV(L) = THE X-COORDINATE OF THE INTERSECTION OF DV(L) AND THE CELL
FACE
FACE
YV(L) = ALL L
Y(J-5) - (THE Y-COORDINATE OF THE INTERSECTION OF DV(L) AND
THE CELL FACE)

IF (ISWEEP.EQ. 1) I=I-1
IF (DUGPKT.LT. 1.E-20) GO TO 10
WRITE (6,8000) M,II,J

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```

10 DO 500 L=LIM1,LIM2
C CONSIDER LENGTHS IN Y-PLANE
  IBACK = 0
  IF (L .GT. L2) GO TO 250
  IF (I .LE. 1) GO TO 50
  IF (PHI(L) .GT. ASIN( X(I-1)/X(I) ) ) GO TO 50
C HITS X(I-1) CIRCLE
  HIT = X(I-1)
  BBB = -1.
  GO TO 340
C HITS X(I) CIRCLE
  HIT = X(I)
  IBACK = 1
  JVV = X(I) * COSPHI(L) * 2.
  GO TO 355
250 IF (I .EQ. IMAX) GO TO 500
C HITS X(I+1) CIRCLE
  HIT = X(I+1)
  BBB = +1.
  340 A = X(I) * COSPHI(L)
  B = SQR( X(I)*X(I) ) * (COSPHI(L)*COSPHI(L)-1.) + HIT*HIT
  DVV = A + B * BBB
  355 IF (UGPNT .LT. 1.E-20) GO TO 360
  K = LMAX + 1 - L
  WRITE (6,8001) K,DVV
C CONSIDER 3-D LENGTHS
  360 IF ( UVV/(SMU(M)/MU(M)) .GT. DT(J)/2. ) GO TO 370
C HITS SIDE FIRST
  UV(LMAX+1-L) = UVV/SMU(M)
  TV(LMAX+1-L) = HIT
  TV(LMAX+1-L) = DVV/(SMU(M)/MU(M))
  GO TO 500
C HITS BOTTOM FIRST
  370 UV(LMAX+1-L) = DT(J)/(2.*MU(M))
  HELP = DT(J)/2. * (SMU(M)/MU(M))
  IF ( IBACK.EQ.1 .AND. HELP.GT.DVV/2. ) HELP=DVV*HELP
  XV(LMAX+1-L) = SQR( X(I)*X(I) + HELP*HELP - 2.*X(I)*HELP*
    . COSPHI(L) )
  TV(LMAX+1-L) = DT(J)/2.
  500 CONTINUE
C
  IF (UGPNT .LT. 1.E-20) GO TO 1000
  LIM1P = LMAX + 1 - LIM2
  LIM2P = LMAX + 1 - LIM1
  WRITE (6,8002) (L,DV(L),L=LIM1P,LIM2P)
  WRITE (6,8003) (L,HIT(L),L=LIM1P,LIM2P)
  WRITE (6,8004) (L,XV(L),L=LIM1P,LIM2P)
  WRITE (6,8005) (L,TV(L),L=LIM1P,LIM2P)
  940 GO TO 1000
C
C UMILJE
C CHARACTERISTIC LENGTH IN FAN (L-M) IN CELL (I,J) MEASURED FROM
C X(1-5),Y(1-J) IN DIRECTION PHI(LMAX+1-L),(PI-ACOS(MU(M))) TO

```


.LENGTHS FOR VERTICAL SURFACE/)
0001 FORMAT (3H L=12.0H LENGTH=.1PE15.6)
0002 FORMAT (1H0/40H CHARACTERISTIC LENGTHS FOR VERTICAL SURFACE/)
0003 FORMAT (1H0/30H X-COORDINATE OF INTERSECTIONS/)
0004 FORMAT (1H0/52H PLANE CHARACTERISTIC LENGTHS FOR HORIZONTAL SURFA
.CE/)
0005 FORMAT (1H0/40H CHARACTERISTIC LENGTHS FOR HORIZONTAL SURFACE/)
0006 FORMAT (1H0/30H T-COORDINATE OF INTERSECTIONS/)
END

Function JO(J, I, L, M)

[illegible]

[illegible]

REAL FUNCTION $\Phi(J,I,L,M)$
 $\Phi(J,I,L,M)$ = RATE AT WHICH ENERGY IS PASSING INTO CELL (I,J,I) THRU
 THE BOTTOM SURFACE ON INTO CELL (I,J) THRU THE TOP SURFACE
 IN $FAN L$ (MEASURED AT $X(I)$ IN DIRECTION MUM) IN THE CASE OF
 ISOTROPIC RADIATION, $MOSSES$ (DIVIDED BY 10)

REAL	LIM1 • LIM2
REAL	II
REAL	NU • LNUA
REAL	MENG

[illegible]

```

COMMON
1  /RAD3/      DUM1(1),      AAL16),
   COSPHI(6),  COUT(6),
   CH(6),      MU(6),        PHI(6),
2  DUM2(1),    SINLP(6),     SMU(6),
   W(6),       CV(6)
3
4  /INDEX/     M,L,LP
COMMON        /CHECK/      KK2,      NOBUFF
COMMON        /DIS/        XV(6),    DV(6),
COMMON        /TH(6),      TH(6),
1  /SUMCHK/    SUMB(50.6),  SUMO(50.6)
COMMON        /BNDT/      BTHETA(50),  ATHTA(50)
2  RAD7
   RAD8
   RAD9
   RAD10
   RAD11
   RAD12
   RAD13
   RAD14
   RAD15
   RAD16
   RAD17

C
C  DIMENSION  OLDTH(1), NOSS(1), BUFF(1), PLANK(1),
   . BUFF2(1)

C
C  DIMENSION  XA(1)
DIMENSION
   ETAIO(6.6), ETAB(6.6), ETABT(6.6),
   . ETAO(6.6), ETAOI(6.6), ETAOT(6.6), ETAOO(6.6)

C
C  EQUIVALENCE  (I,LEFT,UL),      (PL,PR)
EQUIVALENCE  (PL(101),SIGC)
EQUIVALENCE  (S2U,BLANK(12)), (ZP126,BLANK(13)), (ZP136,BLANK(14)),
EQUIVALENCE  (S2U,BLANK(12)), (ZP126,BLANK(13)), (ZP136,BLANK(14)),
1(KMX,BLANK(15)),(MT,BLANK(16)),(ANUMB,BLANK(19)),(LOC2,BLANK(22))
2,(H,BLANK(25)),(NA,BLANK(26))
EQUIVALENCE  (BACC,BLANK(27)),(SACC,BLANK(28)),(TACC,BLANK(29))

C
C  EQUIVALENCE  (FIOUT,OLDTH), (U,ROSS),
   . (CAP(1),BUFF(1),BUFF(1),NOI), (P(1),PLANK(1),BUFF2(1),NOO)
EQUIVALENCE  (V(1),XA(1),Ed(1))

C
C  I1(A,R) = .5*( NO*ASIN(A/R) + A*SORT(R*R-ASA) )

C
C  LL = L
IF (L .GT. L2) LL=MAX+1-L
LIM1 = AMAX1( X(I-1), AALL-1 )
LIM2 = AMAX1( X(I-1), AALL )
J0 = 11(AALL)*X(I) - 11(AALL)*LIM2 - 11(AALL-1)*X(I)
   + 11(AALL-1)*LIM1 + P10*(LIM2-LIM1)
J0 = J0 * FORP1 * W(M) * ABS(MU(M))
RETURN
END

```

Function SIN(A1, B1, EPS1, FUNC)

```

C      FUNCTION SIN(A1,B1,EPS1,FUNC)
C      -----
C      THIS ROUTINE DOES SIMPSON'S INTEGRATION
C      ALGORITHM 182 COMMUNICATIONS OF THE ACM
C      -----
C      A IS THE LOWER LIMIT OF INTEGRATION
C      B IS THE UPPER LIMIT OF INTEGRATION
C      EPS IS USED IN THE ACCURACY CRITERION
C      FUNC IS THE EVALUATING FUNCTION NAME
C      -----
C      DIMENSION UX(30),EPSP(30),X2(30),X3(30),F2(30),F3(30),F4(30),
C      1 FMP(30),FBP(30),EST2(30),EST3(30),LR(30),PVAL(30,3)
C      A = A1
C      B = B1
C      EPS = EPS1
C      L = 0
C      AGR = 1.0
C      L30 = 0
C      EST = 1.0
C      DA = B-A
C      FA = FUNC(A)
C      FM = FUNC((A + B)/2.0)
C      FB = FUNC(B)
C      110 L = L+1
C      UX(L) = DA/3.0
C      SX = UX(L)/6.0
C      F1 = 4.0*FUNC(A+UX(L)/2.0)
C      X2(L) = A + UX(L)
C      F2(L) = FUNC(X2(L))
C      X3(L) = X2(L) + UX(L)
C      F3(L) = FUNC(X3(L))
C      EPSP(L) = EPS
C      F4(L) = 4.0 * FUNC(X3(L) + UX(L)/2.0)
C      FMP(L) = FM
C      EST1 = (FA + F1 + F2(L)) * SX
C      FBP(L) = FB
C      EST2(L) = (F2(L) + F3(L) + FM) * SX
C      EST3(L) = (F3(L) + F4(L) + FB) * SX
C      SUM = EST1 + EST2(L) + EST3(L)
C      IF (ABS(L57-SUM) < LT, LPS(L) *AND, L.NE. 1) GO TO 10
C      IF (L > 64, 30) GO TO 13
C      1 LR(L) = 1
C      UA = UX(L)
C      FM = F1
C      FB = F2(L)
C      LPS = EPSP(L)/1.7
C      EST = EST1
C      20 TO 110
C      2 LR(L) = 2
C      UA = UX(L)
C      FB = F2(L)
C      FM = FMP(L)
C      FB = F3(L)
C      EPS = EPSP(L)/1.7
C      EST = EST2(L)
C      A = X2(L)

```

```

GO TO 110
3  LR(L) = J
   DA = UX(L)
   FA = FJ(L)
   FM = FL(L)
   FR = FR(L)
   EPS = EPS(L)/1.7
   EST = EST(L)
   A = AJ(L)
   GO TO 110
* SUM = PVAL(L,1) + PVAL(L,2) + PVAL(L,3)
  IF (L .GT. 1) GO TO 10
  SINT = SUM
  RETURN
C   DONE AT THIS LEVEL
10  L = L-1
    LT0 = LR(L)
    PVAL(L,LT0) = SUM
    GO TO (2,3,4), LT0
15  IF (100 .EQ. 1) GO TO 10
    I30 = 1
    WRITE (6,100)
    FORMAT (22H SINT REACHED LEVEL 30)
100 GO TO 10
    END

```

268


```

COMMON /RAD2/ ANGL1, ANGL2, ASN, CHKSUM,
1DELTY, FORPI, HELP1, LMDA, PP,
2PI2, PI4, SINF2, TOLER, XMIN,
3XMAX, STEF
COMMON /RAD3/ DUM1(1), AA(16),
1 COSPHI(6), CVIN(6), CVOUT(6),
2 CH(6), MU(6), PHI(6),
3 DUM2(1), SIMPL(6),
4 W(6), CV(6)
COMMON /INDEX/ M,L,LP
COMMON /DRUMH/ LCHECK, KK2, NOBUFF
COMMON /DIS/ DV(6), YV(6), DH(6),
1 XH(6), YH(6)
COMMON /SUMCHK/ SUMB(50,6), SUMO(50,6),
COMMON /BNUT/ BTHETA(50), KTHETA(50), ATHTA(50)
REAL NU, LMDA
EQUIVALENCE (V(1),XA(1),EB(1))
DIMENSION XA(1)
DIMENSION EB(50,6), EV(50,6), ET(50,6),
1 ETAIO(6,6), ETAIT(6,6), ETABI(6,6), ETABT(6,6),
2 ETABO(6,6), ETAOI(6,6), ETAOT(6,6), ETAOO(6,6)
CALCULATE GAMMA(1-1)
LMDA = TANN * (DELTY-2)
IF (XMAX .GT. XMIN+LMDA) GO TO 10
IF (LMDA .GT. XMAX-XMIN) GO TO 20
G2,4 = AMAX1( P12, ACOS( (LMDA+LMDA+HELP1)/(2.*XMIN+LMDA) ) )
GO TO 30
GAM = P1
GO TO 30
GAM = P12
GO TO 30
COMMON /INTEGRAND/ IN = SINI AMAX1(ANGL1, GAM )
30 RETURN
END

```

270

COMMON	/RADJ/	DUM(1),	AA(16),	RAD7
1	COSPHI(6),	CVIN(6),	CVOUT(6),	RAD8
2	CHI(6),	MU(6),	PHI(6),	RAD9
3	DUM2(1),	SINLP(6),	SMU(6),	RAD10
4	W(6),			RAD11
COMMON	/INLEX/	M,L,LP		RAD12
COMMON	/DRUMA/	LCHECK,	NOBUFF	RAD13
COMMON	/DIS/	UV(6),	YV(6),	RAD14
1	XH(6),	YH(6)	DH(6),	RAD15
COMMON	/SUMCHK/		SUM0(50,6),	RAD16
COMMON	/BNDT/	BTHETA(50),	ATHETA(50)	RAD17
REAL	MU , LMDA			
	EQUIVALENCE	(FIOUT,OLDTH) , (U,ROSS) ,		
	(CAP(1),BUFF(1),BUFF(1),NOI) , (P(1),PLANCK(1),BUFF2(1),NOO),			
	DIMENSION	OLDTH(1) , NOSS(1) , BUFF(1) , BUFF(1) , PLANCK(1) ,		
	BUFF2(1)			
	EQUIVALENCE	(V(1),XA(1),EW(1))		
	DIMENSION	XA(1)		
	DIMENSION	EB(50,6) , EV(50,6) , ET(50,6) ,		
	ETA0(6,6) , ETAT(6,6) , ETAB(6,6) , ETABT(6,6) ,			
	ETAB0(6,6) , ETA01(6,6) , ETA0T(6,6) , ETA00(6,6)			
	GO TO (10,20) , ISENO			
	FUNCTION FOR ETAB1			
10	IF (2,LT,XMIN+LMDA .AND. I.NE.1) GO TO 11			
	GAM = 0.			
	GO TO 17			
11	IF (LMDA .GT. SORT((2+2-XMIN+XMIN)) GO TO 15			
	GAM = ACOS((HELP1 + Z+Z) / (2. * Z * LMDA))			
	GO TO 17			
15	GAM = ASIN(XMIN/Z)			
17	NOT = Z+1 AMIN(ASIN(AMIN(1/Z),GAM))			
	GO TO 100			
	FUNCTIONS FOR ETABT AND ETAB0			
	CALCULATE GAMMAS			
	LOWER GAMMA			
20	IF (2,LT.(XMIN+LMDA) .AND. I.NE.1) GO TO 22			
	GAM1 = 0.			
	GO TO 26			
22	IF (LMDA .GT. SORT((2+2-XMIN+XMIN)) GO TO 25			
	GAM1 = ACOS((HELP1+Z+Z)/(2.*Z*LMDA))			
	GO TO 26			
25	GAM1 = ASIN(XMIN/Z)			
	UPPER GAMMA			
26	IF (2 .LE. XMAX+LMDA) GO TO 27			
	IF (2 .LE. LMDA+XMAX) GO TO 28			

```

HELP = (HELP2 + Z*Z)/(2.*Z*LMUA)
IF (HELP .LT. 0.) HELP=ANA.1(HELP,-1.)
IF (HELP .GT. 0.) HELP=AMINI(HELP,1.)
GAM = ACOS(HELP)
GAM2 = AMAX1(GAM,GAM)
GO TO 29
27 GAM2 = PI
GO TO 29
28 GAM2 = GAM1
C
29 GO TO (23,24,31,40,50) , IAI0
C ETABT
23 AAB = ACOS(LMUA/Z)
A = AMINI(AAB,GAM2,ASIN( AMINI(ANGL1/Z,1.) ) )
B = AMAX1( GAM1 , ASIN(AA/Z) )
GO TO 21
24 AAB = ACOS( AMINI(LMUA/Z,1.) )
A = AMINI ( GAM2 , ASIN( AMINI(1.,ANGL1/Z) ) )
B = AMAX1( ASIN(ANGL2/Z),GAM1,AAB )
GO TO 21
31 A = AMINI( (PI-ASIN(ANGL1/Z)), GAM2)
B = (PI - ASIN( AMINI(ANGL2/Z,1.) ) )
21 BOT = Z * AMAX1(A-B,0.)
GO TO 100
C ETAB0
40 A = ASIN( AMINI(1.,ANGL1/Z) )
B = AMAX1( GAM2 , ASIN(ANGL2/Z) )
GO TO 61
50 A = PI - ASIN( ANGL1/Z )
B = AMAX1( GAM2 , (PI - ASIN( AMINI(1.,ANGL2/Z) ) ) )
61 BOT = Z * AMAX1(A-B,0.)
100 RETURN
END

```

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RAD5
RAD6
RAD7
RAD8
RAD9
RAD10
RAD11
RAD12
RAD13
RAD14
RAD15
RAD16
RAD17

PP,
XMIN,

LMDA,
TOLER,

HELP2,
TANN,

HELP1,
SINP2,

FOMPI,
PI4,
STEP,
RAD3/
RAD3/
COSPI(6),
CH(6),
DUM2(1),
V(6),
/INDEX/
/DRUM/
/DIS/
XH(6),
/COMMON
/SUNCHK/
/END/

M.L.P
LCHK,
DV(6),
YH(6),

KK2,
XV(6),
SUMB(50,6),
RTHETA(50),
ATHETA(50)

SUMO(50,6)
DH(6),

C
REAL MU, LMDA

EQUIVALENCE (FIOUT,OLDTH), (U,ROSS),
(CAPI1, BUFF(1)), (NOI), (PI1), PLANCK(1), BUFF2(1), NOI)
.
DIMENSION OLDTH(1), ROSS(1), BUFF(1), BUFF1(1), BUFF(1), PLANCK(1),
. BUFF2(1)

EQUIVALENCE (V(1), XA(1), EB(1))
DIMENSION XA(1)
EB(50,6), EV(50,6), ET(50,6),
. ETAB(6,6), ETALT(6,6), ETAB1(6,6), ETABT(6,6),
. ETABO(6,6), ETAO(6,6), ETABT(6,6), ETABO(6,6)

GO TO (10,20,50), ISEND
C
INTEGRAND FOR ETAOI
C
10 OUT = SIN( AMINI(ANGL1,GAMMO(1,Z)) )
GO TO 70
C
INTEGRAND FOR ETABT
C
20 GO TO (30,40), IABD
C
CASE1 -- EXIT INWARD
A = AMINI( ANGL2, ACOS( (DELT-Y-Z)*TANN/XMAX) )
B = AMAX1( ANGL1, GAMMO(1,Z) )
GO TO 60
C
CASE2 -- EXIT OUTWARD
A = AMINI( ANGL2, GAMMO(2,Z) )
B = AMAX1( ANGL1, GAMMO(1,Z), ACOS( AMINI(1, (DELT-Y-Z)*TANN
, /XMAX) ) )
GO TO 60
C
INTEGRAND FOR ETAOU
C
50 A = ANGL2
B = AMAX1( ANGL1, GAMMO(2,Z) )

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60 OUT = AMAX1(0. , SIN(A)-SIN(B))
70 RETURN
END

Function GAMMO(IG, Z)

FUNCTION GAMMA(IG,Z)
ROUTINE TO CALCULATE GAMMA(I) AND GAMMA(I-1) FOR RADIATION ENTERING
THE OUTSIDE CELL WALL AT Z=I)+Z

H E L I C C O M M O N		D R A B 0 0 7		D R A B 0 0 8	
COMMON	ANN	ATOM	BCATAG	BCBTAG	BCLTAG
COMMON	ANGLN	CHRN	CHOK	CHNE	CO
COMMON	CAPIN	CAPS	CYCLE	DUGPRT	DIANTP
COMMON	CSTOP	CCVV	UT	DTUGR	DTG
COMMON	UNIN	DIR1	DIR2	DUF	EFAC
COMMON	ERRCRT	ETH	EXERO	FDC	FIFT
COMMON	FUDGE	G	HCB	HCP	HNS
COMMON	FHYUB	I	IGOTO	IH	INMAG
COMMON	ISSR	IS	ITRMAX	IU	I2
COMMON	J	JH	JMAX	JU	J5
COMMON	K	KMAX	KMAX1	KMAX2	KMAX
COMMON	MLNGE	MC	N	NK	NPC
COMMON	NH	NT	NTAPE	NI	N3
COMMON	OODC	PABOVE	PALO	PIDTS	PRINTS
COMMON	PHR	P	HADE	RFT	RR
COMMON	SCCOR	SCRE	SCYCLE	SIG	SIGB
COMMON	SPROB	SYNAX	SVS	S3	S5
COMMON	T	TAUTS	THICK	THAX	UT
COMMON	UVMAX	VAPE	VLO	VEL	WSA
COMMON	WLRANK	WLRANK	WLRANK	WLRANK	WLRANK

THE FOLLOWING VARIABLES IN HECTIC COMMON HAVE BEEN RENAMED IN THIS
 ROUTINE -- CV TO CCVV , ISEAD TO ISSSS , L TO LLLL , M TO MMMMM,

[illegible]

REAL	MERGE
/RAD1/ L2P1. LLOW. /RAD2/ COMMON L2. ISHEEP. COMMON	IAID. LMAX. LHI ANG1. /RAD2/ ANG1.2. ASN. CHKSUM.
	ICHECK. MUMAX. ISEND. ISENDD.
	KMAX. M2.
	RAD1 RAD2 RAD RAD3
	DNAB096 DNAB097 DNAB098

[illegible]

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C REAL MERGE
C
COMMON /RAD1/ IAU, KMAK, ICHECK, ISEND,
1L2, L2P1, LMAX, ME, ISENOG,
2ISELP, LLOW, LHI, ANGL2, CKSUM,
COMMON /RAD2/ ANG1, HELP1, PP,
1DELTY, FOP1, HELP2, TANN,
2P12, P10, SINP2,
3JMAX, STLF, DUM1(1),
COMMON /RAD3/ COSPH(6), CVOUT(6),
1COSPH(6), CH(6), NUI(6),
2DUM2(1), CV(6),
3W(6), M, L, LP,
COMMON /INDUX/ LCIRCK, KK2, HOUUFF,
COMMON /UNHUM/ UIS, VV(6), DH(6),
COMMON /XIS/ X(6), SUM(50,6),
1XIS(6), TH(6), RTHTA(50), ATHTA(50)
COMMON /SUMCHK/ SUM(50,6),
COMMON /UNDT/ BHTA(50), RUT(50), RUZ(100)
C
REAL MU, LMDA
C
EQUIVALENCE (PIOUT,OLDTH), (U,ROSS),
1(CAP(1),BUFF1(1),BUFF(1),NO1), (PI1),PLANCK(1),BUFF2(1),NO0),
2(UMENSION, OLUTH(1), ROSS(1), BUFF(1), BUFF(1), PLANCK(1),
3, BUFF2(1))
C
EQUIVALENCE (V(1),XA(1),EB(1))
DIMENSION XA(1)
1EB(50,6), EV(50,6), ET(50,6),
2ETA(6,6), ETAT(6,6), ETAG(6,6), ETABT(6,6),
3ETAB(6,6), ETAB(6,6), ETAGT(6,6), ETAG(6,6)
C
DIMENSION PUR(50), PUZ(100), RUR(50), RUZ(100)
C
FUO = 0.0
ETCP = 0.0
ESID = 0.0
EUTN = 0.0
IF (IMTAG.EQ.U) DIANG=1.
NVEZ=1
NTHVEZ
IF (ITAG.EQ.O) NVEZ=2
VEZ=VEZ
CALL UVCHK (KUMY)
CALCULATE GEOMETRT FACTORS
DO 10 I=1,IMAX
PUR(I)=PI/TAU(I)
DO 20 J=1,JMAX
PUZ(J)=E1/DT(J)
10

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C      20 HUZ(J+1)=2.0E10/(Y(J+1)-Y(J-1))
C      FIND HIGHEST TEMPERATURE
C      THTAMA = .025
C      DO 30 K=2,KMAX
C      IF (UMK(KFIT(K)+2).NE.1) GO TO 30
C      IF (THTA(K).LT.THTAMX) GO TO 30
C      THTAMX=THTA(K)
C      30 CONTINUE
C      THTAMX=MAX1(THTAMX,BCRTAG,BCBTAG,BCATAG)
C      NELNTY POINT FOR SECOND TEMPERATURE ITERATION
C      40 DO 50 K=2,KMAX
C      XA(K) = 0.
C      XB(K) = 0.
C      50 EN(K)=0.0
C      TNN1072
C      TNN1073
C      TNN1074
C      TNN1075
C      TNN1076
C      TNN1077
C      TNN1078
C      TNN1079
C      TNN1080
C      TNN1083
C      TNN1084
C      TNN1085
C      TNN1086
C      TNN1087
C      TNN1088
C      TNN1089
C      TNN1090
C      TNN1091
C      TNN1092
C      TNN1093
C      TNN1094
C      TNN1095
C      TNN1096
C      TNN1097
C      TNN1098
C      TNN1099
C      TNN1100
C      TNN1101
C      TNN1102
C      TNN1103
C      TNN1104
C      TNN1105
C      TNN1107
C      TNN1108
C      TNN1109
C      TNN1111
C      TNN1112
C      TNN1113
C      TNN1114
C      TNN1115
C      TNN1116
C      TNN1117
C      TNN1119
C      TNN1120
C      TNN1121
C      TNN1122
C      TNN1123
C      TNN1124
C      .....
C      DO FREQUENCY BOOKKEEPING
C      .....
C      SET UP MAX FREQ BOUNDARY
C      HNUP=1.0E6
C      HNUB=1.0E24
C      IF (INTAG.F.Q.0) GO TO 180
C      60 HNU=HNU+1
C      CALL KAPPA
C      HNU=HNU**4
C      HNU=HNU/HNU
C      C      MERGE GROUPS WITH HNU MORE THAN (MERGE) TIMES LARGEST THETA
C      IF (MERGE.GT.0.0) GO TO 70
C      S1=7.0225
C      GO TO 965
C      70 IF (THTAMX-HNU/MERGE) 100,80,80
C      80 IF (HNU-1) 90,210,120
C      90 S1=7.0235
C      GO TO 965
C      C      REJECT TAPE IF MORE THAN HALF OF GROUPS MERGE
C      100 IF (HNU/HNU-HNU) 120,110,110
C      110 U=AMUL(MERGL,1.1
C      IF (0.6T.0.4.AND.0.LT.0.0) GO TO 120
C      S1=7.0250
C      GO TO 965
C      120 DO 130 K=2,KMAX
C      IF (UMK(KFIT(K)+2).NE.1) GO TO 130
C      T=THE TALK**4
C      BETAS=U/T/THETA(K)
C      BLTAP=HNU/T/THETA(K)
C      UFU=U*HNU/T/THETA(K)

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TNN1126 IF (OFB+LT-1.E-10) GO TO 130
TNN1127 TEMP(1)=UFB*1%
TNN1128 EMB1=EXP(-BE*TA)
TNN1129 EMB2=EXP(-SE*TAP)
TNN1130 TEMP(2)=UFB+0.034974/T4*(HNU4/(1.0-EMB1)+EMB1-HNUP)/(1.0-EMB2)*E
TNN1131 TNN1129
TNN1132 TEMP(132)
TNN1133 TNN1133

C C C
      FORM NUMERATORS AND DENOMINATORS OF MERGED KAPPAS

      XA(K) = XA(K) + TEMP(1)
      XB(K) = XB(K) + TEMP(2)
      BK(K) = BK(K) + PLANK(K)*TEMP(1)
      EK(K)=EK(K)+TEMP(2)/ROSS(K)
      CONTINUE
      HNUP=HNUP
      HNU4=HNU4
      IF (THETA-X-HNU/MERGE) 60,140,140

      FORM MERGED KAPPAS

      DO 17% K=2,KMAX
      140 IF (XA(K)) 159,17%,160
      159 S1 = 7.0520
      160 GO TO 965
      ROSS(K) = XB(K)/ER(K)
      PLANK(K) = B(K)/XA(K)
      ER(K)=0.
      CONTINUE
      17% HNUP=1.0E6
      HNU4=1.0E7%
      UHNU=HNUP-HNU
      GO TO 230

C C C
      MONOFREQUENCY CALCULATION

      180 HNU=1
      1% HNU=1
      HNU=.001
      GO TO 230
      *****
      BEGIN FREQUENCY LOOP
      *****
      205 IGNU = IGNU + 1
      CALL KAPPA
      UHNU=HNUP-HNU
      HNU=HNU+.4
      CONTINUE
      210 CONTINUE
      230 CONTINUE
      270 CALL UVCHK (KDMY)
      GO TO (280,290), KDMY
      280 S1=7.0522
      GO TO 965

TNN1126 TNN1126
TNN1127 TNN1127
TNN1128 TNN1128
TNN1129 TNN1129
TNN1130 TNN1130
TNN1131 TNN1131
TNN1132 TNN1132
TNN1133 TNN1133

TNN1137 TNN1137
TNN1138 TNN1138
TNN1139 TNN1139
TNN1140 TNN1140
TNN1141 TNN1141
TNN1142 TNN1142
TNN1143 TNN1143
TNN1144 TNN1144

TNN1151 TNN1151
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TNN1161 TNN1161
TNN1166 TNN1166
TNN1167 TNN1167
*****
TNN1168 TNN1168
* TNN1169 *
TNN1170 TNN1170
* TNN1171 *
TNN1171 TNN1171
*****
TNN1174 TNN1174
TNN1175 TNN1175
TNN1176 TNN1176

TNN1205 TNN1205
TNN1206 TNN1206
TNN1207 TNN1207

TNN1209 TNN1209

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C
C      FOMR RUSSELLAND AND PLANCK OPTICAL DEPTHS
C      DOUBLE ON STORAGE FOR ABSORPTION COEFFICIENTS, MU, AND LAMBDA
C      290  DO 304 I=1,IMAX
C            W2(I)=0.
C            K=1
C            M=N-IMAX
C            DO 304 J=1,JMAX
C              FACTOR=1.
C              IF (TAU(I)*TAU(J).GT.1.E-20) FACTOR=(RUSL(K)/PLANK(K))
C              RUSL(K)=AMAX1(RUSL(K),AMAX(K)/(TAU(I)*UY(J)).1.E-20)
C              PLANK(K)=AMAX1(PLANK(K),FACTOR*AMAX(K)/(TAU(I)*UY(J)).1.E-20)
C            M=N-IMAX
C            K=K+IMAX
C          CONTINUE
C      304  *****
C      C      T W U - D I M E N S I O N A L   T R A N S P O R T   S E C T I O N *****
C      *****
C      REMIND J
C      N = HOURUFF
C      NDI = N
C      NUA = N+2
C      NSKP = N+(JMAX-1)
C      NDI = 0
C      200  DO 150 J=1,JMAX
C            DO 150 I=1,IMAX
C              K = 1 + 1 + IMAX*(J-1)
C            150  EK(K) = 0.
C      C
C      C      S T A R T   S W E E P I N G
C      C
C      DO 5679 I=1,IMAX
C      DO 5679 L=1,LMAX
C      E2(I,L) = 0.
C      E2(I,L) = 0.
C      E2(I,L) = 0.
C      5679 CONTINUE
C      DO 2000 M=M+1,M2
C      M = M+M
C      FAC = -1.
C      J = 1
C      JADD = 1
C      CALL ROTBN
C      GO TO 170
C      155  M = M+M+1
C      FAC = 1.
C      J = JMAX
C      JADD = -1
C      CALL ROTBN
C      C      I N W A R D   S W E E P

```



```

170 ISWEEP = 1
    TH2 = Y(J) - DY(J)/2.
    YB = Y(J)
    IF (M.GT. M2) YB=Y(J-1)
    I = IMAX
    CALL OUTBIO
300 K = I + 1 + IMAX(J-1)
    IACT = JNN(KFIT(K)/2)
    SP = 4. * ROSS(K) * STEF * THETA(K) * THETA(K)
    .
    XHM = X(I-1) + DX(I)/2.
    IF (NWL.LE. 0) GO TO 2200
350 IF (I.EQ. 1) GO TO 650
    DO 600 L=1,L2
        C ETAD1
        SUMHB = 0.
        DO 400 LP=1,L2
            SUMHU = SUMHB + EB(I,LP)*BUFF(IBUFF+LP)
            IBUFF = IBUFF + L2
        C ETAD1
        SUMHO = 0.
        DO 500 LP=1,L2
            SUMHO = SUMHO + EV(I,LP)*BUFF(IBUFF+LP)
            IBUFF = IBUFF + L2
        SUM = SUMHB + SUMHO
        U = BUFF(IBUFF+1)
        XAA = BUFF(IBUFF+2)
        YA = TH2 + FAC*BUFF(IBUFF+3)
        C = BUFF(IBUFF+4)
        IBUFF = IBUFF + 4
        EV(I-1,L) = TRANS(ROSS(K),D,XAA,X(I-1),YA,TH2,C,SUM,IACT,SP)
600 CONTINUE
C
650 DO 900 L=1,L2
    C ETAD1
    SUMHB = 0.
    DO 700 LP=1,L2
        SUMHB = SUMHB + EB(I,LP)*BUFF(IBUFF+LP)
        IBUFF = IBUFF + L2
    C ETAD1
    SUMHO = 0.
    DO 800 LP=1,L2
        SUMHO = SUMHO + EV(I,LP)*BUFF(IBUFF+LP)
        IBUFF = IBUFF + L2
    SUM = SUMHB + SUMHO
    D = BUFF(IBUFF+1)
    XAA = BUFF(IBUFF+2)
    YA = TH2 + FAC*BUFF(IBUFF+3)
    C = BUFF(IBUFF+4)
    IBUFF = IBUFF + 4
    ET(I,L) = TRANS(ROSS(K),D,XAA,XHM,YA,YB,C,SUM,IACT,SP)
900 CONTINUE
C
    I = I - 1
    N01 = N01 - 1

```

```

1000 IF (I - L1. 1) 60 TO 1050
60 TO 500
1050 CONTINUE
C
C O U T W A R D S W E E P
C
      LSWEEP = 2
      DO 1050 I=1,IMAX
      K = I + 1 + IMAX*(J-1)
      IACT = JMH(KFIT(K),2)
      SP = 4. * ROSS(K) * STEF * THETA(K) * THETA(K)
      * THETA(K) * THETA(K)
      XMH = X(I-1) + DX(I)/2.
      IF (NOI - LE. 0) 60 TO 2200
1085 DO 1400 L=L2P1,LMAX
C   ETAIO
      SUMI = 0.
      DO 1100 LP=L2P1,LMAX
      1100 SUMI = SUMI + EV(I-1,LP)*BUFF(IBUFF+LP-L2)
      IBUFF = IBUFF + L2
C   ETABO
      SUMHU = 0.
      DO 1200 LP=L1,LMAX
      1200 SUMHU = SUMHU + EV(I,LP)*BUFF(IBUFF+LP)
      IBUFF = IBUFF + LMAX
C   ETABO
      SUMOO = 0.
      DO 1300 LP=L1,L2
      1300 SUMOO = SUMOO + EV(I,LP)*BUFF(IBUFF+LP)
      IBUFF = IBUFF + L2
      SUM = SUMI + SUMOO + SUMHU
      U = BUFF(IBUFF+1)
      XAA = BUFF(IBUFF+2)
      YA = YH2 + FAC*BUFF(IBUFF+3)
      C = BUFF(IBUFF+4)
      IBUFF = IBUFF + 4
      EV(I,L) = TRANS(ROSS(K),D,XAA,X(I),YA,YH2,C,SUM,IACT,SP)
1400 CONTINUE
C
      DO 1800 L=L2P1,LMAX
C   ETABT
      SUMI = 0.
      DO 1500 LP=L2P1,LMAX
      1500 SUMI = SUMI + EV(I-1,LP)*BUFF(IBUFF+LP-L2)
      IBUFF = IBUFF + L2
C   ETABT
      SUMHU = 0.
      DO 1600 LP=L1,LMAX
      1600 SUMHU = SUMHU + EV(I,LP)*BUFF(IBUFF+LP)
      IBUFF = IBUFF + LMAX
C   ETABT
      SUMOO = 0.
      DO 1700 LP=L1,L2
      1700 SUMOO = SUMOO + EV(I,LP)*BUFF(IBUFF+LP)
      IBUFF = IBUFF + L2

```

```

SUM = SUM1 + SUM88 + SUM00
U = BUFF(I, BUFF+1)
XAA = BUFF(I, BUFF+2)
YA = YH2 + FAC*BUFF(I, BUFF+3)
C = BUFF(I, BUFF+4)
IBUFF = IBUFF + 4
E(I, L) = TRANS(ROSS(K), U, XAA, XHH, YA, YB, SUM, IACT, SP)
1800 CONTINUE
NOI = NOI - 1
1850 CONTINUE
C
C WRITE OUT E
1851 IF (MSGPT .LT. 1.E-20) GO TO 1870
WRITE (6, 9005) M, J
IF ((J.NE.1.OR.M.GT.M2) .AND. (J.NE.JMAX.OR.M.LE.M2)) GO TO 1854
WRITE (6, 9050)
DO 1853 I=1, IMAX
1853 WRITE (6, 9002) I, (E(I, L), L=1, LMAX)
1854 WRITE (6, 5001)
DO 1855 I=1, IMAX
1855 WRITE (6, 9002) I, (E(I, L), L=1, LMAX)
WRITE (6, 9003)
DO 1860 I=1, IMAX
1860 WRITE (6, 9002) I, (E(I, L), L=1, LMAX)
1870 CONTINUE
C
C COMPUTE ENERGY DEPOSITION RATES FOR M, J
DO 1874 I=1, IMAX
K = 1 + 1 + IMAX*(J-1)
DO 1874 L=1, LMAX
IF (L.GT. L2) GO TO 1872
HELP = EV(I-1, L)
IF (I.EQ. 1) HELP=0.
ER(K) = ER(K) - HELP + EV(I, L)
GO TO 1873
1872 HELP = EV(I-1, L)
IF (I.EQ. 1) HELP=0.
ER(K) = ER(K) + HELP - EV(I, L)
1873 ER(K) = ER(K) + EV(I, L) - ET(I, L)
1874 CONTINUE
C COMPUTE SYSTEM ENERGY GAIN THRU BOUNDARIES
IF (M.GT.M2 .OR. J.NE.1) GO TO 1905
DO 1900 I=1, IMAX
DO 1900 L=1, LMAX
1900 EBTM = EBTM + E(I, L)
1905 IF (M.GT.M2 .OR. J.NE.JMAX) GO TO 1915
IF (BCATAG .LT. 0.) GO TO 1915
DO 1910 I=1, IMAX
DO 1910 L=1, LMAX
#2(I) = #2(I) + E(I, L)/TAU(I)
1910 ETOP = ETOP - ET(I, L)
1915 IF (M.LE.M2 .OR. J.NE.JMAX) GO TO 1925
IF (BCATAG .LT. 0.) GO TO 1925
DO 1920 I=1, IMAX
DO 1920 L=1, LMAX

```

```

      K2(I) = M2(I) - E(I,L)/TAU(I)
1920 ETOP = ETOP + E(I,L)
1925 IF (M.LE.M2 .OR. J.NE.1) GO TO 1935
      DO 1930 I=1,IMAX
1930 EBTM = EBTM - E(I,L)
1935 DO 1945 L=1,LMAX
      IF (L.GT. L2) GO TO 1940
1940 ESIDE = ESIDE + EV(IMAX,L)
      GO TO 1945
1945 CONTINUE
      DO 1875 I=1,IMAX
      DO 1875 L=1,LMAX
1875 E(I,L) = E(I,L)
C
      J = J + JACU
      IF (J.LE.JMAX .AND. M.LE.M2) GO TO 170
      IF (J .LE. JMAX) GO TO 1890
      DO 1876 IJNK=1,NB1
1876 BACKSPACE 3
      GO TO 155
1890 IF (J .LT. 1) GO TO 1895
      DO 1877 IJNK=1,NB1
1877 BACKSPACE 3
      GO TO 170
1895 IF (M .LE. M2+1) GO TO 2000
      DO 1878 IJNK=1,NSKP
1878 READ(3) (BUFF(I),ILM=1,1023)
2000 CONTINUE
C
C
C
      TAKE INTO ACCOUNT ALL DIRECTIONS
      DO 2005 K=2,KMAX
2005 ER(K) = ER(K)*2.
      WRITE (6,9006)
      DO 2010 J=1,JMAX
      K1 = 2 + IMAX*(J-1)
      K2 = IMAX*J + 1
      WRITE (6,9007) J, (ER(K),K=K1,K2)
2010 CONTINUE
      EBTM = EBTM + 2.
      ETOP = ETOP + 2.
      ESIDE = ESIDE + 2.
      FOO = EBTM + ESIDE + ETOP
      NERIND 3
      GO TO 835
C
C
2200 READ (3) (BUFF(I),ILM=1,1023)
      IF (EOF.3) 2211,2212
2211 IF (ISHELP.EQ.2 .AND. M.EQ.M2 .AND. J.EQ.JMAX .AND. I.EQ.IMAX)
      . GO TO 2212
      IF (ISHELP.EQ.2 .AND. M.EQ.M2+1 .AND. J.EQ.JMAX .AND. I.EQ.IMAX)
      . GO TO 2212

```

```

STOP
2212 16UFF = 1
GO TO (350,1085) , 15WEEP

C
C
835 GROUP = HNU
      HNUP=HNU
      HNUF=HNU
      IF (1HNU-HNU) 205,850,840
840 51=7.1010
      GO TO 965
850 CONTINUE
870 WRITE (6,970) NC
      IF (ABS(UG6PRT) .LT. 1.E-20) GO TO 930
      WRITE (6,980) (ER(K),K=2,KMAX)

C
C
      ADVANCE FREQ, STORE EMERGENT FLUX, TEST FOR COMPLETION OF GROUPS
      *****
      F R E Q U E N C Y   L O O P
      *****
      ALTERNATE ON TEMPERATURE
      GO TO (890,910), NVLZ
880 NVLZ=2
      VCZ=NVLZ
      NYENNVLZ
      1-HNU=0
      DO 905 K=2,KMAX
      WORK, SOURCE TERMS OMITTED
      IF (JARKFIT(K)+2) .NE.1) GO TO 905
      OLUTH(K)=THETA(K)
      E=AI(K)+ER(N)OUT/ANX(K)
      SVETAU(1)=OUT(J)/ANX(K)
      CALL LS (SV,E,TEMP(1),TEMP(2),66)
      THETA(K)=0.5*(THETA(K)+TEMP(1))
      CONTINUE
905 IF (MFTAG.EQ.0) CALL KAPPA
      GO TO 940
      IF (1TAG.EQ.0) GO TO 930
      DO 920 K=2,KMAX
      THETA(K)=OLUTH(K)
920 H A N G E   I N T E R N A L   E N E R G I E S
      TEMP(1)=1.
930 DO 940 K=2,KMAX
      IF (JARKFIT(K)+2) .NE.1) GO TO 940
      DE=ER(K)*OUT/ANX(K)
      Q=SLUQ*AI(K)
      IF (ABS(DE) .LT. ABS(Q)) GO TO 940
      TEMP(1)=AMINI(Q/ABS(DE),TEMP(1))
940 CONTINUE
      OTEMP = OT + TEMP(1)
      IF (IDTEMP .GT. FFB) GO TO 950
      WRITE (6,970) NC
      *****
      T M N 1 3 8 1
      T M N 1 3 8 2
      T M N 1 3 8 4
      T M N 1 3 9 0
      T M N 1 3 9 1
      T M N 1 3 9 2
      T M N 1 3 9 3
      T M N 1 3 9 4
      T M N 1 3 9 5
      * T M N 1 3 9 6
      * T M N 1 3 9 7
      * T M N 1 3 9 8
      T M N 1 3 9 9
      T M N 1 4 0 0
      T M N 1 4 0 1
      T M N 1 4 0 2
      T M N 1 4 0 3
      T M N 1 4 0 4
      T M N 1 4 0 5
      T M N 1 4 0 7
      T M N 1 4 0 9
      T M N 1 4 1 0
      T M N 1 4 1 2
      T M N 1 4 1 3
      T M N 1 4 1 5
      T M N 1 4 1 6
      T M N 1 4 1 7
      T M N 1 4 1 8
      T M N 1 4 1 9
      T M N 1 4 2 0
      T M N 1 4 2 1
      T M N 1 4 2 2
      T M N 1 4 2 3
      T M N 1 4 2 4
      T M N 1 4 2 5
      T M N 1 4 2 7
      T M N 1 4 2 8
      T M N 1 4 3 1

```

TRN1432
TRN1433
TRN1434

```
950 WHITE (6,980) (ER(K),K=2,KMAX)
    WHITE (6,990) DE,K,AIX(K),AMX(K)
    SI=7.1000
    GO TO 965
    T = T - DT
    DT = UTEMP
    T = T + UTEMP
    FOU = FOU + DT
    ETH = ETH + FOU
    UACC = BACC + EBTW*DT
    SACC = SACC + ESID*DT
    TACC = TACC + ETOP*DT
    FUOPLTH,2 ARE IN ENUS
    WHITE (6,1001) NC,DT
    WHITE (6,1010) BACC,SACC,TACC,FOU
    HEAD (4) FIOUT,CAP,P,U,V
    HEWIND,4
    RETURN
    SI = 7.1111
    967 HEAD (4) FIOUT,CAP,P,U,V
    965 ISENG = 2
    CALL EDIT
```

TRN1451

```
C 970 FJMMAT (52MINUT) INITIAL ENERGT GAIN PER UNIT TIME DUE TO RADIATION.8H
    1 CYCLE =15)
    990 FJMMAT (1X,1P10E11.4)
    990 FJMMAT (5H DE =1PE11.4,5H K =110.6H AIX =E11.4,6H AMX =E11.4)
    1001 FJMMAT (1H0 FLUX FOR CYCLE 14,7H DT =1PE10.3)
    1010 FJMMAT (9H BOTTOM =1PE11.4,5H SIDE =E11.4,7H TOP =E11.4,7H FOU TRN1458
    1PE11.4)
    9000 FJMMAT (7H1FOR M=,12.5H J=,12.24H ER FOR INWARD SWEEP ARE)
    9001 FJMMAT (1H0//10X,7HEV(1.1)/50X,1HL/)
    9002 FJMMAT (3H I=,12/16X,1PBE10.6))
    9003 FJMMAT (1H0//10X,7HET(1.1)/50X,1HL/)
    9005 FJMMAT (7H1FOR M=,12.5H J=,12.17H ENERGT RATES ARE)
    9006 FJMMAT (28MINET ENERGT DEPOSITION RATES//50X,1HI/)
    9007 FJMMAT (3H J=,12/16X,1PBE10.6))
    9012 FJMMAT (9H1 BOTTOM=,1PE10.6,5X,5MSIDE=,1PE10.6,5X,5HTOP=,1PE10.6,5
    .X,4HFOO=,1PE10.6)
    9050 FJMMAT (1H0//10X,7HEB(1.1)/50X,1HL/)
    9010 FJMMAT (1H0//1X,1P10E11.4))
    END
```

TRN1460

Function TRANS(SIGMA,D,XA,XB,YA,YB,
C,SUM,IAC,SP)

```

C  FUNCTION TRANS(SIGMA,D,XA,XB,YA,YB,C,SUM,IAC,SP)
C  TRANSPORT FROM (XA,YA) TO (XB,YB)
C  FORM 1 = 12-566371
DATA  FAC/2./
IF (IAC.NE.1) GO TO 100
A1 = SIGMA * D
A2 = 12-566371 * SIGMA
A3 = A2 * SIGMA
A4 = EXP(-A1)
A5 = 1. - A4
SA = S(XA,YA)
SB = S(XB,YB)
HELP = AMIN1(SA,SB,SP)
IF ( ( AMAX1(SA,SB,SP)-HELP)/HELP .LT. FAC ) GO TO 90
SA = SP
SB = SP
90 TRANS = SUM*AC + C*( SA/A2*AS +(SB-SA)/A3*(A1-AS)/D)
RETURN
C  TRANSFER THRU AN INACTIVE ZONE
100 TRANS = SUM
RETURN
END

```


290

```

C RELATIONSHIPS --
C (STEFANS CONST.) = AC/4 = 1.0283E+12
C ENG5/(CM*2 SEC EV*cm)
C SOURCE = 5 = 4*ROSS*(STEFANS CONST.)*(THETA000)
C ENG5/(CM*2 SEC)
C INTENSITY OF RADIATION EMITTED BY A BLACK SURFACE OF TEMP.
C THETA INTEGRATED OVER ALL FREQUENCIES
C = U = S/(4.*PI*ROSS)
C = (STEFANS CONST.)*(THETA000)/PI
C
C COMMON /RAD1/ LAID, KMAX, ENG5/(CM*2 SEC STERADIAN)
C /L2P1/ LMAX, M2, ICHECK, ISENO,
C /LLOW/ LMI, MUMAX, ISEMOO,
C COMMON /RAD2/ ANGL1, ANGL2, ASH, CMKSUM,
C /FURP1/ HELP1, LMDA, PP,
C /P14/ SINP2, TANN, TOLR, XMIN,
C /STLF/ DUM1(1), AA(16),
C /RAD3/ CUPH1(6), CVIN(6), CYOUT(6),
C /CH(6)/ MU(6), PH(6),
C /DUM2(1)/ SINLP(6), SHU(6),
C /W(6)/ CV(6)
C /INDEX/ M,L,L,P
C /LHUMM/ LCHECK, KK2, NOBUFF
C /DIS/ DV(6), YV(6), DH(6),
C /RM(6)/ TH(6)
C /SUMCHK/ SUMW(50*6), RTHETA(50), ATHETA(50)
C /BNDT/
C
C HEAL MU, LMDA
C
C EQUIVALENCE (FIOUT,OLUTH), (U,ROSS),
C (CAP1(1),BUFF1(1),BUFF(1),NO1), (P(1),PLANCK(1),BUFF2(1),NO0),
C DIMENSION OLDTH(1), MOSS(1), BUFF(1), PLANCK(1),
C , BUFF2(1)
C
C EQUIVALENCE (V(1),XA(1),EB(1))
C DIMENSION XA(1)
C DIMENSION EB(50*6), EV(50*6), ET(50*6),
C , ETAIO(6*6), ETAIT(6*6), ETAB(6*6), ETABT(6*6),
C , ETAD0(6*6), ETAO1(6*6), ETAOT(6*6), ETAO0(6*6)
C
C THETA IN EV FROM KAPPA
C JLOW = J-1
C IF (YT .LT. Y(J)-UT(J)/2.) GO TO 100
C JLOW = J
C 100 ILFT = I-1
C IF (XR .LT. X(1)-UX(1)/2.) GO TO 200
C ILFT = I
C 200 KLL = ILFT + 1 + IMAX(JLOW-1)
C KUL = ILFT + 1 + IMAX(JLOW)
C DNP = ( DX(ILFT+1) + DX(ILFT) )/2.
C DYP = ( DY(JLOW+1) + DY(JLOW) )/2.

```

RAD1
RAD2
RAD
RAD3
RAD4
RAD5
RAD6
RAD7
RAD8
RAD9
RAD10
RAD11
RAD12
RAD13
RAD14
RAD15
RAD16
RAD17

```

YUP = Y(JLOW+1) - DY(JLO+1)/2.
YUN = Y(JLOW) - DY(JLO)/2.
XNGT = X(ILFT+1) - DX(ILFT+1)/2.
XLFT = X(ILFT) - DX(ILFT)/2.
IF (JLOW.EQ.0 .AND. ILFT.EQ.IMAX) GO TO 700
IF (JLOW.EQ.JMAX .AND. ILFT.EQ.IMAX) GO TO 600
IF (JLOW.EQ.JMAX .AND. ILFT.EQ.0) GO TO 575
IF (JLOW.EQ.0 .AND. ILFT.EQ.0) GO TO 550
IF (ILFT.EQ.IMAX) GO TO 500
IF (ILFT.EQ.0) GO TO 450
IF (JLOW.EQ.JMAX) GO TO 400
IF (JLOW.EQ.0) GO TO 300
C NORMAL BILINEAR INTERPOLATION
DELT = DXP * DYP
ALL = (YUP-YI)*(XNGT-XX)/DELT
ALK = (YUP-YI)*(XX-XLFT)/DELT
AUL = (YY-YUN)*(XNGT-XX)/DELT
AUR = (YY-YUN)*(XX-XLFT)/DELT
S = SRC(KLL)*ALL + SRC(KLL+1)*ALR + SRC(KUL)*AUL + SRC(KUL+1)*AUR
RETURN
C LOWER BOUNDARY
300 S = SRC(KUL) + ( SRC(KUL+1)-SRC(KUL) )/DXP * (XX-XLFT)
RETURN
C UPPER BOUNDARY
400 S = SRC(KLL) + ( SRC(KLL+1)-SRC(KLL) )/DXP * (XX-XLFT)
RETURN
C AXIS
450 S = SRC(KLL+1) + ( SRC(KUL+1)-SRC(KLL+1) )/DYP * (YY-YUN)
RETURN
C SIDE BOUNDARY
500 S = SRC(KLL) + ( SRC(KUL)-SRC(KLL) )/DYP * (YY-YUN)
RETURN
C LOWER INSIDE CORNER
550 S = SRC(KUL+1)
RETURN
C UPPER INSIDE CORNER
575 S = SRC(KLL+1)
RETURN
C UPPER OUTSIDE CORNER
600 S = SRC(KLL)
RETURN
C LOWER OUTSIDE CORNER
700 S = SRC(KUL)
RETURN
END

```

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```

COMMON
1  /RADJ/      DUM1(1),      AA(6),
   CUSPHI(6),  CVIN(6),      CVOUT(6),
   CH(6),      MU(6),      PHI(6),
2  DUM2(1),    SINLP(6),    SMU(6),
3  b(6),      CV(6)
4  /INDEX/     M-L-LP
COMMON        /LCHK/      KK2,      NOBUFF
COMMON        /DIS/      DV(6),      YV(6),      DH(6),
COMMON        XH(6),      YH(6)
1  COMMON      /SUNCNK/    SUMB(50,6),  RTHETA(50),  ATHETA(50)
COMMON        /BNDT/      BTHETA(50),
C  REAL        MU , LMDA
C
EQUIVALENCE (FIOUT,OLDTH) , (U,ROSS) ,
(CAP(1),BUFF1(1),BUFF1(1),NOI) , (P(1),PLANCK(1),BUFF2(1),NOO)
DIMENSION OLDTH(1) , ROSS(1) , BUFF1(1) , BUFF(1) , PLANCK(1) ,
. BUFF2(1)
C
EQUIVALENCE (V(1),XA(1),EB(1))
DIMENSION XA(1)
. EB(50,6) , EV(50,6) , ET(50,6) ,
. ETAIO(6,6) , ETAT(6,6) , ETAB(6,6) , ETABT(6,6),
. ETABO(6,6) , ETAOI(6,6) , ETAOT(6,6) , ETAOO(6,6)
C
C
C  IF (JMK(KF1(KKK),2) .EQ. 1) GO TO 50
C  INACTIVE CELL
SRC = 0.
RETURN
C  ACTIVE CELL
50  SRC = n. * ROSS(K) * STEF * THETA(KKK) * THETA(KKK)
. THETA(KKK) * THETA(KKK)
C  IF (MFTAS .EQ. 0) GO TO 100
MULTIFREQUENCY CASE
DFB = PLANKUT (HNU/THETA(KKK),HNU/THETA(KKK))
SRC = SRC * DFB
100 RETURN
END

```

295

[illegible]

297

```

RAD6
RAD7
RAD8
RAD9
RAD10
RAD11
RAD12
RAD13
RAD14
RAD15
RAD16
RAD17

JXMAX,
COMMON
1  STEF
2  /RAD3/
3  COSPHI(6),
4  CH(6),
5  DUM2(1),
6  W(6),
7  /INDEX/
8  M,L,LP
9  /LCHK/
10 /LCHK/
11 /LCHK/
12 /LCHK/
13 /LCHK/
14 /LCHK/
15 /LCHK/
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AA(16),
CVOUT(6),
PHI(6),
SMU(6),
DUM1(1),
CVIN(6),
MU(6),
SINLP(6),
CV(6),
KK2,
NOUFF,
XV(6),
YV(6),
DH(6),
SUMB(50,6),
RTHETA(50),
ATHETA(50)

REAL MU, LMUA

EQUIVALENCE (FIOUT,OLDTH), (U,ROSS),
(CAP(1),BUFF1(1),NOI), (P(1),PLANCK(1),BUFF2(1),NOO)
DIMENSION OLDTH(1), ROSS(1), BUFF1(1), BUFF(1), PLANCK(1),
BUFF2(1)

EQUIVALENCE (V(1),XA(1),EU(1))
DIMENSION XA(1)
DIMENSION
ETAB0(6,6), ETAB1(6,6), ETAB2(6,6), ETAB3(6,6),
ETAB4(6,6), ETAB5(6,6), ETAB6(6,6), ETAB7(6,6),
ETAB8(6,6), ETAB9(6,6), ETAB10(6,6), ETAB11(6,6),
ETAB12(6,6), ETAB13(6,6), ETAB14(6,6), ETAB15(6,6),
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ETAB96(6,6), ETAB97(6,6), ETAB98(6,6), ETAB99(6,6),
ETAB100(6,6)

REAL JO

EV(1MAX,L) = THE RATE AT WHICH ENERGY IS TRANSFERRED ACROSS FACE
X(1MAX) IN FAN L FOR A PARTICULAR Y(J)
DO 1000 L=1,L2
BINT = STEF/P1 * RTHETA(J) * RTHETA(J)
1000 EV(1MAX,L) = BINT * JO(J,1MAX,L,M)
RETURN
END

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NONEQUILIBRIUM DIFFUSION ROUTINES

Subroutine TDRAD (ADI)

Modifications to Subroutine TDRAD
(ADI) for OLIPHANT

Modifications to Subroutine TDRAD
(ADI) for SOR

Subroutine SPEW

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Subroutine TDRAD (ADI)

TDH2000	TDH2001	TDH2002	TDH2003	TDH2004	TDH2005	TDH2006	TDH2007	TDH2008	TDH2009	TDH2010	TDH2011	TDH2012	TDH2013	TDH2014	TDH2015	TDH2016	TDH2017	TDH2018	TDH2019	TDH2020	TDH2021	TDH2022	TDH2023	TDH2024	TDH2025	TDH2026	TDH2027	TDH2028	TDH2029	TDH2030	TDH2031	TDH2032	TDH2033	TDH2034	TDH2035	TDH2036	TDH2037	TDH2038	TDH2039	TDH2040	TDH2041	TDH2042	TDH2043	TDH2044	TDH2045	TDH2046	TDH2047	TDH2048	TDH2049	TDH2050							
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COMMON	AMN	AMBIEN	ANN	ATON	BCATAG	BCBTAG	BCLTAG	HECCOM01	HECCOM02	HECCOM03	HECCOM04	HECCOM05	HECCOM06	HECCOM07	HECCOM08	HECCOM09	HECCOM10	HECCOM11	HECCOM12	HECCOM13	HECCOM14	HECCOM15	HECCOM16	HECCOM17	HECCOM18	HECCOM19	HECCOM20	HECCOM21	HECCOM22	HECCOM23	HECCOM24	HECCOM25	HECCOM26	HECCOM27	HECCOM28	HECCOM29	HECCOM30	HECCOM31	HECCOM32	HECCOM33	HECCOM34	HECCOM35	HECCOM36	HECCOM37	HECCOM38	HECCOM39	HECCOM40	HECCOM41	HECCOM42	HECCOM43	HECCOM44	HECCOM45	HECCOM46	HECCOM47	HECCOM48	HECCOM49	HECCOM50
1BCRTAG	CAPIN	CAPS	CAUT	CMK	CMXOM	CMDE	CO	HECCOM01	HECCOM02	HECCOM03	HECCOM04	HECCOM05	HECCOM06	HECCOM07	HECCOM08	HECCOM09	HECCOM10	HECCOM11	HECCOM12	HECCOM13	HECCOM14	HECCOM15	HECCOM16	HECCOM17	HECCOM18	HECCOM19	HECCOM20	HECCOM21	HECCOM22	HECCOM23	HECCOM24	HECCOM25	HECCOM26	HECCOM27	HECCOM28	HECCOM29	HECCOM30	HECCOM31	HECCOM32	HECCOM33	HECCOM34	HECCOM35	HECCOM36	HECCOM37	HECCOM38	HECCOM39	HECCOM40	HECCOM41	HECCOM42	HECCOM43	HECCOM44	HECCOM45	HECCOM46	HECCOM47	HECCOM48	HECCOM49	HECCOM50
2CUE	CSTOP	CV	CYLE	DGPRT	DTH	DTH	DTH	HECCOM01	HECCOM02	HECCOM03	HECCOM04	HECCOM05	HECCOM06	HECCOM07	HECCOM08	HECCOM09	HECCOM10	HECCOM11	HECCOM12	HECCOM13	HECCOM14	HECCOM15	HECCOM16	HECCOM17	HECCOM18	HECCOM19	HECCOM20	HECCOM21	HECCOM22	HECCOM23	HECCOM24	HECCOM25	HECCOM26	HECCOM27	HECCOM28	HECCOM29	HECCOM30	HECCOM31	HECCOM32	HECCOM33	HECCOM34	HECCOM35	HECCOM36	HECCOM37	HECCOM38	HECCOM39	HECCOM40	HECCOM41	HECCOM42	HECCOM43	HECCOM44	HECCOM45	HECCOM46	HECCOM47	HECCOM48	HECCOM49	HECCOM50
3UMIN	UNN	DT	DTUGR	DT	DTUGR	DT	DT	HECCOM01	HECCOM02	HECCOM03	HECCOM04	HECCOM05	HECCOM06	HECCOM07	HECCOM08	HECCOM09	HECCOM10	HECCOM11	HECCOM12	HECCOM13	HECCOM14	HECCOM15	HECCOM16	HECCOM17	HECCOM18	HECCOM19	HECCOM20	HECCOM21	HECCOM22	HECCOM23	HECCOM24	HECCOM25	HECCOM26	HECCOM27	HECCOM28	HECCOM29	HECCOM30	HECCOM31	HECCOM32	HECCOM33	HECCOM34	HECCOM35	HECCOM36	HECCOM37	HECCOM38	HECCOM39	HECCOM40	HECCOM41	HECCOM42	HECCOM43	HECCOM44	HECCOM45	HECCOM46	HECCOM47	HECCOM48	HECCOM49	HECCOM50
4UTH	UTH1	DT2	DT	DT	DT	DT	DT	HECCOM01	HECCOM02	HECCOM03	HECCOM04	HECCOM05	HECCOM06	HECCOM07	HECCOM08	HECCOM09	HECCOM10	HECCOM11	HECCOM12	HECCOM13	HECCOM14	HECCOM15	HECCOM16	HECCOM17	HECCOM18	HECCOM19	HECCOM20	HECCOM21	HECCOM22	HECCOM23	HECCOM24	HECCOM25	HECCOM26	HECCOM27	HECCOM28	HECCOM29	HECCOM30	HECCOM31	HECCOM32	HECCOM33	HECCOM34	HECCOM35	HECCOM36	HECCOM37	HECCOM38	HECCOM39	HECCOM40	HECCOM41	HECCOM42	HECCOM43	HECCOM44	HECCOM45	HECCOM46	HECCOM47	HECCOM48	HECCOM49	HECCOM50
5E11	EMCRT	ETH	HC	HC	HC	HC	HC	HECCOM01	HECCOM02	HECCOM03	HECCOM04	HECCOM05	HECCOM06	HECCOM07	HECCOM08	HECCOM09	HECCOM10	HECCOM11	HECCOM12	HECCOM13	HECCOM14	HECCOM15	HECCOM16	HECCOM17	HECCOM18	HECCOM19	HECCOM20																														

[illegible]

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C C C      SET UP MAX FREQ BOUNDARY
C C C      HNU*2=1.0E0
C C C      HNU*4=1.0E2*
C C C      IF (HNU*6.E0.0) GO TO 160
C C C      MULTIFREQUENCY CALCULATION
C C C
C C C      60 HNU=HNU*1
C C C      CALL KAPPA
C C C      HNU=HNU**4
C C C      HNU=HNU*HNU
C C C
C C C      MERGE GROUPS WITH HNU MORE THAN (MERGE) TIMES LARGEST THETA
C C C
C C C      IF (HNU*GT.0.0) GO TO 70
C C C      S1 = 0.0225
C C C      GO TO 8*5
C C C      70 IF (THETA*HNU/MERGE) 100.00.80
C C C      80 IF (HNU-1) 90.216.120
C C C      90 S1 = 0.0235
C C C      GO TO 8*5
C C C
C C C      REJECT TAPE IF MORE THAN HALF OF GROUPS MERGE
C C C
C C C      100 IF (HNU*HNU-HNU) 120.110.110
C C C      110 IF (AMU*(MERGE,1.).E0.0.5) GO TO 120
C C C      S1 = 0.0250
C C C      GO TO 8*5
C C C
C C C      FORM NUMERATORS AND DENOMINATORS OF MERGED KAPPAS
C C C      STORE IN XA, XB, XC, XALP
C C C
C C C      120 GO 130 K=2,KMAX
C C C      IF (HNU*(F1(K),2).NE.1) GO TO 130
C C C      T4=THETA(K)**4
C C C      U*2=HNU/THETA(K)
C C C      B*2=HNU/THETA(K)
C C C      D*2=PLMKUT(B*2*U*2)
C C C      IF (U*4.E0.0) GO TO 130
C C C      TEMP(1)=D*2*U*4
C C C      EMU1=EXP(-BETA)
C C C      EMU2=EXP(-BETA*P)
C C C      TEMP(2)=D*2*U*4*0.0384974/T4*(HNU*(1.0-EMU1)+EMU1-HNU*P/(1.0-EMU2)+EMU2)
C C C      130
C C C      XA(K)=XA(K)+TEMP(1)
C C C      XB(K)=XB(K)+TEMP(2)
C C C      XC(K)=XC(K)+PLANK(K)*TEMP(1)
C C C      XALP(K)=XALP(K)+TEMP(2)/ROSS(K)
C C C      130 CONTINUE
C C C      HNU*2=HNU
C C C      IF (THETA*HNU/MERGE) 60.140.140

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FORM MERGED KAPPAS

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140 DO 170 K=2,KMAX
   IF (XA(K)) 150,170,160
150 S1 = 0.0320
   GO TO 845
160 HOSS(K)=XB(K)/XALP(K)
   PLANK(K)=XC(K)/XA(K)
170 CONTINUE
   HNU=1.0E6
   HNU4=1.0E24
   DHNU=HNU-HNU
   GO TO 230

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MONOFREQUENCY CALCULATION

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180 HNU=1
   HNU=1
   HNU=.001
   DO 190 K=2,KMAX
   IF (JH(KFIT(K),2).NE.1) GO TO 190
   B(K)=THETA(K)**4
190 CONTINUE
   GO TO 230

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TYPICAL GROUP CALCULATION OF SOURCES

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200 IHNU=IHNU+1
   CALL KAPPA
   DHNU=HNU-HNU
   HNU4=HNU**4
210 DO 220 K=2,KMAX
   IF (JH(KFIT(K),2).NE.1) GO TO 220
   DF=PLNKUT(IHNU/THETA(K),HNU/THETA(K))
   B(K)=DF*BTHETA(K)**4
220 CONTINUE

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SET HLACBODY CONJITIONS

INNER CYLINDRICAL RADIUS NOT ASSUMED ZERO

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230 BUL=0.
   IF (BCLTAG.LE.0.0) GO TO 240
   UF=PLNKUT(IHNU/BCLTAG,HNU/BCLTAG)
   BUL=BCLTAG**4*UF
240 BURE=0.
   IF (BCRTAG.LE.0.0) GO TO 250
   UFB=PLNKUT(IHNU/BCRTAG,HNU/BCRTAG)
   BURE=BCRTAG**4*UFB
250 BUB=0.
   IF (BCBTAG.LE.0.0) GO TO 260
   UFB=PLNKUT(IHNU/BCBTAG,HNU/BCBTAG)
   BUB=BCBTAG**4*UFB
260 BUA=0.
   IF (BCATAG.LE.0.0) GO TO 290

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TUR2244
TUR2245
TUR2246
TUR2247
TUR2248
TUR2249
TUR2250
TUR2251
TUR2252
TUR2253
TUR2254
TUR2255
TUR2256

DUE=PLANKUT (IMU/BCATAG,MMUP/BCATAG)
BBA=BCATAG**DFU

      FORM ROSSELAND AND PLANK OPTICAL DEPTHS

290 DO 340 I=1,IMAX
    K=1+I
    MEK=IMAX
    W2(I)=0.0
    DO 340 J=1,JMAX
      FACTOR=1.
      IF (HPLAG,ME,0) FACTOR=ROSS(K)/PLANK(K)
      ROSS(K)=AMAX1(ROSS(K),AMX(K)/(TAU(I)*DY(J)),1.E-20)
      PLANK(K)=AMAX1(PLANK(K),FACTOR*AMX(K)/(TAU(I)*DY(J)),1.E-20)
    END DO
  END DO

      FORM LAMUUA, MEAN FREE PATH AT CELL EDGES

      IF (1.EQ.1) GO TO 310
      DTAU=ROSS(K)*DX(I)
      DTAU=ROSS(K-1)*DX(I-1)
      IF (ABS(DTAU-DTAU)/(DTAU+DTAU).LE.000C) GO TO 300
      ALAMV(K)=AMAX1(1./ROSS(K),1./ROSS(K-1))
      GO TO 310
300 ALAMV(K)=(DX(I)*DX(I-1))/(DTAU+DTAU)
310 IF (J.EQ.1) GO TO 330
      DTAU=ROSS(K)*DY(J)
      DTAU=ROSS(K)*DY(J-1)
      IF (ABS(DTAU-DTAU)/(DTAU+DTAU).LE.000C) GO TO 320
      ALAMH(K)=AMAX1(1./ROSS(K),1./ROSS(M))
      GO TO 330
320 ALAMH(K)=(DY(J)*DY(J-1))/(DTAU+DTAU)
330 MEH=IMAX
    MEK=IMAX
  340 CONTINUE

      BEGIN NONEQUILIBRIUM DIFFUSION TREATMENT

      FORM MATRIX ELEMENTS XA, XB, XC, XD, XALP, XGAM
      BOUNDARY CONDITION MODIFIES XB AND XD.
      THIS CODING ASSUMES DX(1) = X(1+1) - X(1), SAME FOR Y.

      DO 440 I=1,IMAX
        K=1+I
        MEK=IMAX
        DO 440 J=1,JMAX
          XA(K)=0.
          IF (J.EQ.1) GO TO 350
          XA(K)=PUM(I)-PUM(I+1)*ALAMV(K)
          GO TO 360
350 XA(K)=0.0
          INNER CYLINDRICAL RADIUS ASSUMED ZERO
360 IF (1.EQ.IMAX) GO TO 370
          XC(K)=PUM(I)*PUM(I+1)*ALAMV(K+1)
          GO TO 380

```

```

370 AC(K)=0.0
   IF (BCATAG.LT.0.0) GO TO 380
   METAB(1)=1.0/12.-EXP(-OY(1)*SORT(0.75*ROSS(K)*PLANK(K)))
   XU(K)=XU(K)+PUZ(1)*3.E10*OCTAB(1)
   XD(K)=XD(K)-4.104E12*PUZ(1)*X(1)*OOR
380 IF (J.EQ.1) GO TO 390
   XALP(K)=PUZ(1)*NUZ(1)*ALAMH(K)
   GO TO 400
390 XALP(K)=0.0
   IF (BCATAG.LT.0.0) GO TO 400
   METAB(1)=1.0/12.-EXP(-OY(1)*SORT(0.75*ROSS(K)*PLANK(K)))
   XU(K)=XU(K)+PUZ(1)*1.5E10*OCTAB(1)
   XD(K)=XD(K)-PUZ(1)*2.052E12*OOR
400 IF (J.EQ.JMAX) GO TO 410
   XGAM(K)=PUZ(1)*NUZ(1)*ALAMH(K)
   GO TO 420
410 XGAM(K)=0.0
   IF (BCATAG.LT.0.0) GO TO 420
   METAB(1)=1.0/12.-EXP(-OY(1)*SORT(0.75*ROSS(K)*PLANK(K)))
   XU(K)=XU(K)+PUZ(1)*1.5E10*OCTAB(1)
   XD(K)=XD(K)-PUZ(1)*2.052E12*OOR
420 XU(K) = XU(K) + 3.E10*PLANK(K)
   XD(K)=XD(K)-4.104E12*PLANK(K)*B(K)

   INITIALIZE ERAD
   ERAD(K)=U(K)*137.
430 K=K+IMAX
   KEN=IMAX
440 CONTINUE

OPTIONAL PRINTOUT
IF (AMOD(CYCLE,PRINTL).NE.0) GO TO 442
WRITE (6,870) CYCLE,DT
CALL SPE(XA,IMAX,JMAX,2HXA)
CALL SPE(XB,IMAX,JMAX,2HXB)
CALL SPE(XC,IMAX,JMAX,2HXC)
CALL SPE(XD,IMAX,JMAX,2HXD)
CALL SPE(XGAM,IMAX,JMAX,4HIGAM)
CALL SPE(XALP,IMAX,JMAX,4HALP)
CALL SPE(ROSS,IMAX,JMAX,4HROSS)
CALL SPE(PLANK,IMAX,JMAX,6HPLANK)
CALL SPE(B,IMAX,JMAX,1HB)

START ADI SOLUTION
442 ERADM = 0.
   ERRORH = 0.
   DO 460 J=1,JMAX
   CALCULATE XDI AND XGI FOR VERTICAL SWEEP
   M = (J-1)*IMAX + 2
   BEE = (1. + XLAM)*XB(M) - XA(M) - XC(M)

```

TDR2257
TDR2258
TDR2259
TDR2260
TDR2261
TDR2262
TDR2263
TDR2264

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TDR2273
TDR2274
TDR2275
TDR2276
TDR2277
TDR2278
TDR2279

TDR2282
TDR2291
TDR2292
TDR2293

TDR2294

```

IF(BEE.NE.0.) GO TO 443
XG1(1) = 0.
XU1(1) = 0.
GO TO 444
443 DEE = -XD(M) - XALP(M)*(ERAD(M-1MAX) - ERAD(M))
      - XGAM(M)*(ERAD(M+1MAX) - ERAD(M))
      + XLAM*XB(M)*ERAD(M)
      XG1(1) = -XC(M)/DEE
      XU1(1) = DEE/BEE
444 DO 446 I=2,1MAX
      K = (J-1)*1MAX + I + 1
      BEE = (1. + XLAM)*XB(K) - XA(K) - XC(K)
      DEE = -XD(K) - XALP(K)*(ERAD(K-1MAX) - ERAD(K))
      - XGAM(K)*(ERAD(K+1MAX) - ERAD(K))
      + XLAM*XB(K)*ERAD(K)
      DENOM = XA(K)*XG1(I-1) + BEE
      IF(DENOM.NE.0.) GO TO 445
      XG1(I) = 0.
      XU1(I) = 0.
      GO TO 446
445 XG1(I) = -XC(K)/DENOM
      XU1(I) = (DEE - XA(K)*XD(I-1))/DENOM
446 CONTINUE

      CALCULATE EH (HALF STEP ERAU)
      EH(K) = XD(1MAX)
      1MAXM1 = 1MAX - 1
      DO 447 I=1,1MAXM1
      K = K - 1
      II = 1MAX - I
      EH(K) = XG1(II)*EH(K+1) + XU1(II)
447 CONTINUE
448 CONTINUE
      DO 500 I=1,1MAX

      CALCULATE XD1 AND XG1 FOR HORIZONTAL SWEEP
      M = I + 1
      BEE = (1. + XLAM)*XB(M) - XALP(M) - XGAM(M)
      IF(BEE.NE.0.) GO TO 470
      XG1(1) = 0.
      XU1(1) = 0.
      GO TO 480
470 DEE = -XD(M) - XA(M)*(EH(M-1) - EH(M))
      - XC(M)*(EH(M+1) - EH(M))
      + XLAM*XB(M)*EH(M)
      XG1(1) = -XGAM(M)/BEE
      XU1(1) = DEE/BEE
480 DO 485 J=2,1MAX
      K = (J-1)*1MAX + I + 1
      BEE = (1. + XLAM)*XB(K) - XALP(K) - XGAM(K)
      DEE = -XD(K) - XA(K)*(EH(K-1) - EH(K))
      - XC(K)*(EH(K+1) - EH(K))
      + XLAM*XB(K)*EH(K)

```

```

DENOM = XALP(K)*XG1(J-1) + DEE
IF (DEE.NE.0.) GO TO 482
XG1(J) = 0.
XG1(J) = 0.
GO TO 485
482 XG1(J) = - XGAM(K)/DENOM
483 XJ(J) = (DEE - XALP(K)*XG1(J-1))/DENOM
485 CONTINUE

      CALCULATE ERAD
      ERADP = ERAD(K)
      ERAD(K) = XDI(JMAX)
      ERADM = AMAX1(ERADM,ERAD(K))
      ENNOH4 = AMAX1(ENNOH4,ABS(ERAD(K) - ERADP))
      JMAXN1 = JMAX - 1
      DO 490 J=1,JMAXN1
      K = K - IMAX
      JJ = JMAX - J
      ERADP = ERAD(K)
      ERAD(K) = XG1(JJ)*ERAD(K+IMAX) + XDI(JJ)
      ERADM = AMAX1(ERADM,ERAD(K))
      ENNOH4 = AMAX1(ENNOH4,ABS(ERAD(K) - ERADP))
490 CONTINUE
500 CONTINUE

      SET NEGATIVE ERAD TO 0.
      DO 545 K=2,KMAX
      IF (ERAD(K).LT.0.) ENAD(K) = 0.
545 CONTINUE
      ENNOH = ENNOH4/ERADM

      OPTIONAL PRINTOUT

      IF (AMOD(CYCLE,PRINTS) .NE. 0.) GO TO 560
      WRITE (6,1010) ITER,ERROR
      IF (MOD(ITER,1000) .NE. 1) GO TO 560
      CALL SPEW(ERAD,IMAX,JMAX,4*HERAD)
      TEST FOR CONVERGENCE

560 ITER = ITER + 1
      IF (ITER .LE. ITRMAX) GO TO 570
      CALL SPEW(ERAD,IMAX,JMAX,4*HERAD)
      S1 = 8.0560
      GO TO 845
570 IF (ENNOH.GT.ENNOHRT) GO TO 442

      CALCULATE FLUXES
      DO 640 I=1,IMAX
      K2=1
      NER=IMAX
      DO 640 J=1,JMAX

```

TDR2336

TDR2354
TDR2355
TDR2356
TDR2357


```

IF (JLU.1) GO TO 740
XAKI=RUZ(J)*ALAM(K)*(LHAD(M)-ERAD(K))
GO TO 610
580 IF (INHTAG) 590,600,600
590 RAKI=0.0
GO TO 610
600 RAKI=2.0*PI*1.2*INHTAG*1.5*10*EMAD(K)*EMETAK(I)
FOUT=OUTRAK(I)*TAU(I)
610 IF (I.1.1) GO TO 620
RAKI=RUZ(I)*ALAM(K)*(LHAD(K-1)-EMAD(K))
GO TO 630
620 RUKI=0.0
630 RUKI=RAK
640 CONTINUE
DO 720 I=1,IMAX
K=1
NEK=IMAX
DO 720 J=1,JMAX
FAUVEAAT(I)
IF (J.NE.-JMAX) GO TO 680
IF (UCATAG) 650,660,660
650 FAUVEU=0
GO TO 670
660 FAVU=-2.052*1.2*EMAD(K)*EMETAK(I)
FOUFOO=FAVU*TAU(I)
670 F2(I)=FAVU*F2(I)
680 FRTXUK(I)
IF (I.NE.IMAX) GO TO 710
IF (UCRTAG) 690,700,700
690 FRTU=0
GO TO 710
700 FRTI=2.052*1.2*EMAD(K)*EMETAK(J)*X(IMAX)
FOUFCO=FRT*PI*DY(J)
710 ER(K)=X(K)*(XAK(K)-FAVU)*TAU(I)+(X(K)-FRT)*PI*DY(J)
KEK=IMAX
NEK=IMAX
720 CONTINUE

OPTIONAL PRINTOUT
IF (AMOUICYLE.PRINTL).NE.0.) GO TO 750
CALL SPB(XB,IMAX,JMAX,SHV FLUX)
CALL SPB(XA,IMAX,JMAX,SHV FLUX)
CALL SPB(ER,IMAX,JMAX,2*HER)

ADVANCE FREQ, STORE EMERGENT FLUX, TEST FOR COMPLETION OF GROUPS
750 HARP=HARP+HARP
HARP=HARP
IF (IMOU-HARP) 200,770,750

END FREQUENCY LOOP

760 SI = 0.1050

```

TIM2354
 TIM2355
 TIM2360
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 TIM2393
 TIM2394
 TIM2395
 TIM2396
 TIM2398

 TDR2404
 TDR2405
 TDR2406
 TDR2409
 TDR2410
 TDR2411
 TDR2413
 TDR2414
 TDR2415
 TDR2417

```

60 TO 845
ITERATE ON TEMPERATURE
770 GO TO (790,800), NVEZ
780 NVEZ=2
NVEZ=2
NVEZ=2
NVEZ=2
DO 790 K=2,KMAX
*ORK, SOURCE TERMS OMITTED
IF (JMK(KFIT(K),2).NE.1) GO TO 790
OLDTH(K)=THETA(K)
E=AI(K)*LR(K)*DT/AMX(K)
SV=TAU(1)*DT(J)/AMX(K)
CALL ES (SV,E,TEMP(1),TEMP(2),66)
THETA(K)=0.5*(THETA(K)+TEMP(1))
790 CONTINUE
IF (MTAG.EQ.0) CALL KAPPA
GO TO 40
800 IF (ITAG.EQ.0) GO TO 820
DO 810 K=2,KMAX
810 THETA(K)=OLDTH(K)

CHANGE INTERNAL ENERGIES
820 TEMP(1)=1.
DO 830 K=2,KMAX
IF (JMK(KFIT(K),2).NE.1) GO TO 830
DE=ER(K)*OT/AMX(K)
Q=SLU*AI(K)
IF (ABS(UE) .LT. ABS(Q)) GO TO 830
TEMP(1)=AMINI(Q/ABS(UE),TEMP(1))
830 CONTINUE
DTEMP = DT*TEMP(1)
IF (DTEMP .GT. FFB) GO TO 850
840 WRITE (6,960) (C,DT
CALL SPEWIER,IMAX,JMAX,2*HER)
S1 = 8.1668
845 ISENU = 2
READ (6) FIOUT,P,UVV
CALL EUIT
850 T = T - DT + DTEMP
DT = DTEMP
ETH = ETH + FOO*DT
WRITE (6,970) NC,DT
READ (4) FIOUT,P,UVV
NEWNU 4
RETURN

870 FORMAT (7H1CYCLE F6.0,7H TIME 1PE13.6,5H DT 1PE13.6/)
950 FORMAT (4H K =14.5H DE =1PE11.4,6H AMX =11.4,6H AIX =1PE11.4)
960 FORMAT (52H1INTERNAL ENERGY GAIN PER UNIT TIME DUE TO RADIATION,8HTDR2460
1 CYCLE =15.5H DT =1PE11.4)

```

TDR2419
TDR2420
TDR2421
TDR2422
TDR2423
TDR2424
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TDR2433
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TDR2438

TDR2440
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TDR2444
TDR2446
TDR2450
TDR2453
TDR2455

TDR2464

TDR2466
TDR2467
TDR2468
TDR2478
TDR2479
TDR2480
TDR2481

TDR2482
TDR2483

970 FORMAT (19H0, FLUX FOR CYCLE 19,7H, DT 217,10,3)
1010 FORMAT(1H0,9HITERATION,16,10X,21HMAX. ERROR/ MAX. ERAD/F16.8)
END

Modifications to Subroutine TDRAD (ADI) for OLIPHANT

```

FOR OLIPHANT
  REMOVE TUR2054 - TUR2055
  REPLACE TUR2278 - TUR2279 BY
420 XB(K)=XB(K)+3.E10*PLANK(K)-XA(K)*XC(K)-XALP(K)*XGAM(K)
  XD(K)=XD(K)+4.10E12*PLANK(K)*B(K)
  REPLALL TUR2294 - TUR2296 BY
  FIRST PASS -- 1. 10, 100, 4/-49 OF OLIPHANT
490 DO 495 K=2,KMAX
  NEX=IMAX
  XB(K)=XB(K)+XALP(K)*XGAM(K)-XA(K)*XC(K-1)
  XGAM(K)=XGAM(K)+CMXK/XB(K)
  XC(K)=XC(K)+CMXK/XB(K)
495 CONTINUE

  SECOND PASS -- EQS. 50, 51 OF OLIPHANT
495 DO 510 K=2,KMAX
  NEX=IMAX
  NEX=IMAX
  XH=XU(K)+XALP(K)*XC(K)+XGAM(K)+XA(K)*XGAM(K-1)+ERAD(N-1)+(CMXK-
  11.)*XB(K)/CMXK*(XGAM(K)+ERAD(N)*XC(K)+ERAD(K+1))
  XG(K)=(XH-XALP(K)*XG(K)-XA(K)*XG(K-1))/XB(K)
510 CONTINUE

  BACKWARD PASS -- EQ. 52 OF OLIPHANT
  ERADM = 0.
  ERRORH = 0.
  DO 530 L=2,KMAX
  K=KMAX+1-L
  NEX=IMAX
  ENEN=CMXKH*(XG(K)-XGAM(K)+ERAD(N)-XC(K)+ERAD(K+1))+(1.-CMXKH)*ERAD
  1(K)
  ERADM = AMAX1(ERADM,ENEN)
  ERRORH = AMAX1(ERRORH,ABS(ERAD(K) - ENEN) )
520 ERAD(K)=ENEN
530 CONTINUE

```

TUR2278
TUR2279TUR2294
TUR2295
TUR2296TUR2296
TUR2297
TUR2298
TUR2299

TUR2312

TUR2314
TUR2315
TUR2316
TUR2317
TUR2318
TUR2319

TUR2320

TUR2323
TUR2324
TUR2325
TUR2326
TUR2327TUR2335
TUR2336

Modifications to Subroutine TDRAD (ADI) for SOR

FOR SUCCESSIVE OVER RELAXATION

REPLACE TOR2054 - TOR2055 BY

TO RUN IN HEFTIC A WAY MUST BE PROVIDED TO STORE THE
EF ARRAY ON A DRUM ----- TOO LARGE TO HOLD IN CORE.

```

DIMENSION IL(50),IR(50),EF(15000),RES(1200),EP(1)
DIMENSION LBASE(50)
EQUIVALENCE (X5,EP)
INTEGER R,NMAX,RXN1
INTEGER RXIW

```

A WAY MUST BE FOUND TO CALCULATE THE PANEL BOUNDARIES

```

DATA IL/1.4,7.10,13.16,4400/, IR/3.8,11.19,17.20,4400/, LMAX/6/
DATA XLAM/1.3/

```

REPLACE TOR2278 - TOR2279 BY

```

420 XB(K) = XB(K) + 3.E10*PLANK(K) - XA(K) - XC(K) - XALP(K)
      1 - XGAM(K)
      XD(K) = XD(K) - 4.104E12*PLANK(K)*B(K)

```

TOR2276

TOR2279

REPLACE TOR2294 - TOR2336 BY

START RELAXATION SOLUTION

INITIALIZE

```

490 ENADM = 0.
ERRORH = 0.
LBASE(1) = 0

```

START LOOP ON L

```

DO 480 L=1,LMAX
IRL = IRL
INL = IRL
ILL = ILL
IW = ILL - ILL + 1
IMP1 = IW + 1
IMP2 = IW + 2
KI = IMAX - IW
NFI = IW + IMP1
NFN1 = NFI - 1
NMAX = IMAX
RXN1 = NMAX - 1
RXIW = (NMAX + 2)*IMP1
NFMAX = NMAX*NFI - IW
LB = LBASE(L)
LBASE(L+1) = LB + NFMAX + IW

```

```

IF(ITER.GT.1) GO TO 461
LOAD EF ARRAY
DO 452 KF=1,KFMAX
  EF(LB+KF) = 0.
452 CONTINUE
  K = ILL - KI
  KF = LB - IW
  DO 456 J=1,JMAX
    K = K + KI
    DO 458 I=ILL,IRL
      K = K + 1
      KF = KF + KFI
      EF(KF-IW) = XALP(X)
      EF(KF) = XB(K)
      EF(KF+IW) = XGAM(K)
      IF(I.EQ.ILL) GO TO 456
      EF(KF-1) = XA(K)
      IF(I.NE.IRL) GO TO 456
      EF(KF+1) = 0.
      GO TO 458
456 EF(KF-1) = 0.
456 EF(KF+1) = XC(K)
458 CONTINUE
DO ELIMINATION ON EF ARRAY
KF = LB - IW
DO 460 N=1,NXMI
  KF = KF + KFI
  KFN = KF
  NX = MIN(IW,IMAX-R)
  DO 460 N=1,NX
    KFN = KFN + KFI
    IF(IEF(KF).EQ.0.) STOP
    EF(KFN) = -EF(KF)/EF(KF)
    FACT = EF(KFN)
    IF(FACT.EQ.0.) GO TO 460
    DO 459 M=1,NX
      EF(KFN+M) = EF(KFN+M) + FACT*EF(KF+M)
459 CONTINUE
460 CONTINUE
COMPUTE RESIDUALS
461 K = ILL - KI
  R = 0
  DO 462 J=1,JMAX
    K = K + KI
    DO 462 I=ILL,IRL
      R = R + 1
      K = K + 1
  RES(R) = - XO(K) - XALP(K)*ERAD(K-IMAX) - XGAM(K)*ERAD(K+IMAX)
  - XA(K)*ERAD(K-1) - XB(K)*ERAD(K) - XC(K)*ERAD(K+1)
1

```

```

462 CONTINUE
DO ELIMINATION ON RESIDUALS
KFN = LB - IW
DO 464 I=1,NX
NX = MIN(IW,RMAX-R)
KFN = KFN + KFI
DO 463 I=1,NX
KFN = KFN + KFI
RES(R+N) = RES(R+N) + RES(I)*EF(KFN)
463 CONTINUE
KFN = KFN - KFI*NX
464 CONTINUE

DO BACK SUBSTITUTION
KF = LB +KFN*MAX + KFI
DO 466 I=1,RMAX
KF = KF - KFI
KR = RMAX - R + 1
NX = MIN(IW,RMAX-KR)
IF (NX.EQ.0) GO TO 467
DO 466 I=1,NX
KFN = KF + N
RES(KR) = RES(KR) - EP(KR+N)*EF(KFN)
466 CONTINUE
467 EP(KR) = RES(KR)/EF(KF)
468 CONTINUE

INCREMENT EFAD
R = 0
K = ILL - KI
DO 470 J=1,JMAX
K = K + KI
DO 470 I=ILL,IRL
R = R + 1
K = K + 1
ERADP(K) = ERAD(K)
ERAD(K) = ERAD(K) + ILAM*EP(R)
ERRORN = MAX1(ERRORN,ABS(ERAD(K) - ERADP))
ERADN = MAX1(ERADN,ERAD(K))
470 CONTINUE
480 CONTINUE
END LOOP ON L

```


Subroutine SPEW

```

PRINTS 2-D ARRAYS(I,J)
1-JMAX ... IMAX,JMAX
.
.
.
1..1 ... IMAX,1
.
.
.
SUBROUTINE SPEW(ARRAY,IMAX,JMAX,BCD)
DIMENSION ARRAY(1)
IL = 1
IR = NINT(IMAX/10)
JT = JMAX
JB = MAX0(JMAX-99,1)
10 PRINT 1000, BCD
15 PRINT 1010, (I=IL,IR)
DO 20 JJ=JB,JT
J = JT - JJ + JB
K1 = IMAX(J-1) + IL + 1
K2 = IMAX(J-1) + IR + 1
PRINT 1020, J, (ARRAY(K1),K2)
20 CONTINUE
IF(JB-EG,1) GO TO 30
JT = JB - 1
JB = MAX0(JT-99,1)
GO TO 10
30 IF(IR-EG,IMAX) GO TO 40
IL = IR + 1
IR = NINT(IMAX/10)
JT = JMAX
JB = MAX0(JMAX-99,1)
IF(JMAX-ST,20-OR,IMAX,ST,40) GO TO 10
PRINT 1030, BCD
GO TO 15
40 RETURN
1000 FORMAT(1H,AB)
1010 FORMAT(1H0,6X,10I2)
1020 FORMAT(1H,16,10E12,3)
1030 FORMAT(1H0,AB)
END

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REFERENCES

1. Lax, P. D., and B. Wendroff, "Systems of Conservation Laws," Comm. Pure Appl. Math. 13, 217 (1960).
2. Richtmyer, R. D. and K. W. Morton, Difference Methods for Initial Value Problems, 2d ed., Interscience Publishers, New York, 1967.
3. Burstein, Samuel Z., "Finite-Difference Calculations for Hydrodynamic Flows Containing Discontinuities," J. Computational Phys. 2, 198 (1967).
4. "Nuclear Explosion Interaction Studies," v. I, "Methods for Analysis of Radiative Transfer," Air Force Weapons Laboratory Report AFWL-TR-66-108, General Dynamics, General Atomic Division, Sections III and IV, to be issued.
5. Freeman, B. E., and C. G. Davis, "Fireball Phenomenology and Code Development," v. III, "SPUTTER Subroutines for Radiation Transport in Spheres," Air Force Weapons Laboratory Report AFWL-TR-65-143, General Dynamics, General Atomic Division, August 1965.
6. Carlson, B., "Numerical Solution of Neutron Transport Problems," Proc. Symp. Appl. Math. XI, 219 (1961).
7. Leiner, A. L. et al, "PILOT - A New Multiple Computer System," Assoc. Comp. Mach. 6, 315 (1963).
8. Freeman, B. E., "The Non-Equilibrium Diffusion Equation," USAEC Report LA-3377, Los Alamos Scientific Laboratory, 1965.
9. Oliphant, T. A., "An Extrapolation Procedure for Solving Linear Systems of Equations," Quart. Appl. Math. 20, 257 (1962).
10. Wachpress, E. L., "Iterative Solution of Elliptic Systems and Applications to the Neutron Diffusion Equations of Reactor Physics," Prentice-Hall, Englewood Cliffs, New Jersey, 1966.

11. Peaceman, D. W., and H. H. Rachford, Jr., "The Numerical Solution of Parabolic and Elliptic Differential Equations," J. Soc. Ind. Appl. Math. 3, 28 (1955).
12. Southwell, R. V., Relaxation Methods in Theoretical Physics, The Clarendon Press, Oxford, 1946.

UNCLASSIFIED

Security Classification

DOCUMENT CONTROL DATA - R & D

(Security classification of title, body of abstract and indexing annotation must be entered when the overall report is classified)

1. ORIGINATING ACTIVITY (Corporate author)

Gulf General Atomic Incorporated
San Diego, California 92112

2a. REPORT SECURITY CLASSIFICATION

UNCLASSIFIED

2b. GROUP

3. REPORT TITLE

NUCLEAR EXPLOSION INTERACTION STUDIES, Volume II, Two-Dimensional Code Development

4. DESCRIPTIVE NOTES (Type of report and inclusive dates)

29 September 1966-27 October 1967

5. AUTHOR(S) (First name, middle initial, last name)

J. R. Triplett et al.

6. REPORT DATE

April 1968

7a. TOTAL NO. OF PAGES

332

7b. NO. OF REFS

12

8a. CONTRACT OR GRANT NO.

F29601-67-C-0014

b. PROJECT NO.

5710

c. Subtask No.

07.017

ARPA Order No. 313

d.

9a. ORIGINATOR'S REPORT NUMBER(S)

AFWL-TR-67-131, Vol II

9b. OTHER REPORT NO(S) (Any other numbers that may be assigned this report)

Contractor's report No. GA-7764, Vol II

10. DISTRIBUTION STATEMENT This document is subject to special export controls and each transmittal to foreign governments or foreign nationals may be made only with prior approval of AFWL (WLRT), Kirtland AFB, NM, 87117. Distribution is limited because of the technology discussed in the report.

11. SUPPLEMENTARY NOTES

12. SPONSORING MILITARY ACTIVITY

AFWL (WLRT)

Kirtland AFB, NM 87117

13. ABSTRACT

(Distribution Limitation Statement No. 2)

The HECTIC code is a two-dimensional Eulerian radiative hydrodynamics code designed for nuclear and laser phenomenology applications involving primarily vapor-phase materials, with provisions for heat conduction and vaporization of condensed matter. In this report, a formulation is presented for an improved method of solving the hydrodynamic equations. Solution of the radiative transfer equations by the long-characteristic method is discussed, and computer codes utilizing this approach are presented. The short-characteristic method is also discussed. The nonequilibrium diffusion approximation in two dimensions is considered, and a study of algorithms for solving the implicit difference equations which arise is reported. Experimental codes utilizing the nonequilibrium diffusion and short-characteristic methods are presented.

DD FORM 1473
1 NOV 66

UNCLASSIFIED

Security Classification

14	KEY WORDS	LINK A		LINK B		LINK C	
		ROLE	WT	ROLE	WT	ROLE	WT
	Radiative hydrodynamics codes Two-dimensional radiative transfer codes Two-dimensional hydrodynamics equations						